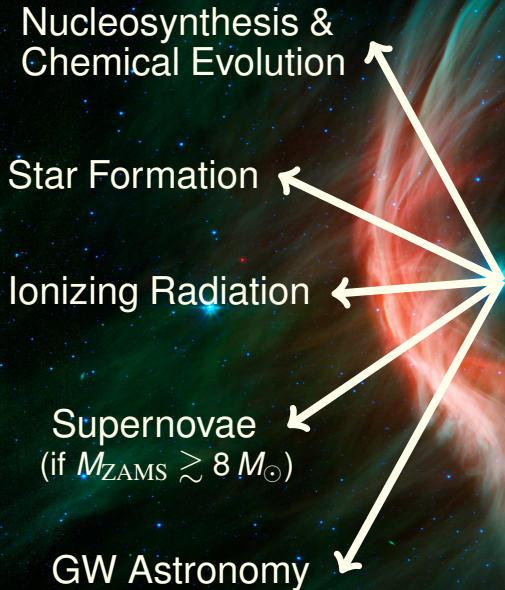


Mathieu Renzo  
PhD in Amsterdam

# Massive stars and binaries: why & how?

# Why are Massive Stars Important?



# Why are Massive Stars Important?

Nucleosynthesis &  
Chemical Evolution

Star Formation

Ionizing Radiation

Supernovae  
(if  $M_{ZAMS} \gtrsim 8 M_{\odot}$ )

GW Astronomy

**~ 70% of O type stars are  
in close binaries**

(e.g. Mason *et al.* '09, Sana & Evans '11,  
Sana *et al.* '12, Kiminki & Kobulnicky '12,  
Kobulnicky *et al.* '14)

**~ 10% of O type stars are  
runaways!**

(e.g. Blaauw '61, Gies '87, Stone '91)

# 30 Doradus



ANTON PANNEKOEK  
INSTITUTE

$$Z = Z_{\text{LMC}}$$

# Massive stars have companions



Average number of companions  $f_0 \simeq 2.8$

# Physical processes

- Irradiation
- Mass Transfer
- Tidal effects



## Physical processes

- Irradiation
- Mass Transfer
- Tidal effects

⇒ many more parameters:

$$q \stackrel{\text{def}}{=} \frac{M_2}{M_1}, P \text{ (or } a), e, \dots$$



## Physical processes

- Irradiation
- Mass Transfer
- Tidal effects

⇒ many more parameters:

$$q \stackrel{\text{def}}{=} \frac{M_2}{M_1}, P \text{ (or } a), e, \dots$$



## Astrophysical outcome

- Stripped stars
- Contact binaries
- Runaways & “walkaway” stars
- Mergers



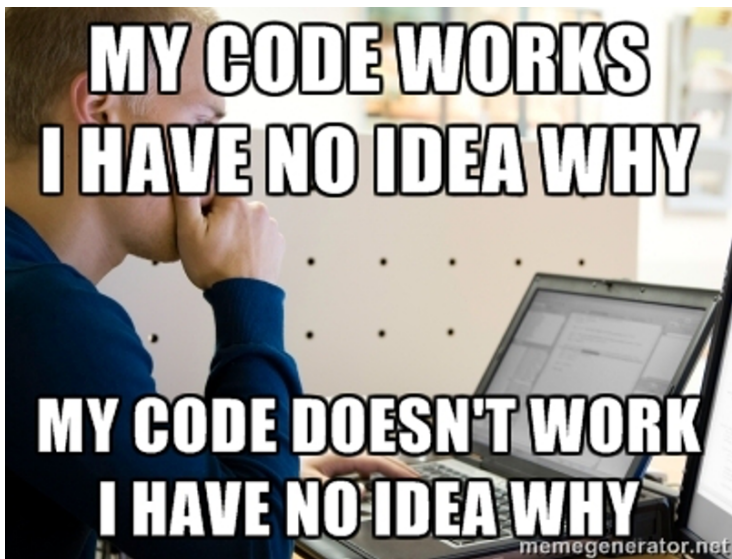
## Introduction: Massive Stars

### Computational astrophysics

- Stellar evolution & structure
- Binary Population Synthesis

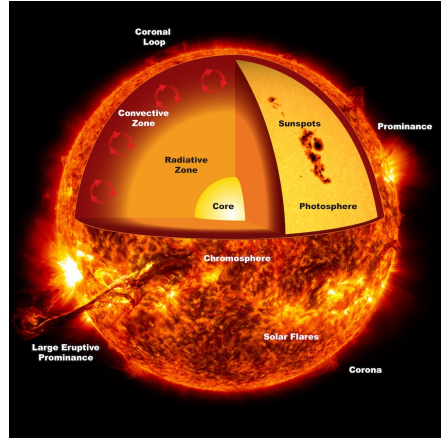
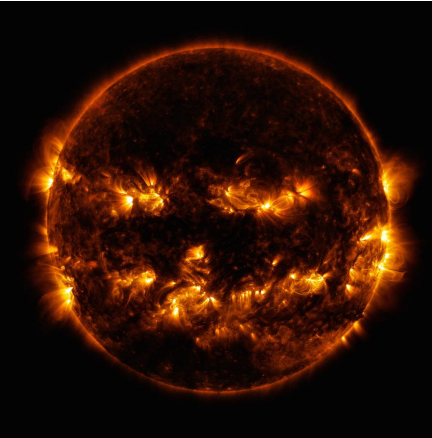
### (If you care) preliminary results

- Can stellar wind change the final fate of a massive star?
- What physics can we learn from breaking apart binaries?



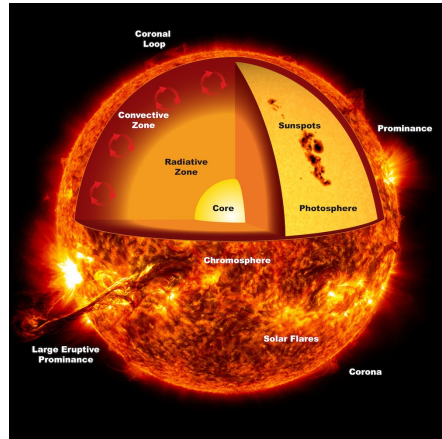
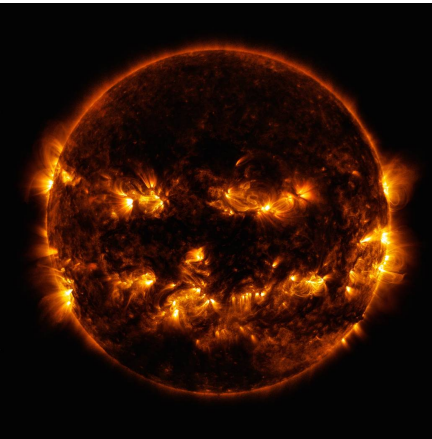
# How can we “look” inside a star?

Figures Credits: NASA



# How can we “look” inside a star?

Figures Credits: NASA



## We simply can't!!

Other Q: How can we observe how *one* star evolves?

- 1 Build a theory from first principles;
- 2 Plug it in a computer;
- 3 Get out a *model*;
- 4 Find a smart way to compare it to what we *can* observe.

## Advantages

- Full control over the parameters  
⇒ Numerical Experiments;
- Allow to focus on interesting things (e.g. no reddening!);
- Allow to deal with long-lasting, rare, inaccessible phenomena;

## Drawbacks

- Numerical errors;
- Limited computational resources;
- **Nature**  $\gg$  **Theory**  $\gg$  **Model**.

*“All models are wrong, but some are useful”* – G. Box

# MESA

is a *tool*, not a theory!

What does it stand for?

**Modules for  
Experiments in  
Stellar  
Astrophysics**

References:

Paxton *et al.* 2011, ApJs192,3  
Paxton *et al.* 2013, ApJs208,4  
Paxton *et al.* 2015, ApJs220,15  
`mesa.sourceforge.net`  
`mesastar.org`

Open Source  $\Leftrightarrow$  Open Know How

*“An algorithm must be seen to be believed”* – D. Knuth

Prohibitive computational cost of 3D  
 $\Rightarrow$  1D, but stars are *not* spherical-symmetric!

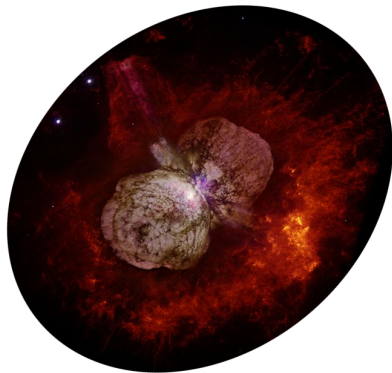
Need of parametric approximations for:

- Rotation  $\Rightarrow$  “Shellular Approximation”;
- Magnetic Fields;
- Convection  $\Rightarrow$  Mixing Length Theory (MLT);
- (Some) mixing processes;
- ...

**Beware of systematic errors!**

$$\frac{dP}{dr} = - \frac{Gm(r)\rho}{r^2}$$

... but stars are not necessarily static!



Other examples:

- He flash,
- Outburst and Eruptions,
- Impulsive mass loss,
- RLOF,
- ...

Figure:  $\eta$  Car, APOD.



Dynamical correction to **static** equilibrium

$$\frac{dP}{dr} = -\frac{Gm(r)\rho}{r^2} - a(r)\rho$$

$$a(r) \stackrel{\text{def}}{=} \frac{dv}{dt} = \frac{d^2r}{dt^2} \ll \frac{Gm(r)}{r^2} \quad \text{“Calculated Passively”}$$

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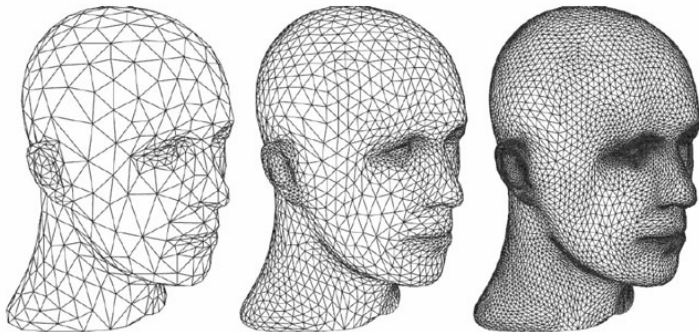
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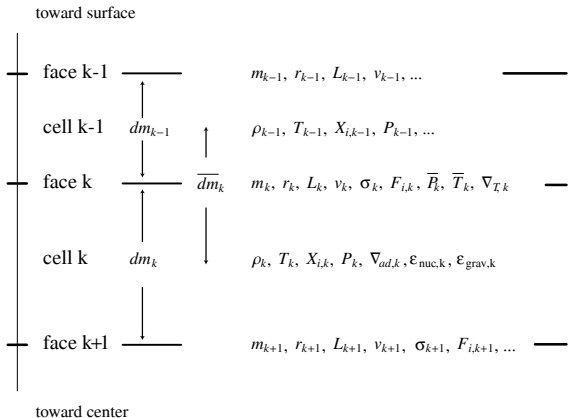
**Explicitly** time-dependent reformulation

$$\frac{\partial v}{\partial t} = -\frac{Gm(r)}{r^2} - \frac{4\pi r^2}{3} \frac{dP}{dm} + g_{\text{visc}}$$

(Euler eq. + reformulation of all stellar structure equations)

$$\frac{df}{dx} \rightarrow \frac{f(x_{k+1}) - f(x_k)}{x_{k+1} - x_k}$$

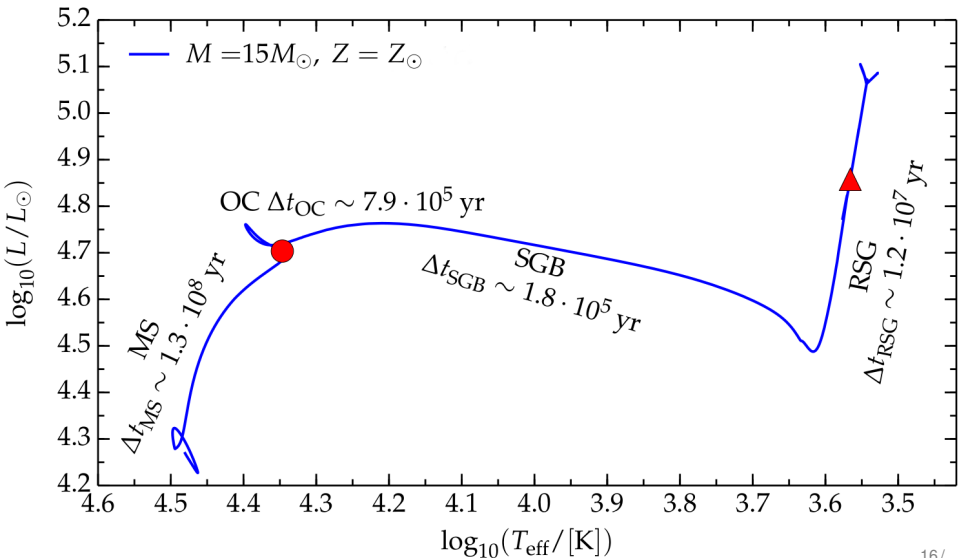




- Intensive quantities (e.g.  $T, \rho$ ) averaged by mass within each cell;
- Extensive quantities (e.g.  $m, L$ ) calculated at outer cell boundary.

Check that physical results do *not* depend on discretization

allow for computation, but resolve physical processes



Physical Theory:

$$\frac{dP}{dr} = -\frac{Gm(r)\rho}{r^2} \quad (+a\rho) \quad \Leftrightarrow$$

$$\frac{dm}{dr} = 4\pi r^2 \rho \quad \Leftrightarrow$$

$$\frac{dT}{dr} = -\frac{3}{16\pi ac} \frac{\kappa\rho L}{r^2 T^3} \quad \Leftrightarrow$$

$$\frac{dL}{dr} = 4\pi r^2 \rho \varepsilon \quad \Leftrightarrow$$

$$P \equiv P(\rho, \mu, T) \quad \Leftrightarrow$$

Numerical Implementation:

$$\left. \frac{dX_i}{dt} \right|_r = \left[ \sum_j \mathcal{P}_{j,i}(T, \rho) - \sum_k \mathcal{D}_{i,k}(T, \rho) \right] + \left[ \sigma_i \nabla^2 X_i \right]$$

$\Updownarrow$

Physical Theory:

Numerical Implementation:

$$\frac{dP}{dr} = -\frac{Gm(r)\rho}{r^2} \quad (+a\rho)$$

 $\Leftrightarrow$ 

$$\frac{P_{k-1} - P_k}{0.5(dm_{k-1} - dm_k)} = -\frac{Gm_k}{4\pi r_k^4} - \frac{a_k}{4\pi r_k^2}$$

$$\frac{dm}{dr} = 4\pi r^2 \rho$$

 $\Leftrightarrow$ 

$$\ln(r_k) = \frac{1}{3} \ln \left[ r_{k+1}^3 + \frac{3}{4\pi} \frac{dm_k}{\rho_k} \right]$$

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$$\frac{T_{k-1} - T_k}{(dm_{k-1} - dm_k)/2} = -\nabla_{T,k} \left( \frac{dP}{dm} \Big|_k \right) \frac{\tilde{T}_k}{\tilde{P}_k}$$

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## Lagrangian



$$\mathbf{v} \equiv \mathbf{v}(m, t)$$

## Eulerian



$$\mathbf{v} \equiv \mathbf{v}(\mathbf{r}, t)$$

Physical Theory:

Numerical Implementation:

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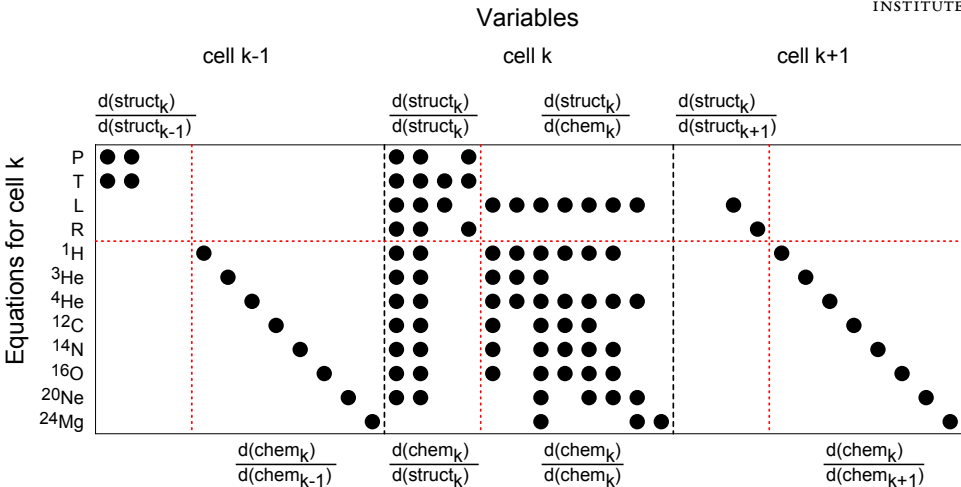


Figure: From Paxton *et al.* 2013, *ApJs*, 208, 4. Black dots are non-zero entries.

- **Henvey** code: varies all the quantities in each zone until an acceptable solution is found ( $\neq$  Shooting Method);
- Generalized **Newton-Raphson** solver ( $\Rightarrow$  FIRST ORDER):

$$0 = \mathbb{F}(y) \simeq \mathbb{F}(y_i + \delta y_i) = \mathbb{F}(y_i) + \left[ \frac{d\mathbb{F}(y)}{dy} \right]_i \delta y_i + O((\delta y_i)^2) ;$$

$$\delta y_i \simeq - \frac{\mathbb{F}(y_i)}{\left[ \frac{d\mathbb{F}(y)}{dy} \right]_i}$$

$$\Downarrow$$

$$y_{i+1} = y_i + \delta y_i$$

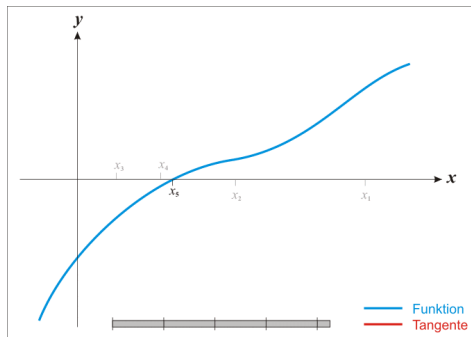


Figure: From Wikipedia

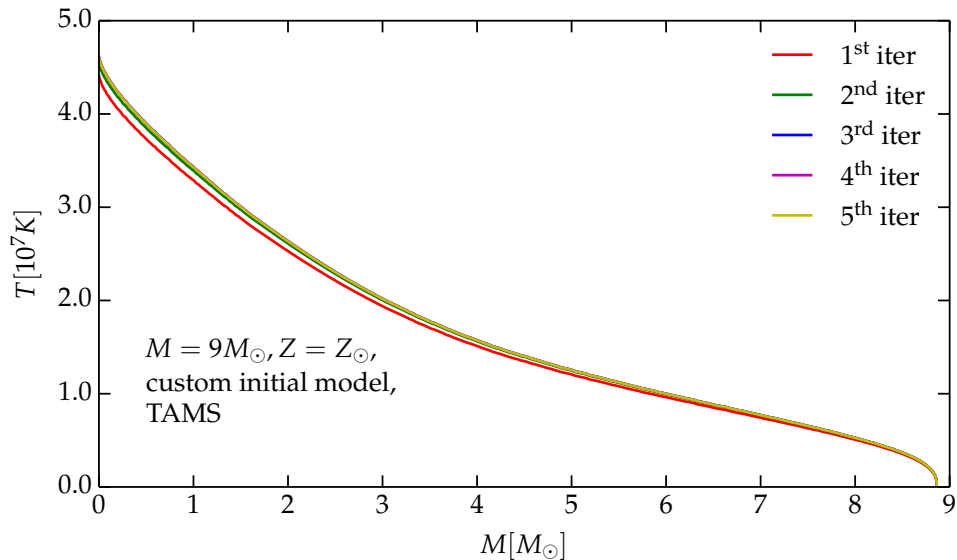


Figure: Two models after the end of core hydrogen burning

## Introduction: Massive Stars

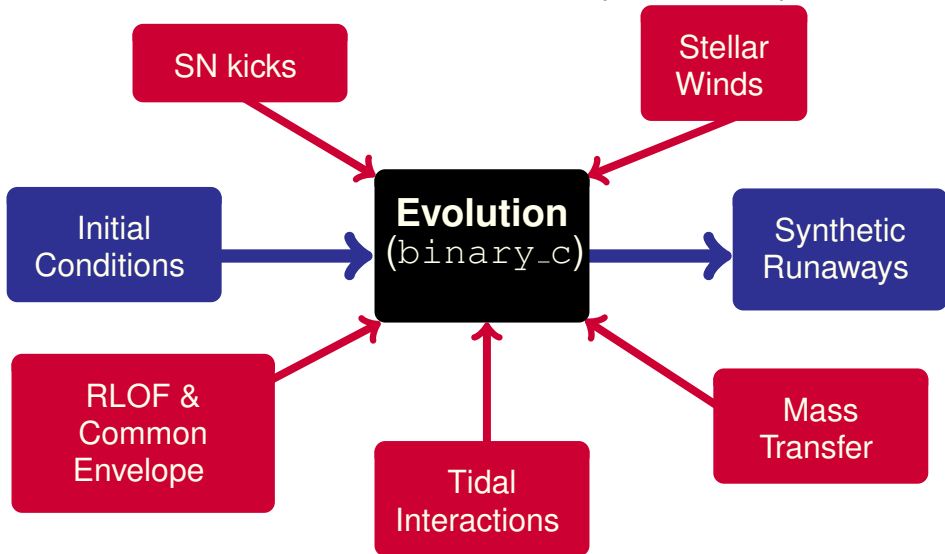
### Computational astrophysics

- Stellar evolution & structure
- Binary Population Synthesis

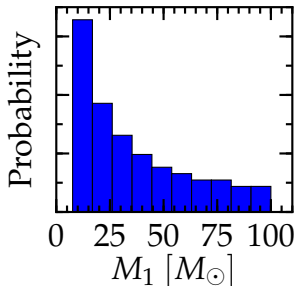
### (If you care) preliminary results

- Can stellar wind change the final fate of a massive star?
- What physics can we learn from breaking apart binaries?

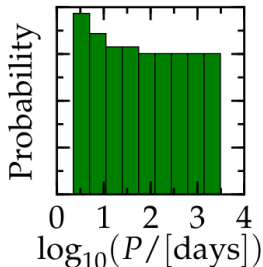
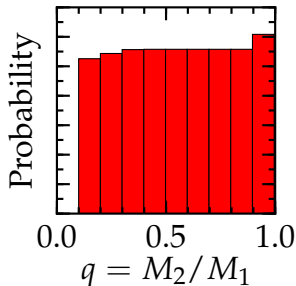
Fast  $\Rightarrow$  Allows statistical tests of the inputs & assumptions





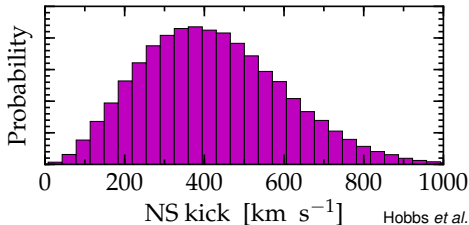


Kroupa '01

Sana *et al.* '12

Total Population:  $2 \times 10^6$  stars

Maxwellian  $\sigma_{v_{\text{kick}}} = 265 [\text{km s}^{-1}]$

Hobbs *et al.* '05

## Introduction: Massive Stars

### Computational astrophysics

- Stellar evolution & structure
- Binary Population Synthesis

### **(If you care) preliminary results**

- Can stellar wind change the final fate of a massive star?
- What physics can we learn from breaking apart binaries?

# Why are Massive Stars Important?

Nucleosynthesis & Chemical Evolution

Star Formation

Ionizing Radiation

Supernovae  
(if  $M_{ZAMS} \gtrsim 8 M_{\odot}$ )

GW Astronomy

## Mass loss for the environment:

- Pollution of ISM
- Tailoring of CSM
- Trigger for Star Formation

## Mass loss for the star

- Evolutionary Timescales
- Appearance & Classification (e.g. WR)
- Light Curve and Explosion Spectrum
- Final Fate: BH, NS or WD?

Radiative Driving



Stellar Winds

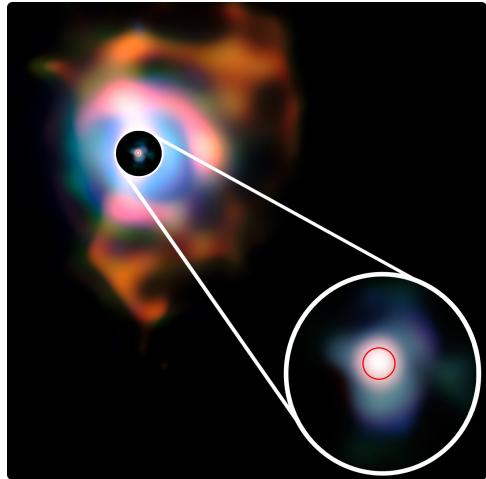


Figure: Betelgeuse

## Dynamical Instabilities



LBVs, Impulsive Mass Loss,  
Pulsations,  
Super-Eddington Winds

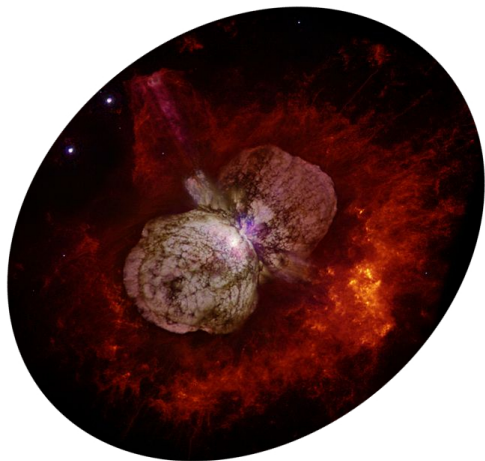


Figure:  $\eta$  Carinae.

## Binary interactions



Roche Lobe Overflow, Common  
Envelope, Fast rotation

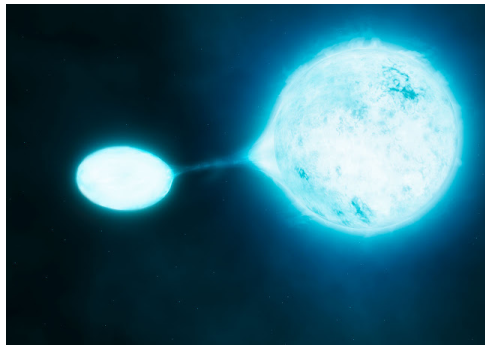
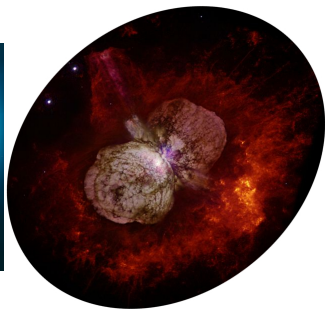


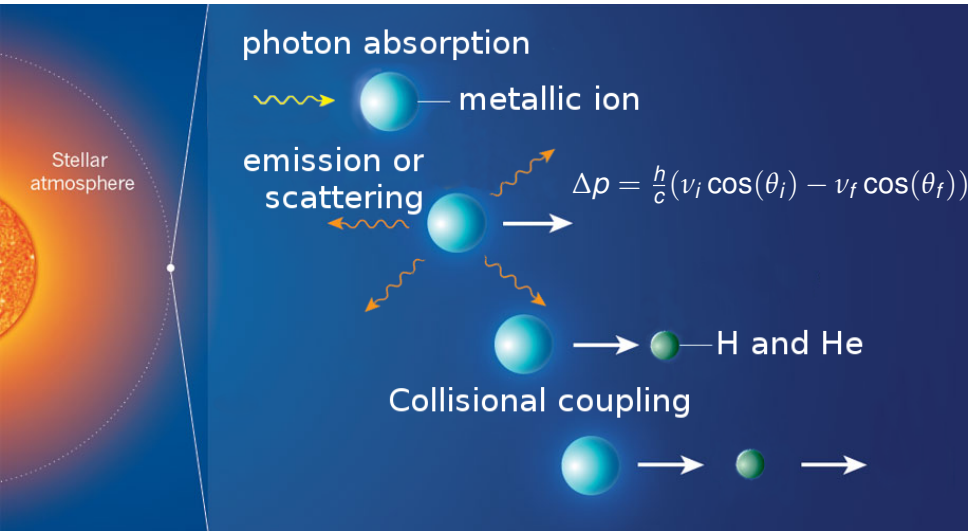
Figure: Artist Impression



... but stellar evolution codes assume hydrostatic equilibrium:

$$\frac{dP}{dr} = - \frac{Gm(r)\rho}{r^2}$$

Open question: **Which dominates in term of total mass lost?**



Problems: High Non-Linearity and Clumpiness



## Inhomogeneities:

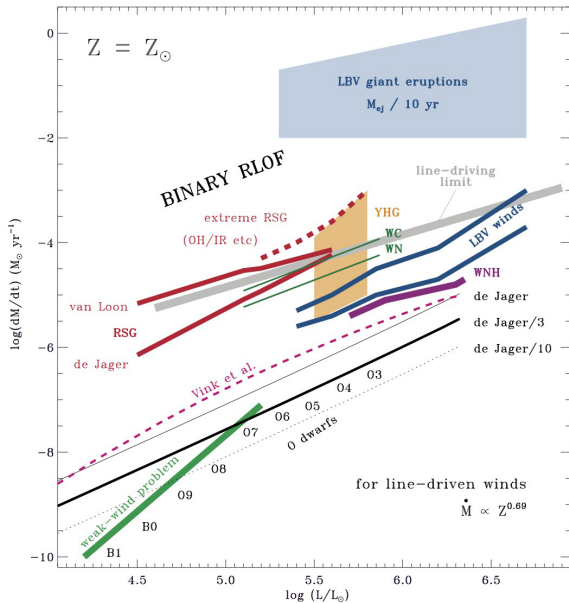
$$f_{\text{cl}} \stackrel{\text{def}}{=} \frac{\langle \rho^2 \rangle}{\langle \rho \rangle^2} \neq 1 \Rightarrow \dot{M} \neq 4\pi r^2 \rho v(r)$$

## Inhomogeneities:

$$f_{\text{cl}} \stackrel{\text{def}}{=} \frac{\langle \rho^2 \rangle}{\langle \rho \rangle^2} \neq 1 \Rightarrow \dot{M} \neq 4\pi r^2 \rho v(r)$$

## Risk:

Possible overestimation of the  
wind mass loss rate



(Semi-)Empirical  
parametric models.

Efficiency factor:

$$\dot{M}(L, T_{\text{eff}}, Z, R, M, \dots)$$

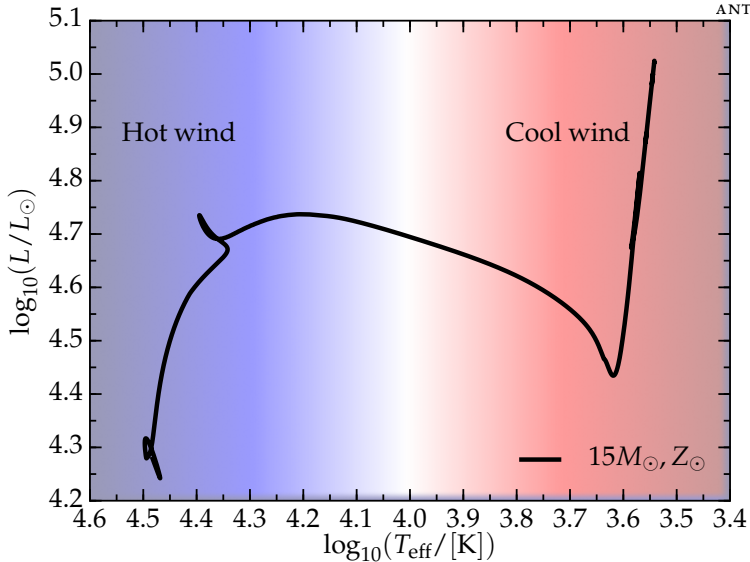


$$\eta \dot{M}(L, T_{\text{eff}}, Z, R, M, \dots)$$

$\eta$  is a **free** parameter:

$$\eta \in [0, +\infty)$$

Figure: From Smith 2014, ARA&A, 52, 487S



WR wind  $\Leftrightarrow X_s < 0.4$

Grid of  $Z_{\odot} \simeq 0.019$ , non-rotating stellar models:

- Initial mass:

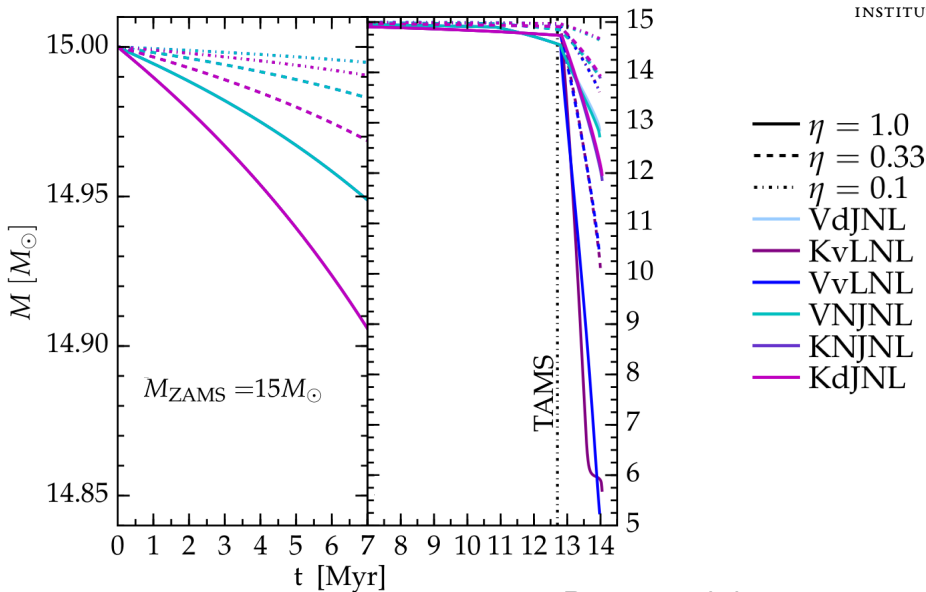
$$M_{\text{ZAMS}} = \{15, 20, 25, 30, 35\} M_{\odot};$$

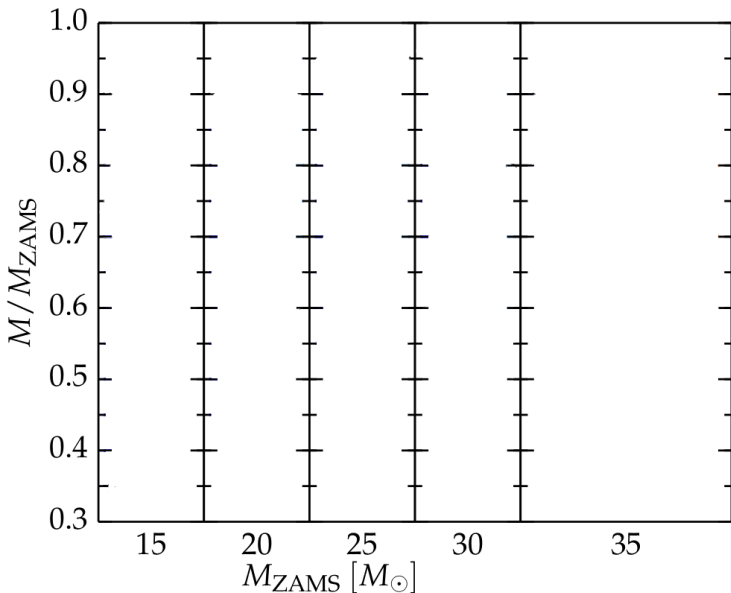
- Efficiency:

$$\eta = \left\{1, \frac{1}{3}, \frac{1}{10}\right\};$$

- Combinations of wind mass loss rates for “hot” ( $T_{\text{eff}} \geq 15$  [kK]), “cool” ( $T_{\text{eff}} < 15$  [kK]) and WR:

Kudritzki *et al.* '89; Vink *et al.* '00, '01;  
 Van Loon *et al.* '05; Nieuwenhuijzen *et al.* '90;  
 De Jager *et al.* '88;  
 Nugis & Lamers '00; Hamann *et al.* '98.

Renzo *et al.*, in prep.

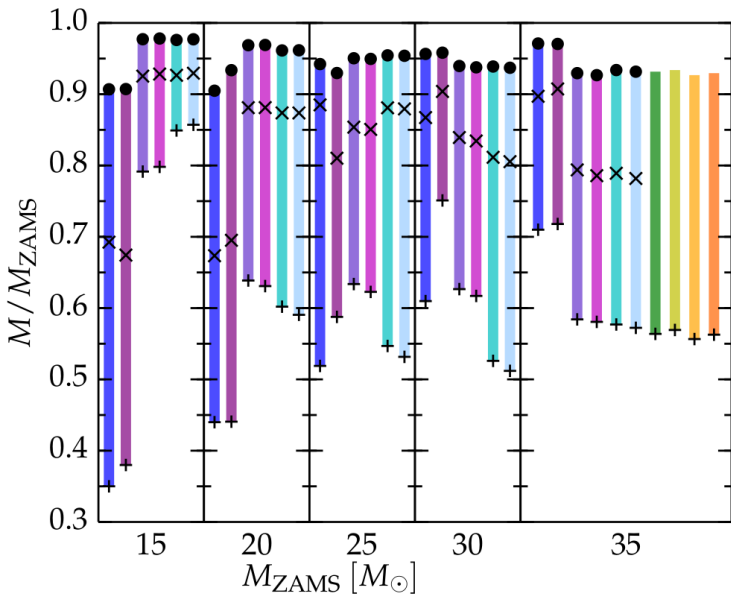


MESA

Legend:

- $\eta = 0.1$
- ×  $\eta = 0.33$
- +  $\eta = 1.0$

Renzo *et al.*, in prep.



MESA

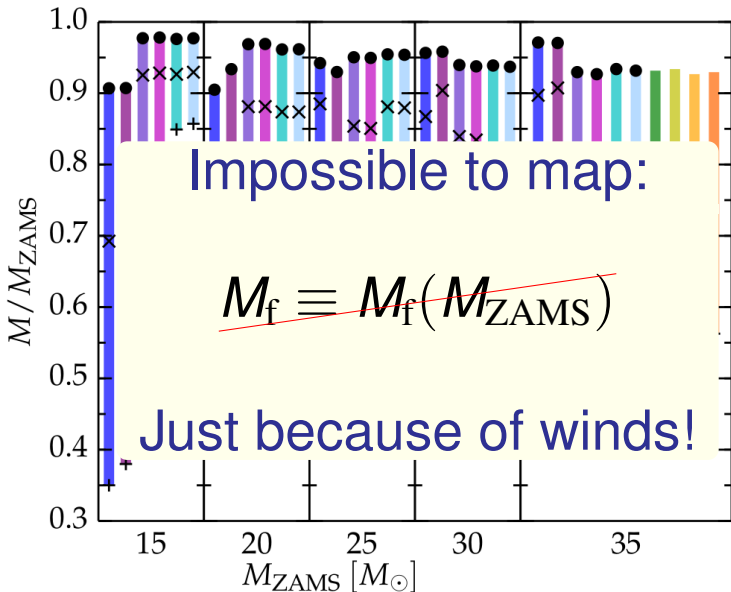
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$\eta \rightarrow$  largest  
uncertainty

Renzo *et al.*, in prep.





MESA

**Legend:**

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- ✕  $\eta = 0.33$
- +

$\eta \rightarrow$  largest  
uncertainty

Renzo *et al.*, in prep.

# “Explodability” & Compactness



$$\tilde{\zeta}_{\mathcal{M}}(t) \stackrel{\text{def}}{=} \frac{\mathcal{M}/M_{\odot}}{R(\mathcal{M})/1000 \text{ km}}$$

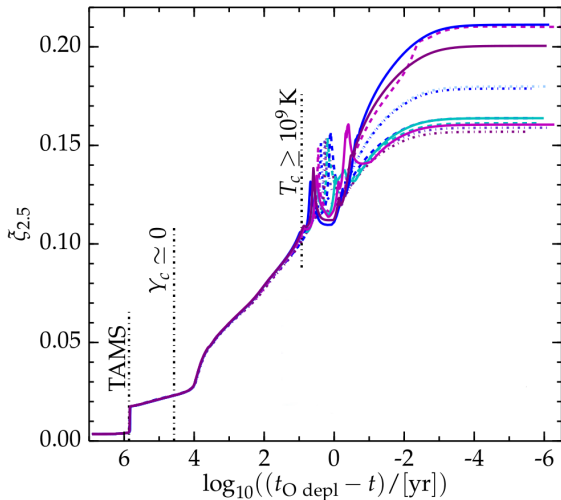
- “Large”  $\tilde{\zeta}_{2.5} \Rightarrow$  harder to explode  $\Rightarrow$  BH formation
- “Small”  $\tilde{\zeta}_{2.5} \Rightarrow$  easier to explode  $\Rightarrow$  NS formation

(e.g. O’Connor & Ott 2011,  
Ugliano *et al.* 2012,  
Sukhbold & Woosley 2014)

$\mathcal{M} = 2.5 M_{\odot}$

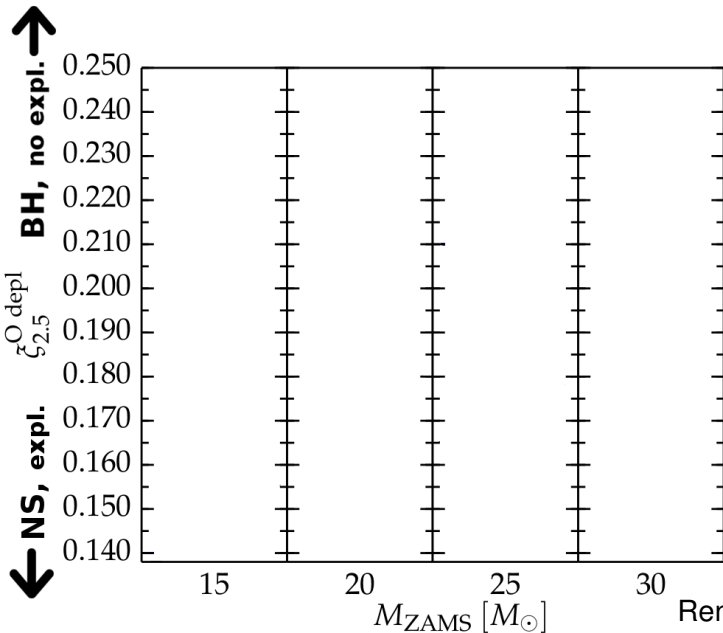
not to scale!

$R(\mathcal{M})$

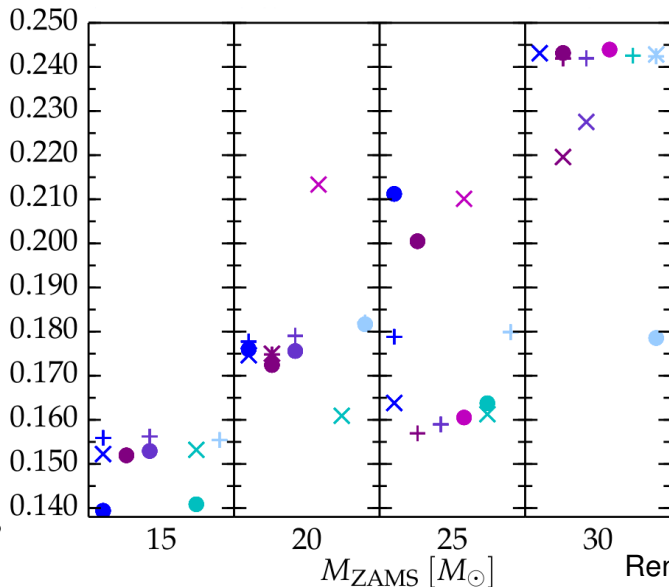
$M_{\text{ZAMS}} = 25 M_{\odot}$  MESA models


Critical point: Ne core burning/C shell burning

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**BH, no expl.**  $\zeta_{2.5}^{\text{depl}}$   
**NS, expl.**



**Legend:**

- $\eta = 0.1$
- ✕  $\eta = 0.33$
- +  $\eta = 1.0$

Post O burning  
evolution



Core contraction



**Amplification of  
the differences.**

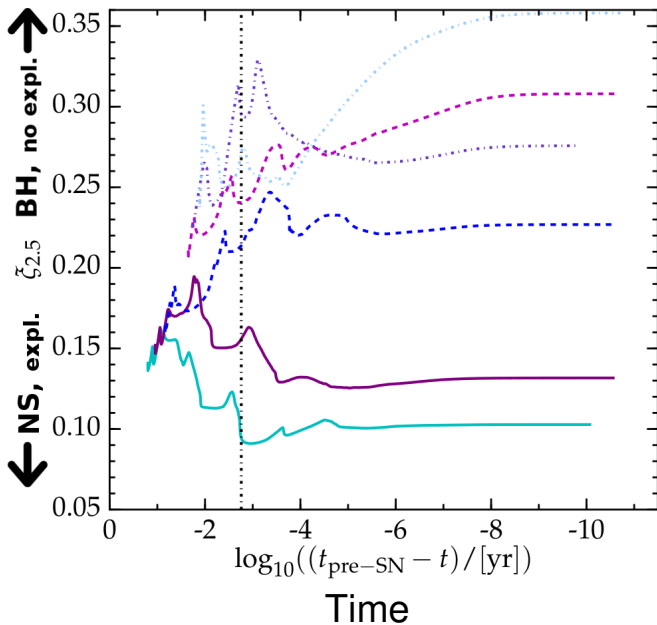
Renzo *et al.*, in prep.

- Initially small effect  $\Rightarrow N_{\text{zones}} \gtrsim 20\,000$ ;
- Complex nuclear burning  $\Rightarrow N_{\text{iso}} \gtrsim 200$ ;



SurfSara's **Cartesius** Computer.

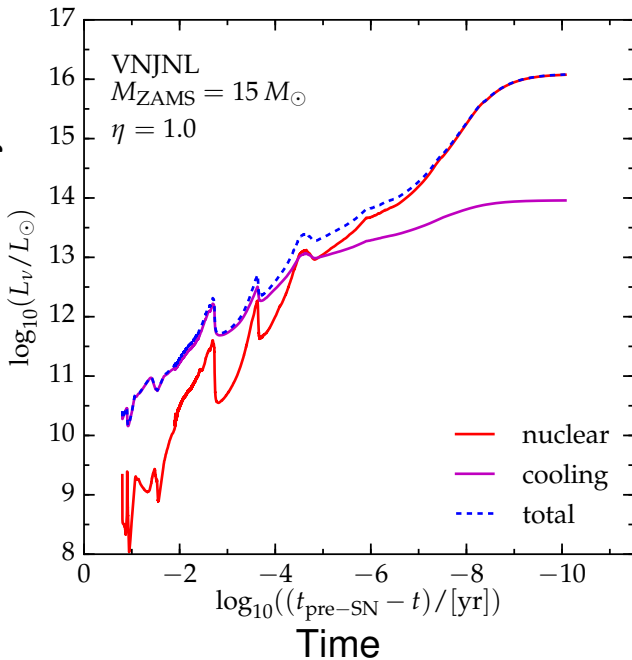
Si shell burning →



- $15M_{\odot}, \eta = 1.0$
- - -  $25M_{\odot}, \eta = 0.33$
- ⋯  $30M_{\odot}, \eta = 0.33$

Renzo *et al.*, in prep.

Neutrino Luminosity



Fuel ignition in  
 (partially)  
 degenerate  
 environment



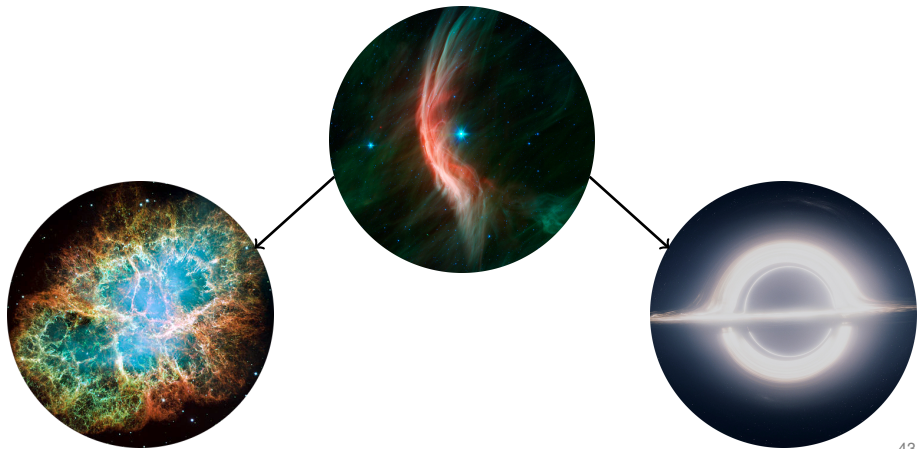
Flash

Renzo *et al.*, in prep.



## Uncertainties in stellar winds:

- pre-SN mass  $\Rightarrow$  no  $M_f \equiv M_f(M_{ZAMS})$  map;
- core structure  $\Rightarrow$  “explodability” & remnant.



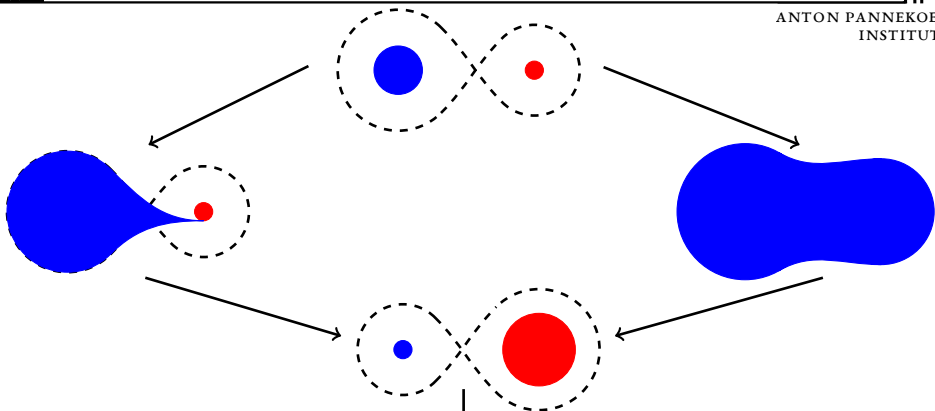
## Introduction: Massive Stars

### Computational astrophysics

- Stellar evolution & structure
- Binary Population Synthesis

### **(If you care) preliminary results**

- Can stellar wind change the final fate of a massive star?
- What physics can we learn from breaking apart binaries?



- Unbinding Matter

(e.g. Blaauw '61)

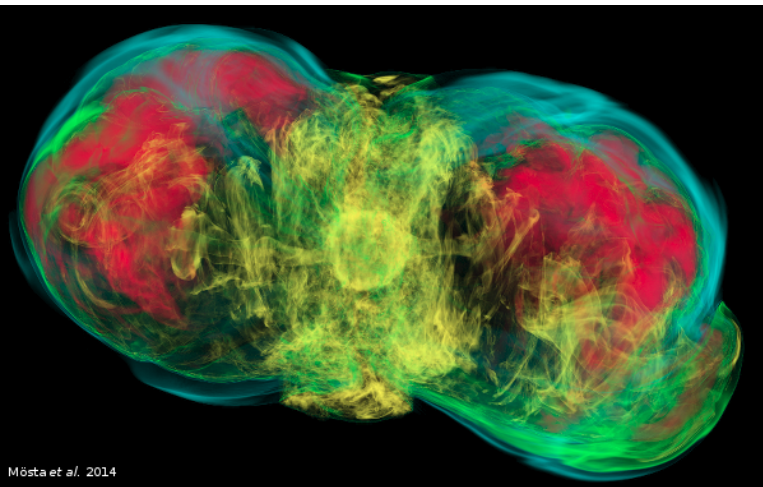
- Ejecta Impact

(e.g. Tauris & Taken '98)

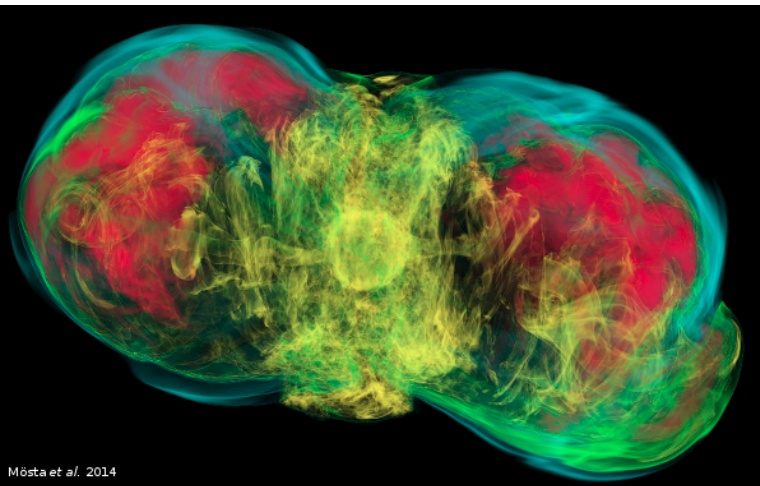
- SN Natal Kick

(e.g. Cordes *et al.* '93)

$$v_{RW} \approx v_2^{\text{orb}}$$



$\nu$  emission and/or ejecta anisotropies

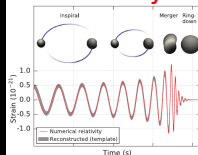


Mösta et al. 2014

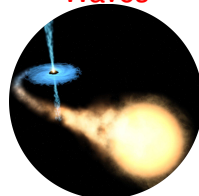
$\nu$  emission and/or ejecta anisotropies



Runaways



Gravitational Waves



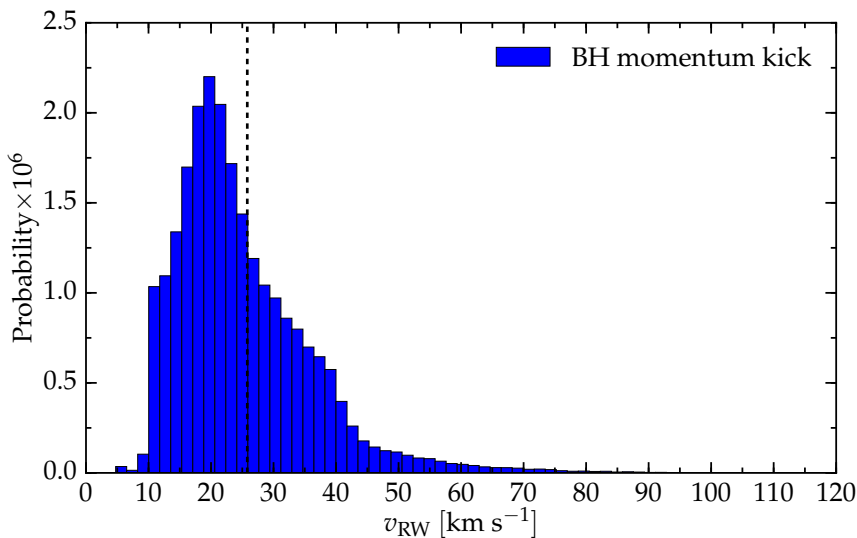
XRBs 46/

## 30 Doradus

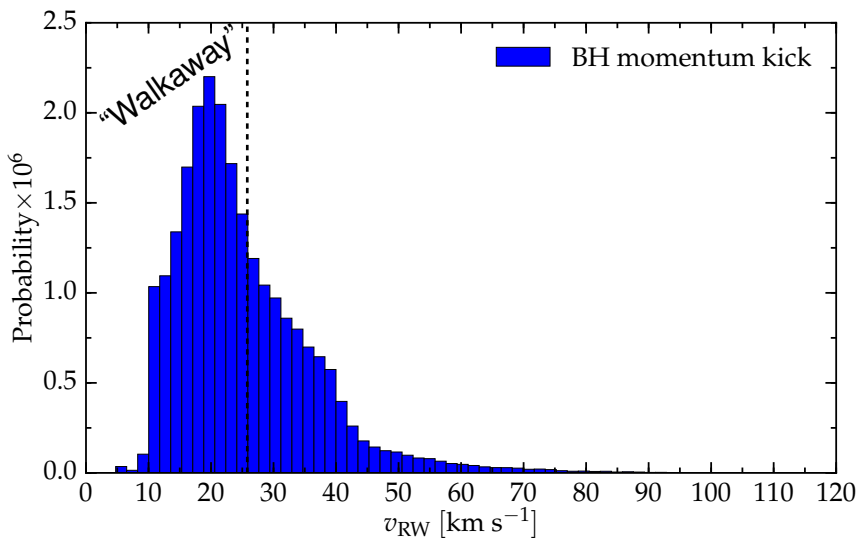
ANTON PANNEKOEK  
INSTITUTE

$$Z = Z_{\text{LMC}}$$

## O-type from disrupted binaries only

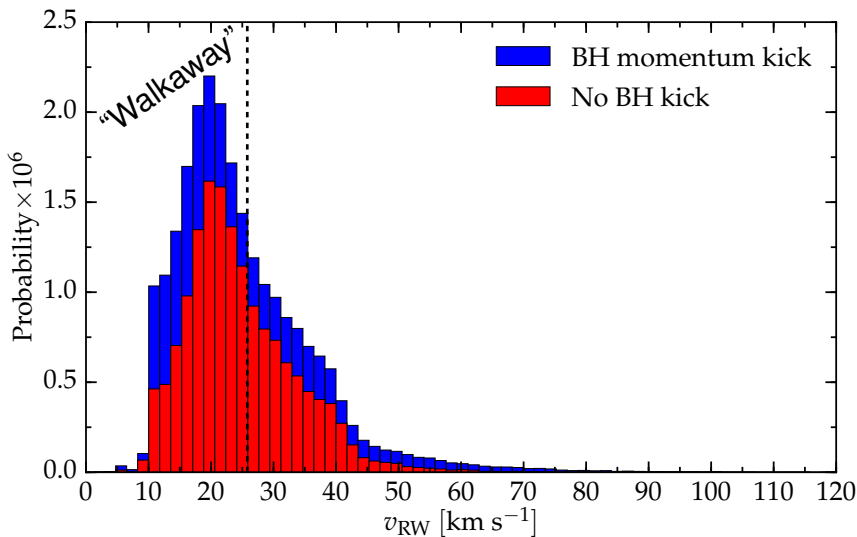


## O-type from disrupted binaries only

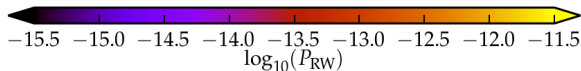
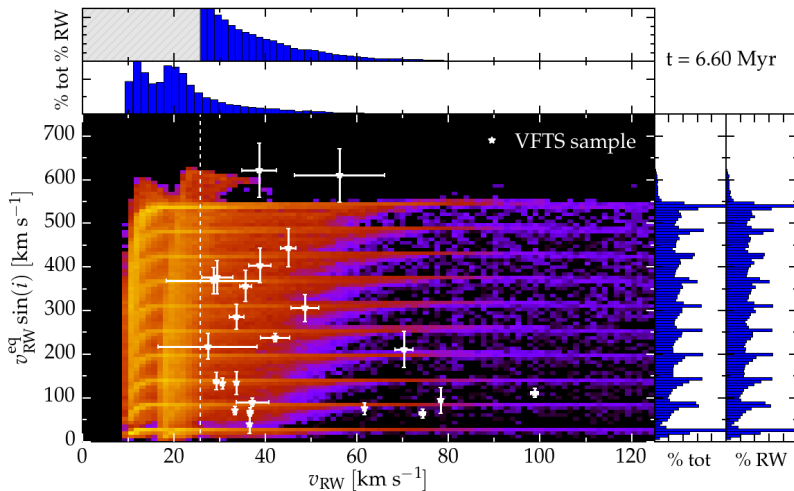




## O-type from disrupted binaries only



~ Rotational velocity



~ Line of sight velocity

