

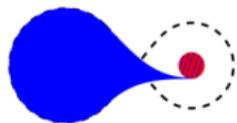
# Mass transfer in binary systems

“Widowed” stars and “living” gravitational wave sources

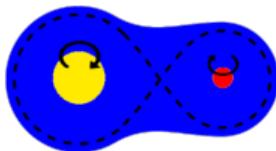


## Mass transfer in binary systems can be

**dynamically stable:**

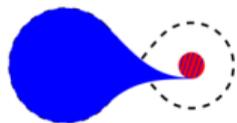


**dynamically unstable:**



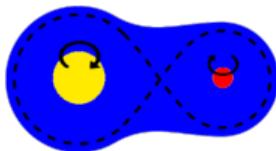
### **dynamically stable:**

“Widowed” accretor stars



### **dynamically unstable:**

can LISA detect GW from common envelope evolution?



# Why care about the accretor?

## Stellar populations



accretors hide in samples

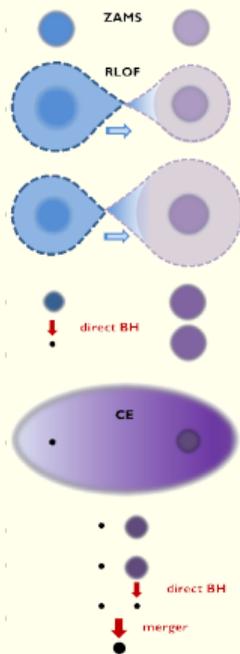
de Mink *et al.* 2013, Renzo *et al.* 2019b

+

Oe/Be stars, stragglers

Pols *et al.* 1991, Wang *et al.* 2021

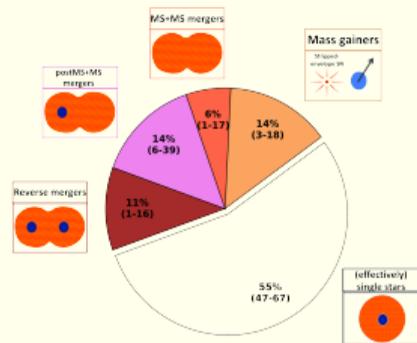
## GW progenitors



e.g., Belczynski *et al.* 2016

## Transients

type II supernovae



Zapartas *et al.* (incl. MR) 2019

+

long GRB

Cantiello *et al.* 2007, MR & Götzberg 2021

# Most common massive binary evolution path: stable case B RLOF

---

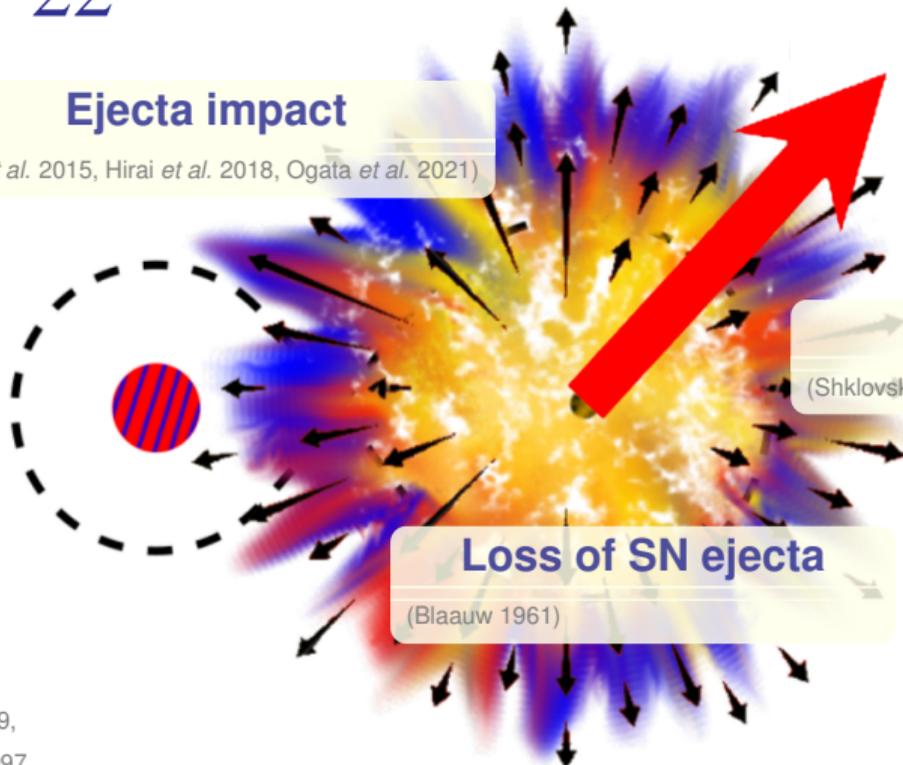
Credits: ESO, L. Calçada, M. Kornmesser, S.E. de Mink

# SN natal kicks disrupt the binary

$86^{+11}_{-22}\%$  of massive binaries are disrupted

## Ejecta impact

(Liu *et al.* 2015, Hirai *et al.* 2018, Ogata *et al.* 2021)



## SN natal kick

(Shklovskii 1970, Katz 1975, Janka 2013, 2017)

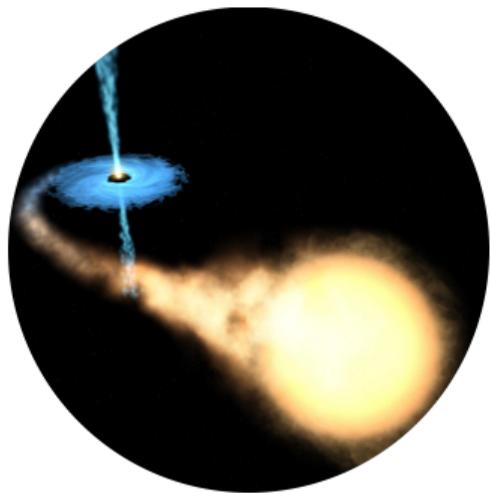
## Loss of SN ejecta

(Blaauw 1961)

## Do BHs receive kicks ?

NO

⇒ most remain bound to companion



YES

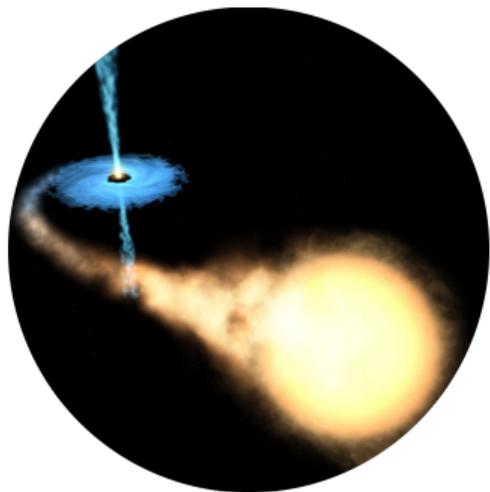
⇒ most are single and we can't see them...



## Do BHs receive kicks ?

NO

⇒ most remain bound to companion



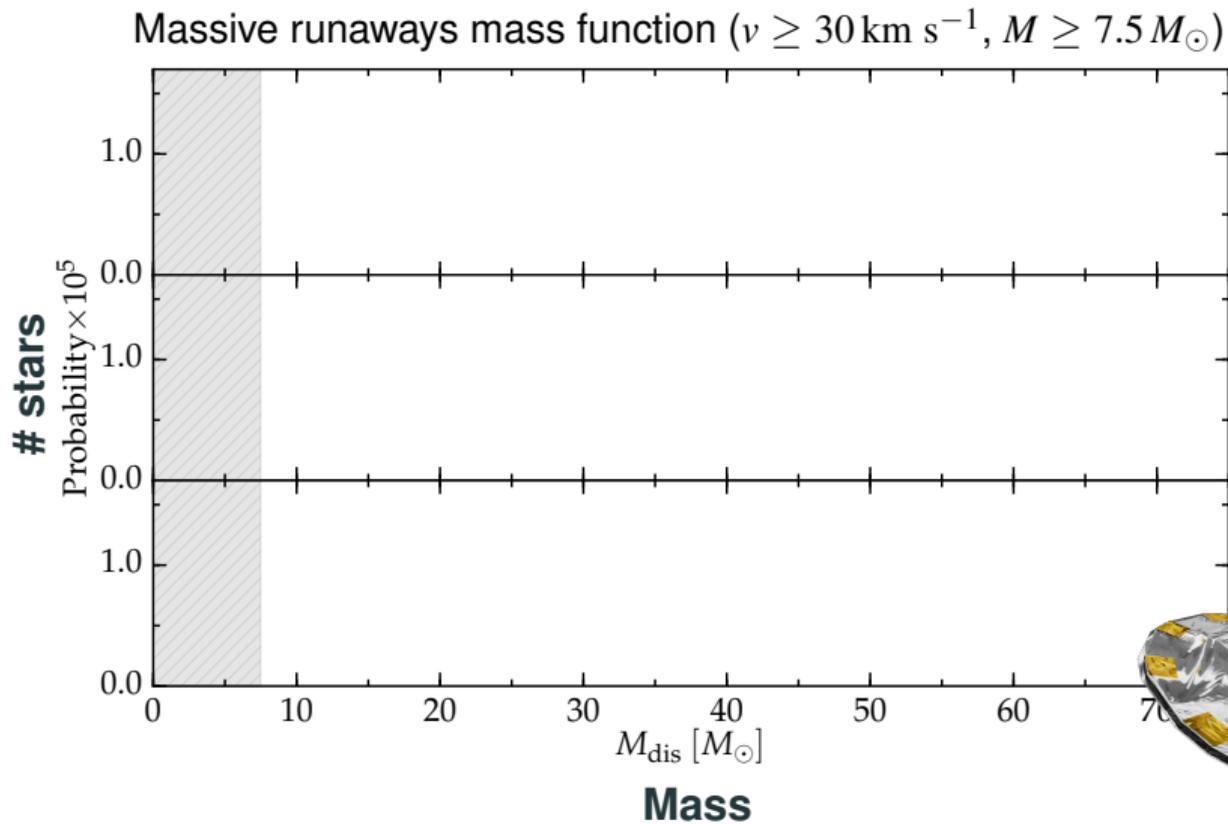
YES

⇒ most are single and we can't see them...



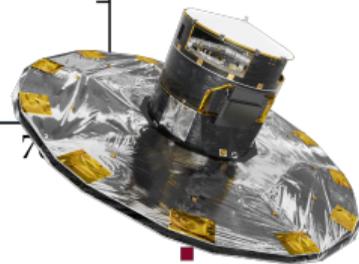
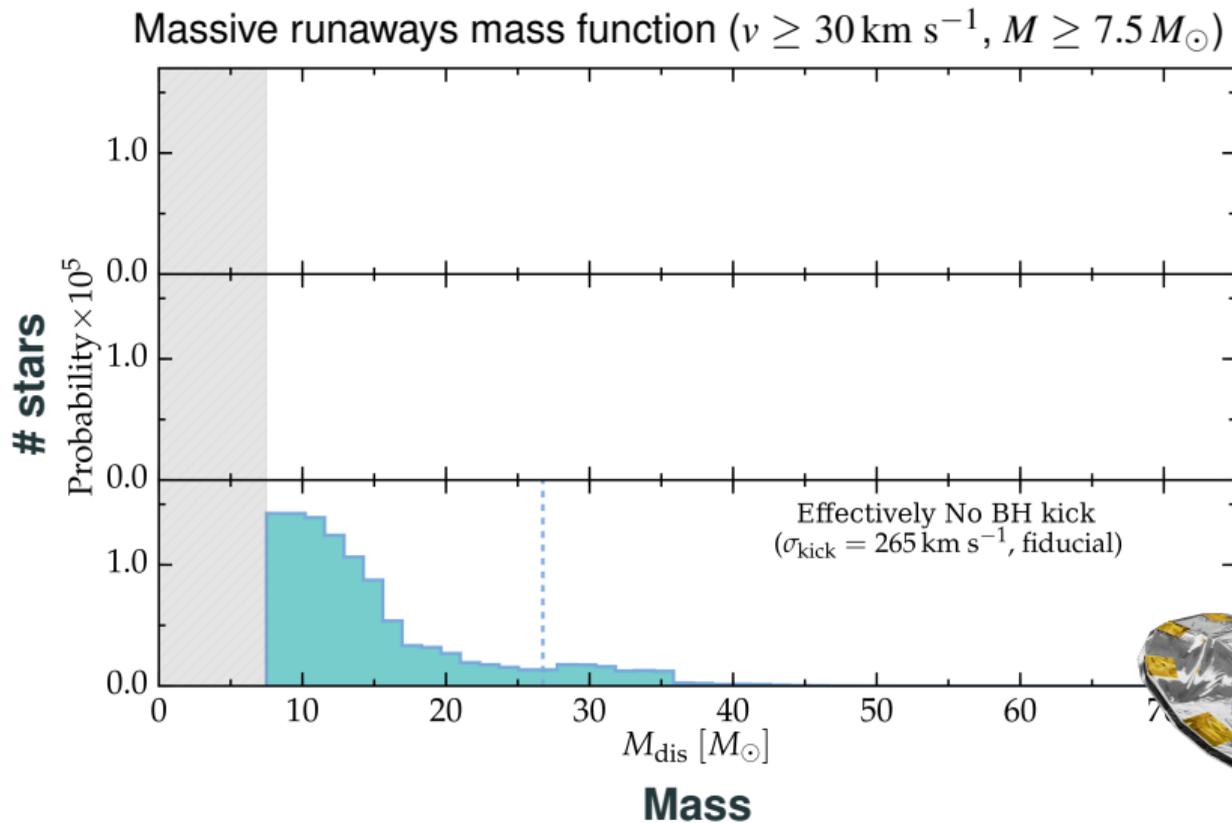
...but we can see the  
“widowed” companions

# Constrain BH kicks looking at the former companion with *Gaia*



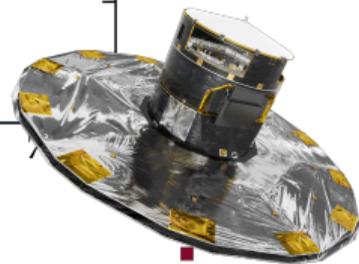
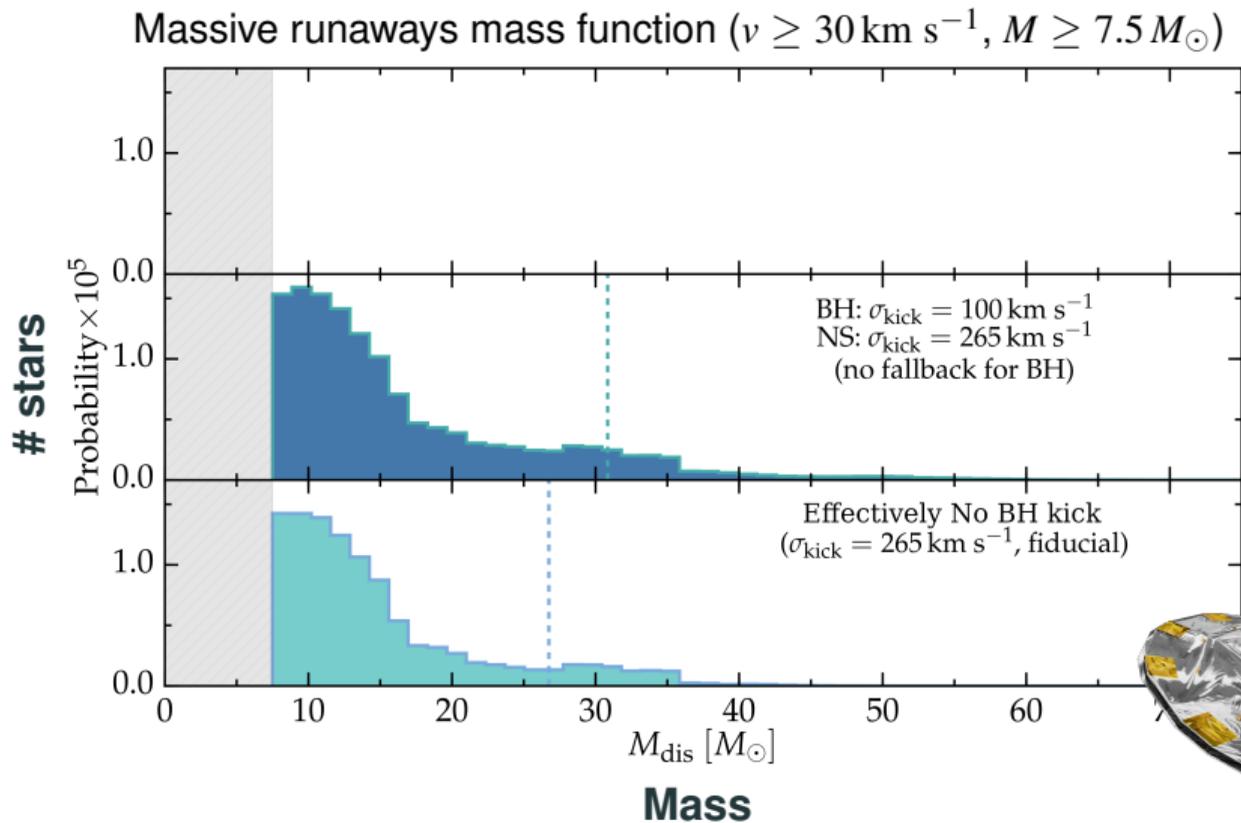
gaia

# Constrain BH kicks looking at the former companion with *Gaia*



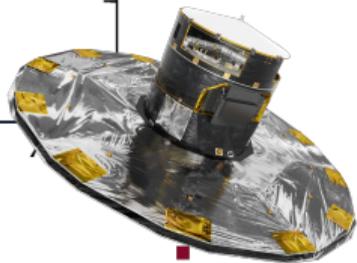
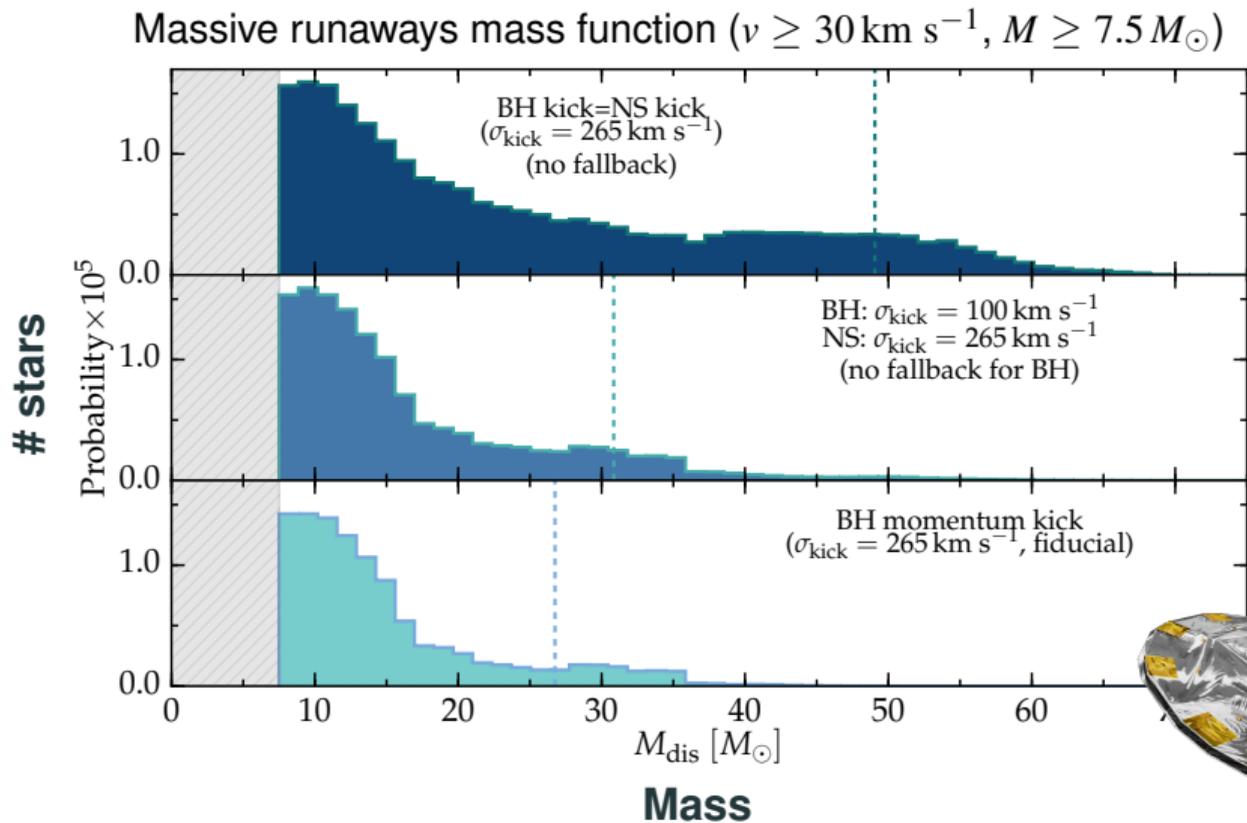
gaia

# Constrain BH kicks looking at the former companion with *Gaia*



gaia

# Constrain BH kicks looking at the former companion with *Gaia*



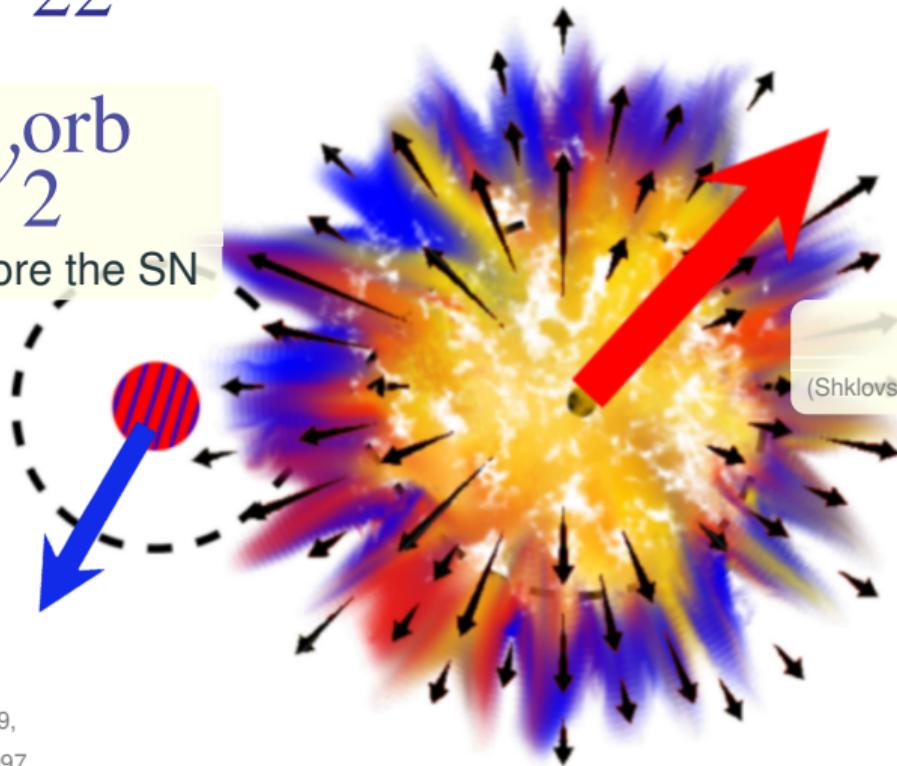
gaia

## Kicks do not change the velocity of the widowed star

$86^{+11}_{-22}\%$  of massive binaries are disrupted

$$v_{\text{dis}} \simeq v_{\text{orb}}^{\text{orb}}$$

before the SN



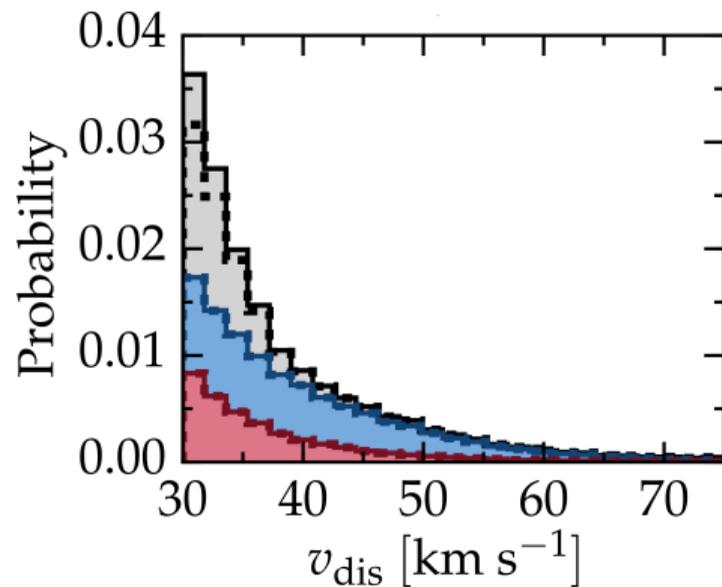
**SN Natal kick**

(Shklovskii 1970, Katz 1975, Janka 2013, 2017)

## **Kinematics of the widowed stars**

---

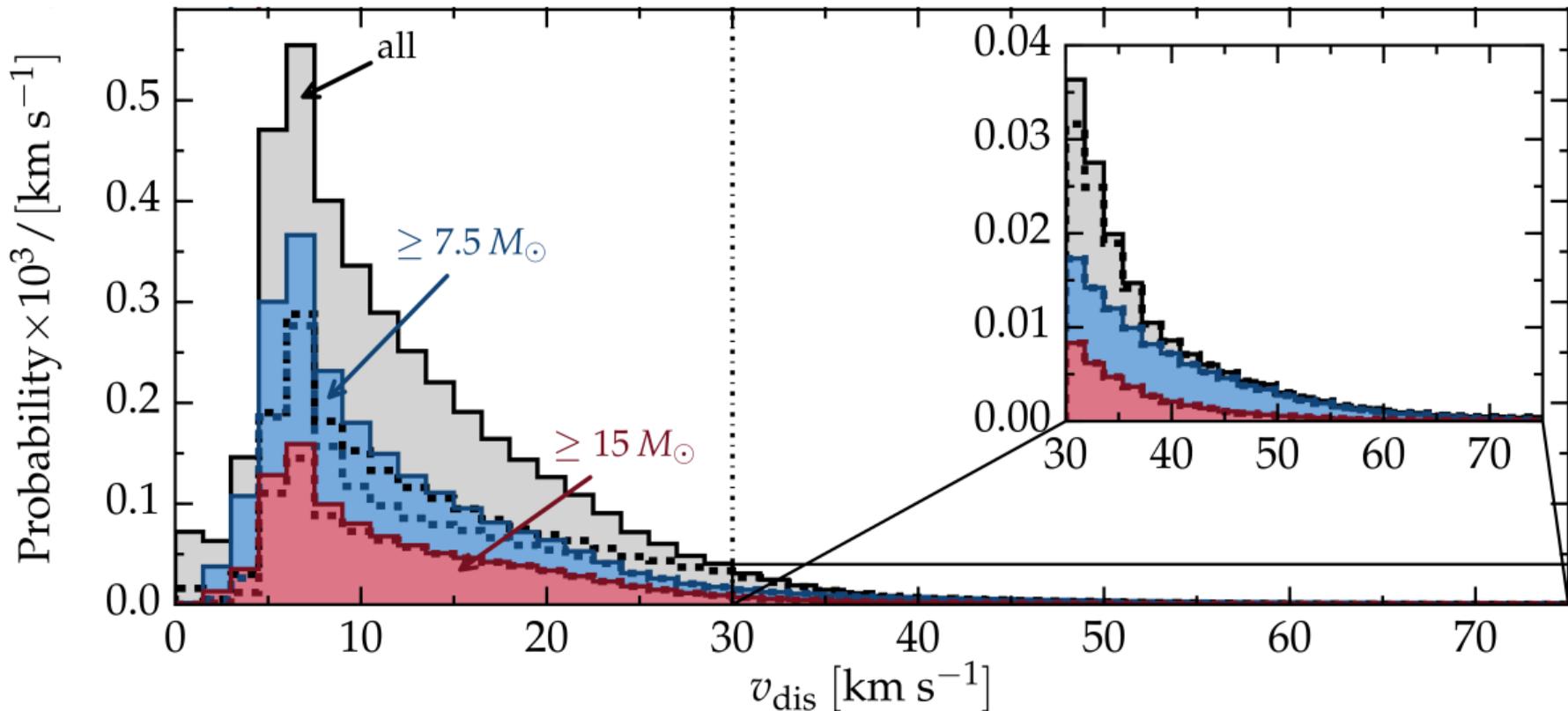
## Widowed stars can be *runaways*...



Velocity w.r.t. pre-explosion binary center of mass

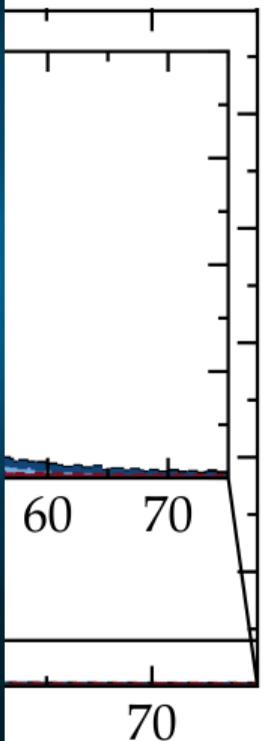
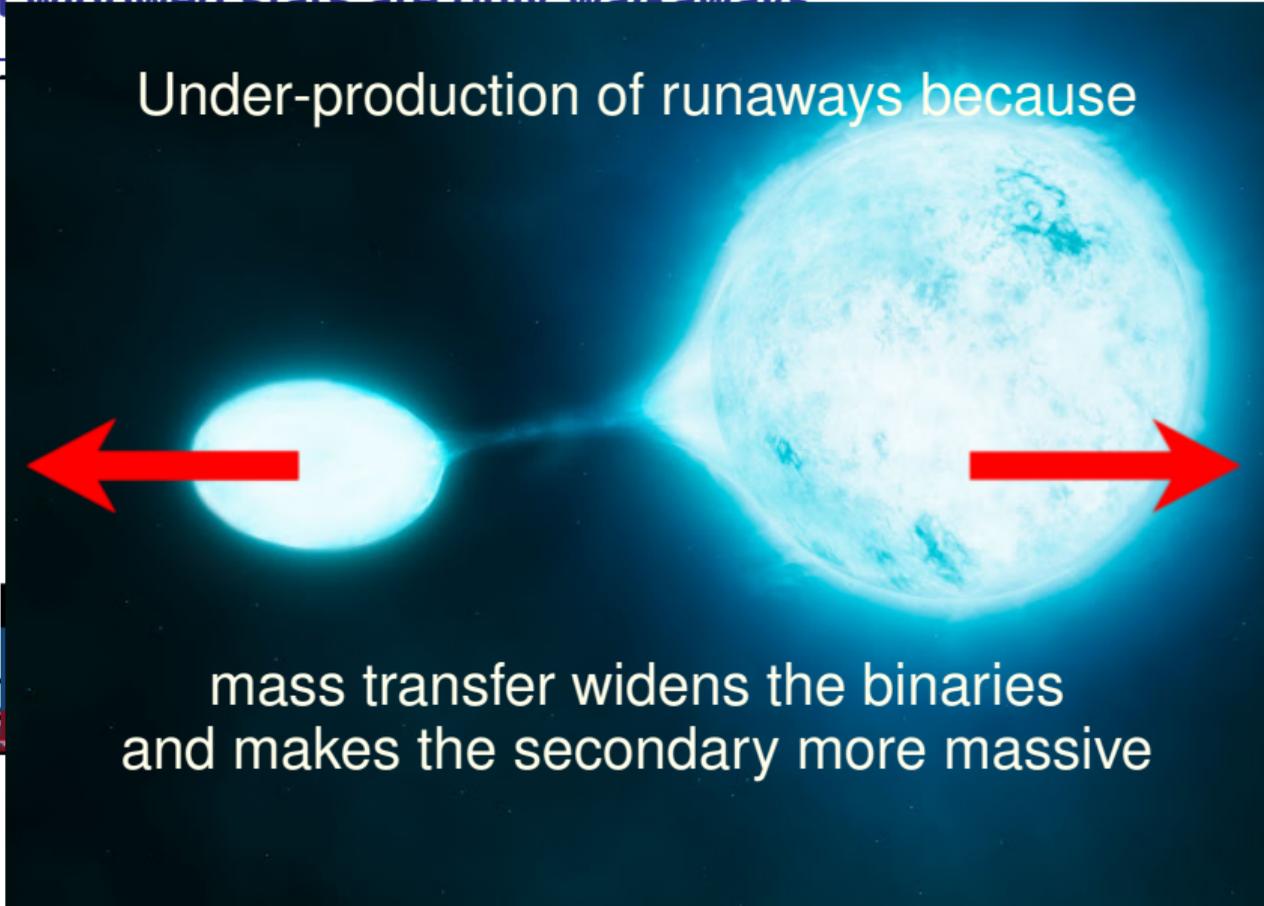
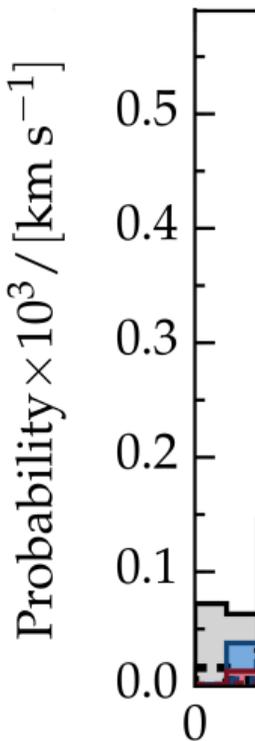
Numerical results: <http://cdsarc.u-strasbg.fr/viz-bin/qcat?J/A+A/624/A66>

## ...but most widowed stars are only *walkaways*



Velocity respect to the pre-explosion binary center of mass

...but most widowed stars are only walkaways



velocity respect to the pre-explosion binary center of mass

## **Appearance of “widowed” stars**

---

# Spin up, pollution, and rejuvenation of the second star

The binary disruption shoots out  
the accretor



Spin up: Packet 1981, Cantiello *et al.* 2007, de Mink *et al.* 2013

Pollution: Blaauw 1993

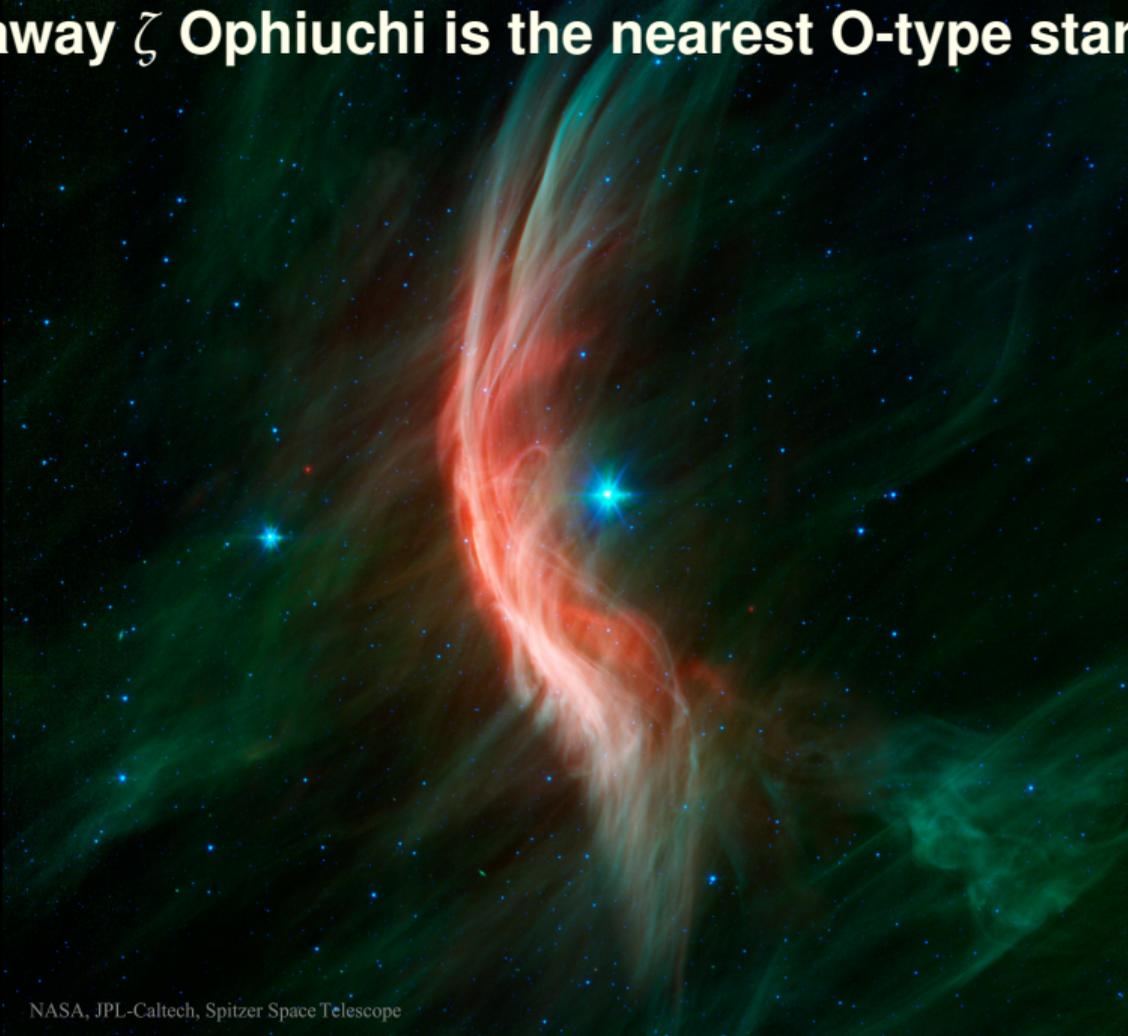
Rejuvenation: Hellings 1983, Schneider *et al.* 2015

## **Appearance of “widowed” stars**

---

**Constraints from the nearest O-type star**

# The runaway $\zeta$ Ophiuchi is the nearest O-type star to Earth



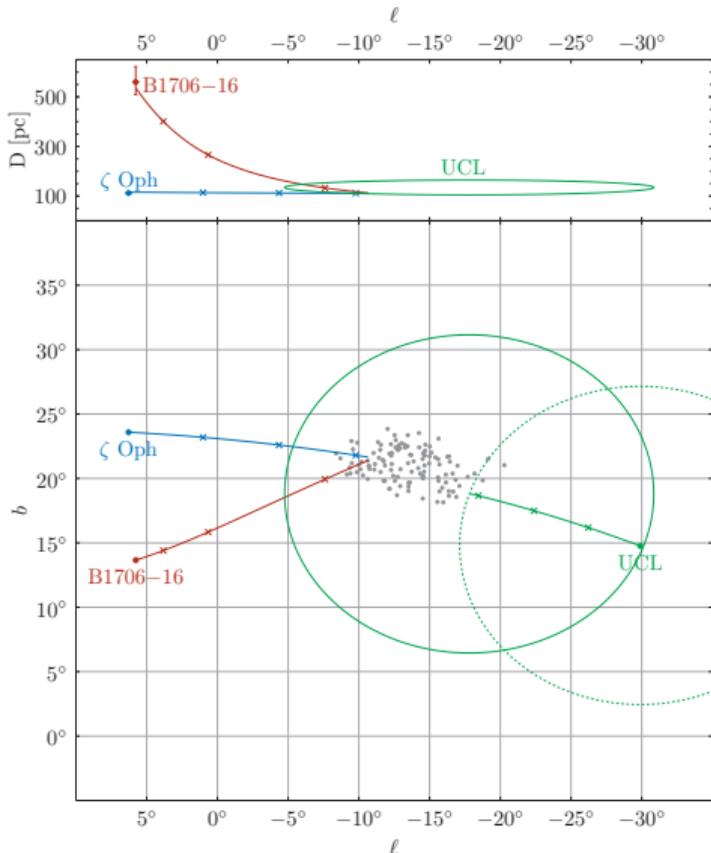
# The runaway $\zeta$ Ophiuchi is the nearest O-type star to Earth

## Observational constraints

- $d \simeq 107 \pm 4$  pc
- $M \simeq 20 M_{\odot}$
- $20 \text{ km s}^{-1} \lesssim v_{\text{sys}} \lesssim 50 \text{ km s}^{-1}$
- $v \sin(i) \gtrsim 350 \text{ km s}^{-1}$
- $(T_{\text{eff}}, L)$  position
- $Z \lesssim Z_{\odot}$ ,  ${}^4\text{He}$ - and  ${}^{14}\text{N}$ -rich, normal  ${}^{12}\text{C}$  and  ${}^{16}\text{O}$
- ✗ Weak wind problem:

$$|\dot{M}_{\text{obs}}| \simeq 10^{-8.8} \ll |\dot{M}_{\text{th}}| \simeq 10^{-6.8} [M_{\odot} \text{yr}^{-1}]$$

# $\zeta$ Oph is a “widowed” star: we can trace it back to a neutron star



A nearby recent supernova that ejected the runaway star  $\zeta$  Oph, the pulsar PSR B1706-16, and  $^{60}\text{Fe}$  found on Earth

R. Neuhäuser,<sup>1\*</sup> F. Gießler<sup>1</sup>, and V.V. Hambaryan<sup>1,2</sup>

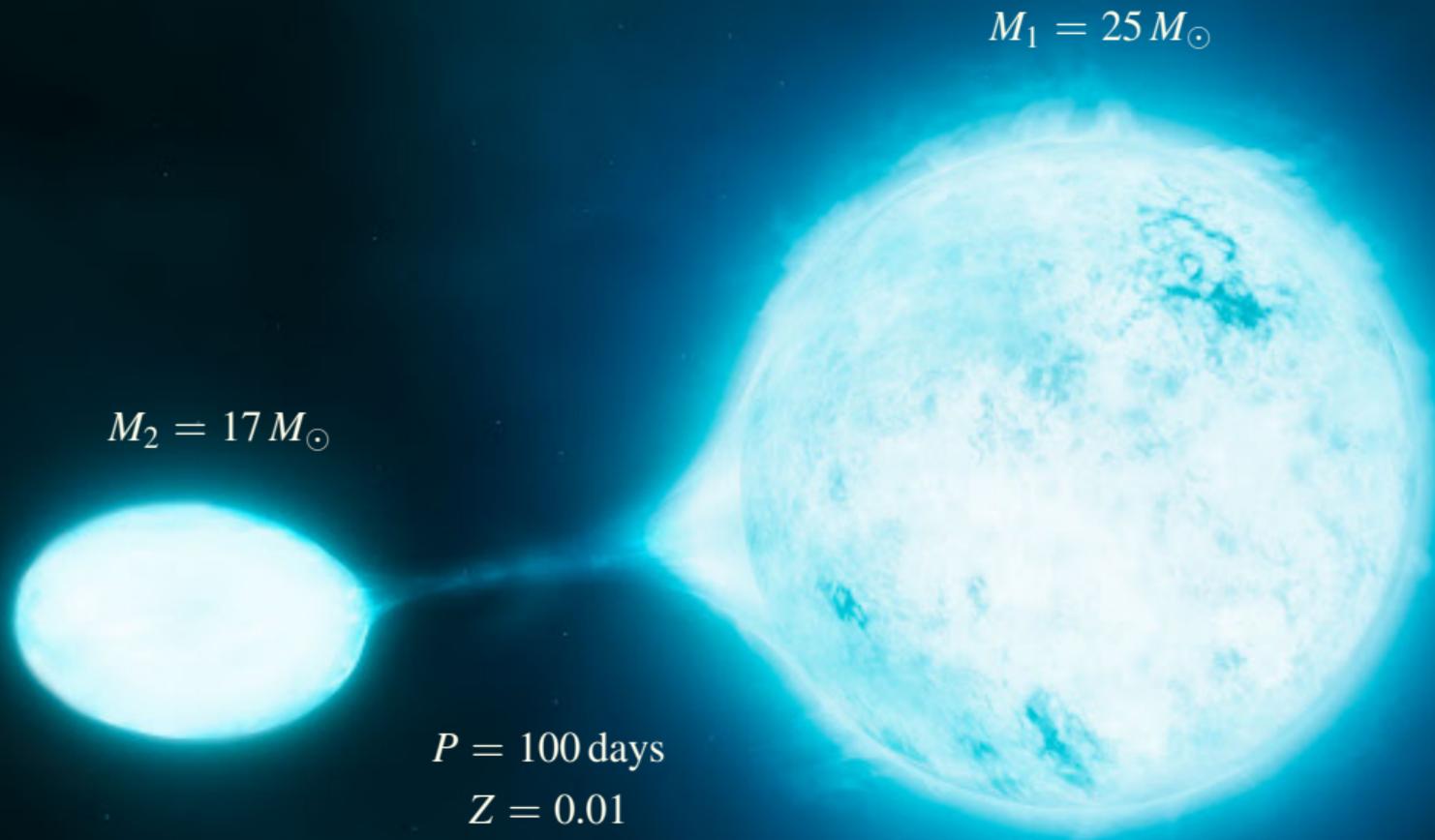
<sup>1</sup> *Astrophysikalisches Institut und Universitäts-Sternwarte Jena, Schillergäßchen 2-3, 07745 Jena, Germany*

<sup>2</sup> *Byurakan Astrophysical Observatory, Byurakan 0213, Aragatsozn, Armenia*

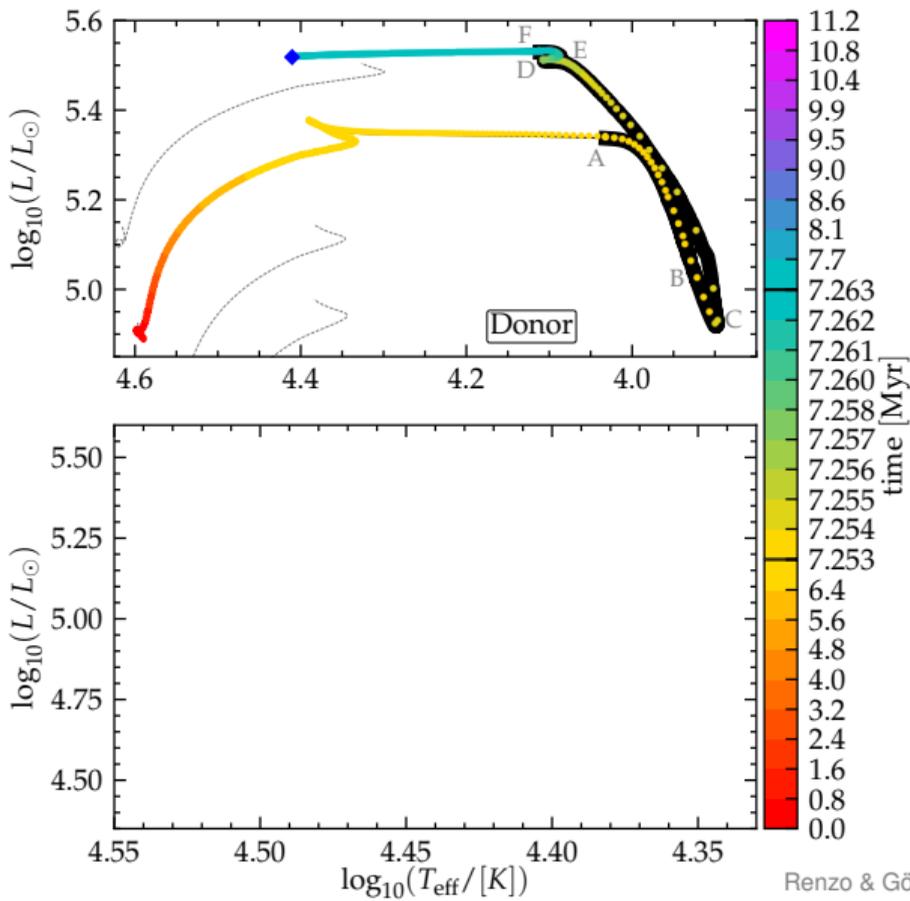
Accepted 2019 Sep 10. Received 2019 Sep 3; in original form 2019 July

SN explosion  $\sim 1.78 \pm 0.21$  Myr ago

# Self-consistent MESA model

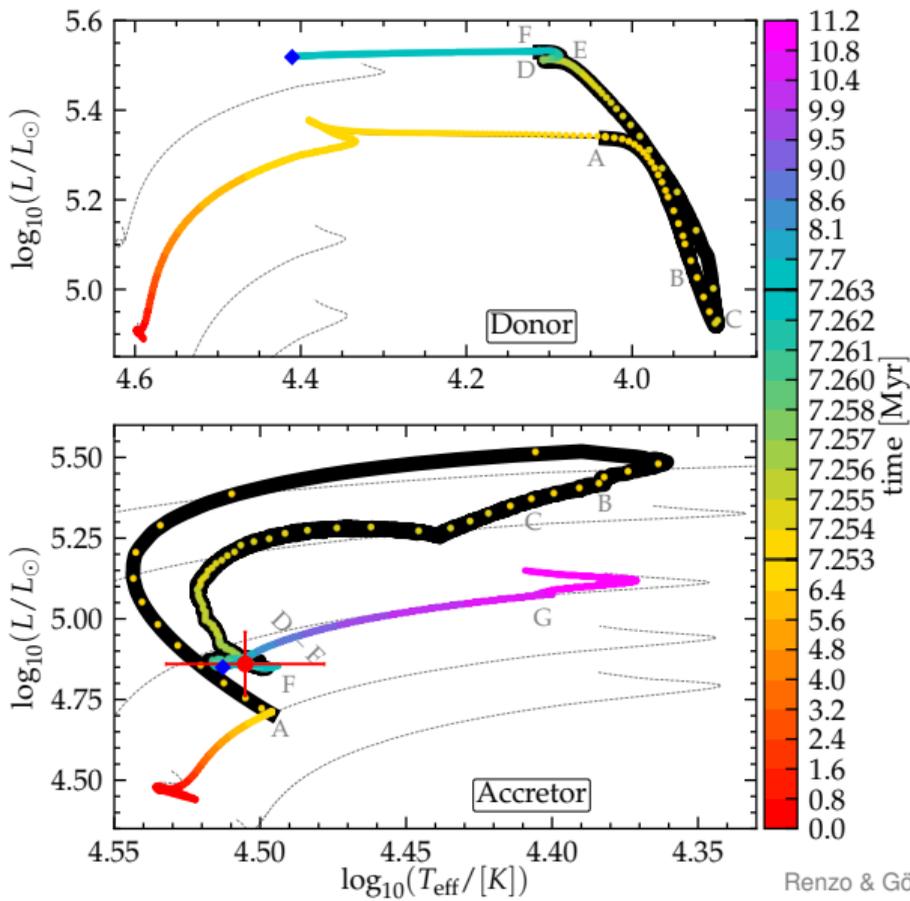


# Hertzprung-Russel diagram of both stars: the donor



**Roche lobe overflow is short**  
But has long-lasting impact on **both** stars.

# Hertzprung-Russel diagram of both stars: the donor & the accretor



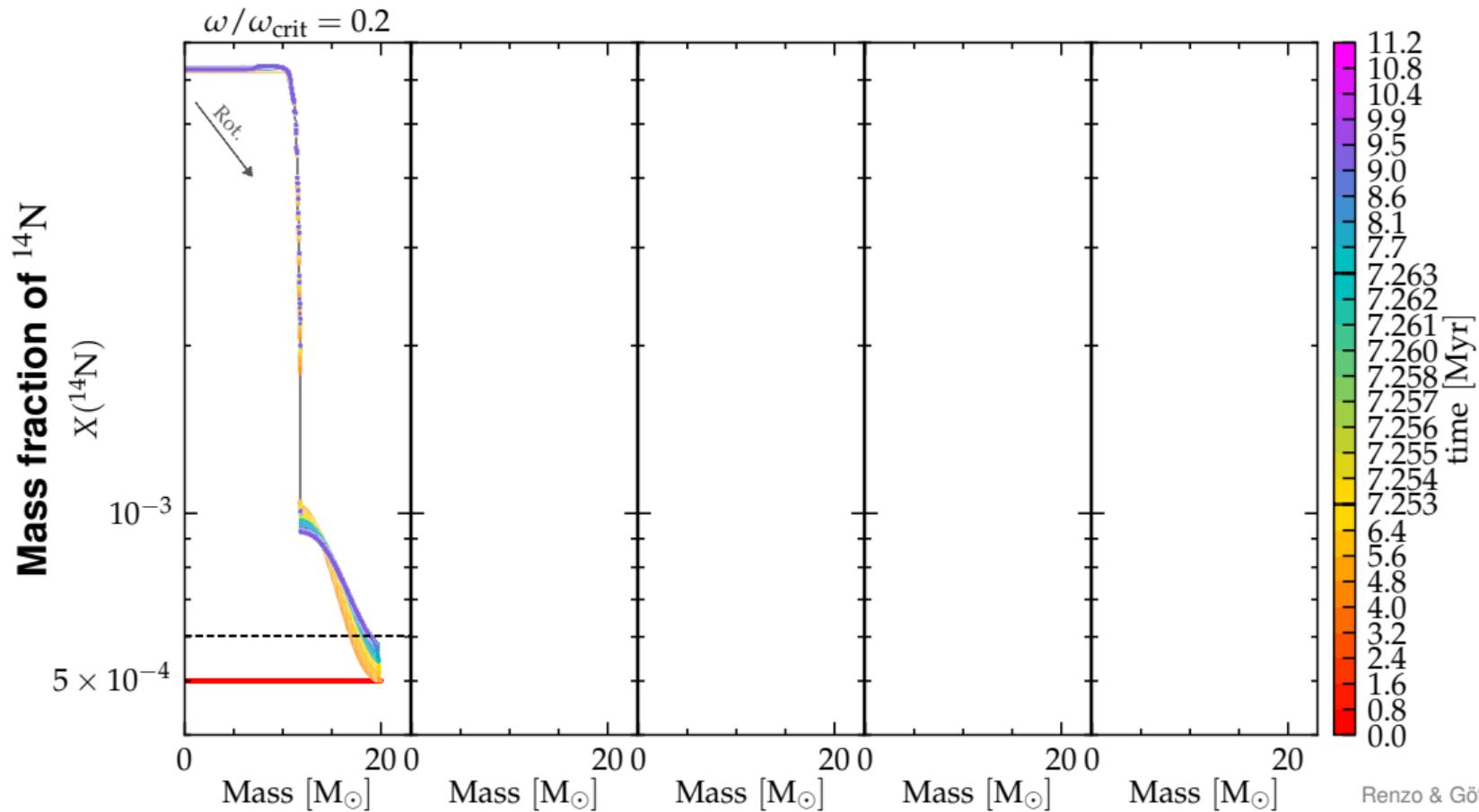
**Roche lobe overflow is short**  
But has long-lasting impact on **both** stars.

## Appearance of “widowed” stars

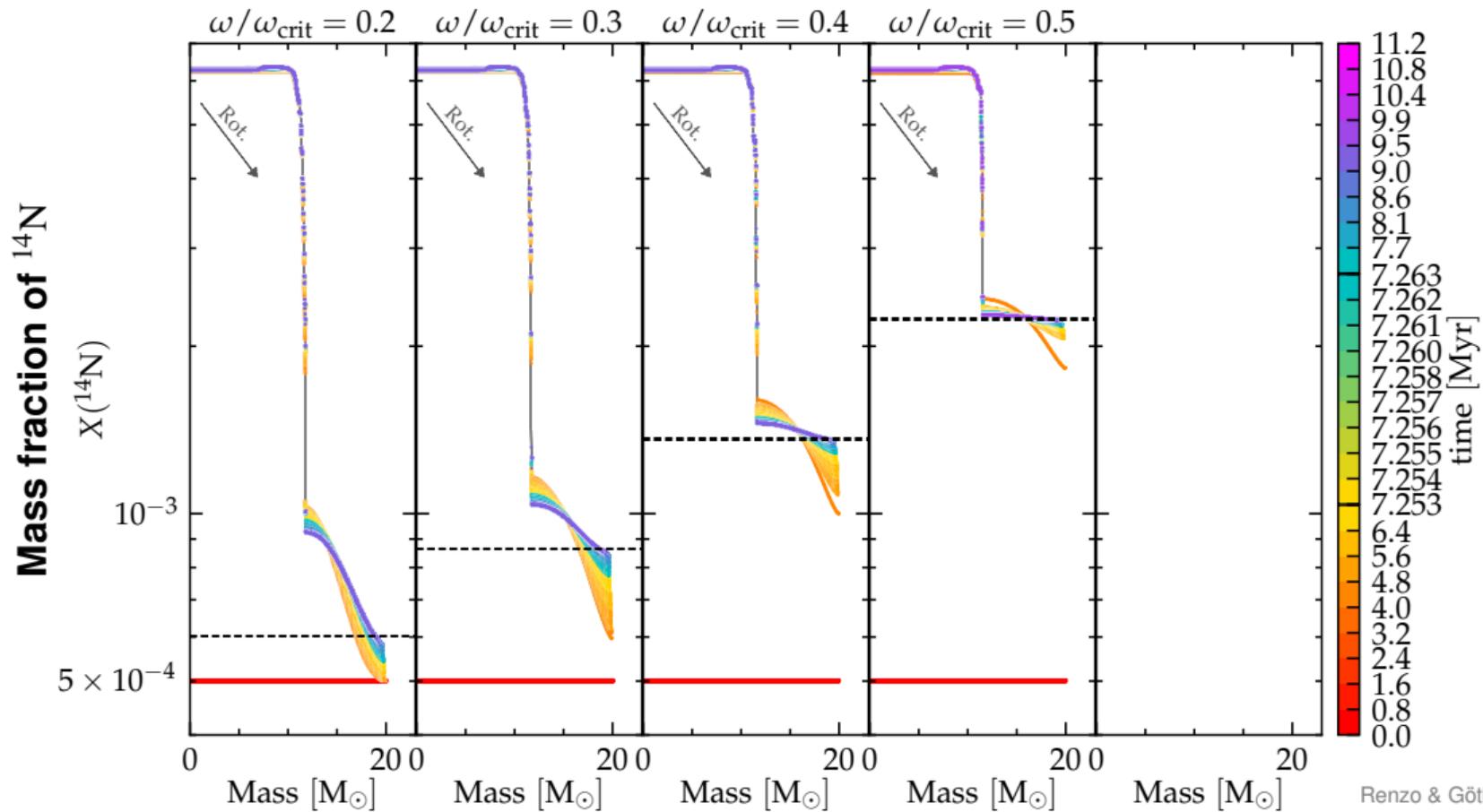
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$^{14}\text{N}$  as a tracer of chemical composition

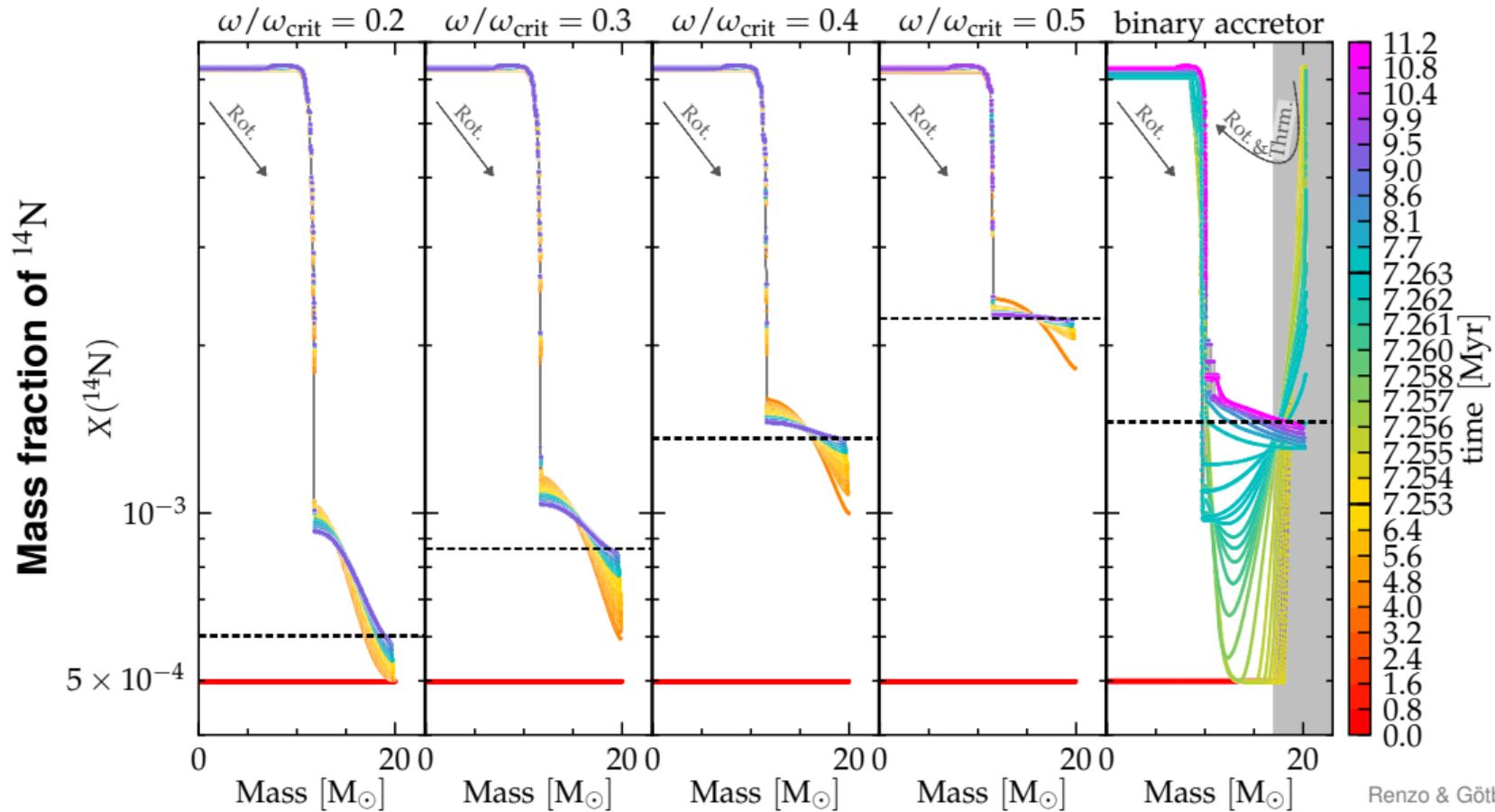
# Composition profile: comparison with rotating single stars



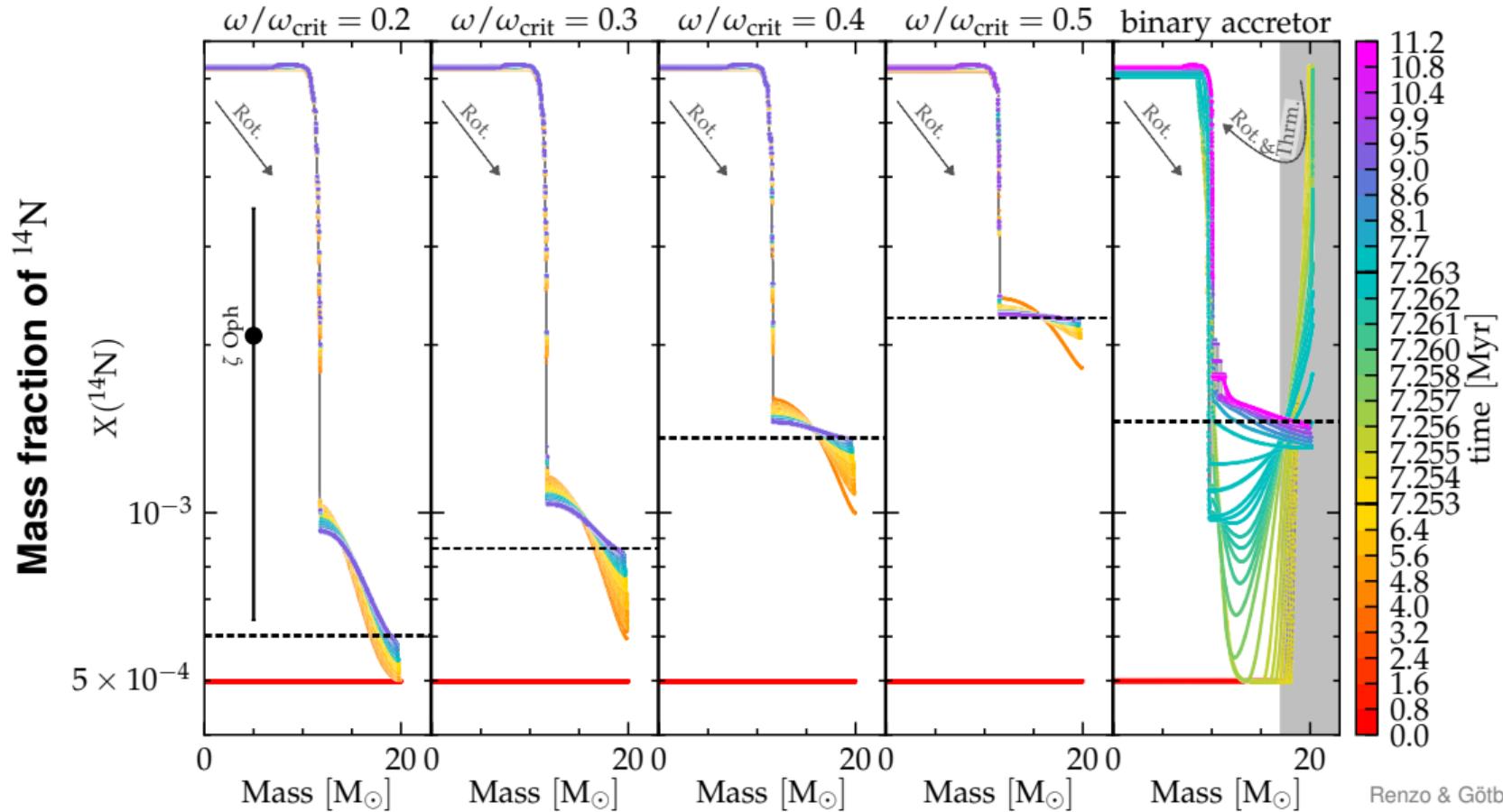
# Composition profile: comparison with rotating single stars



# Composition profile: comparison with rotating single stars



# Composition profile: comparison with rotating single stars

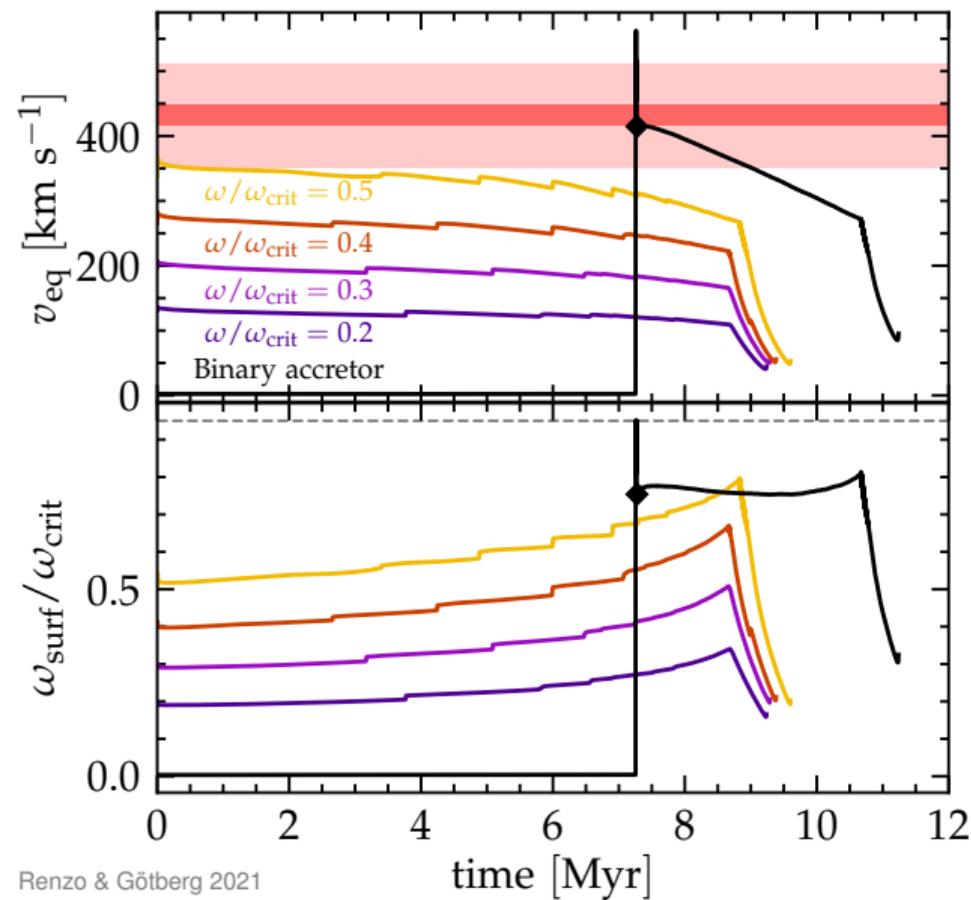


# Appearance of “widowed” stars

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Rotation

## Surface rotation rate

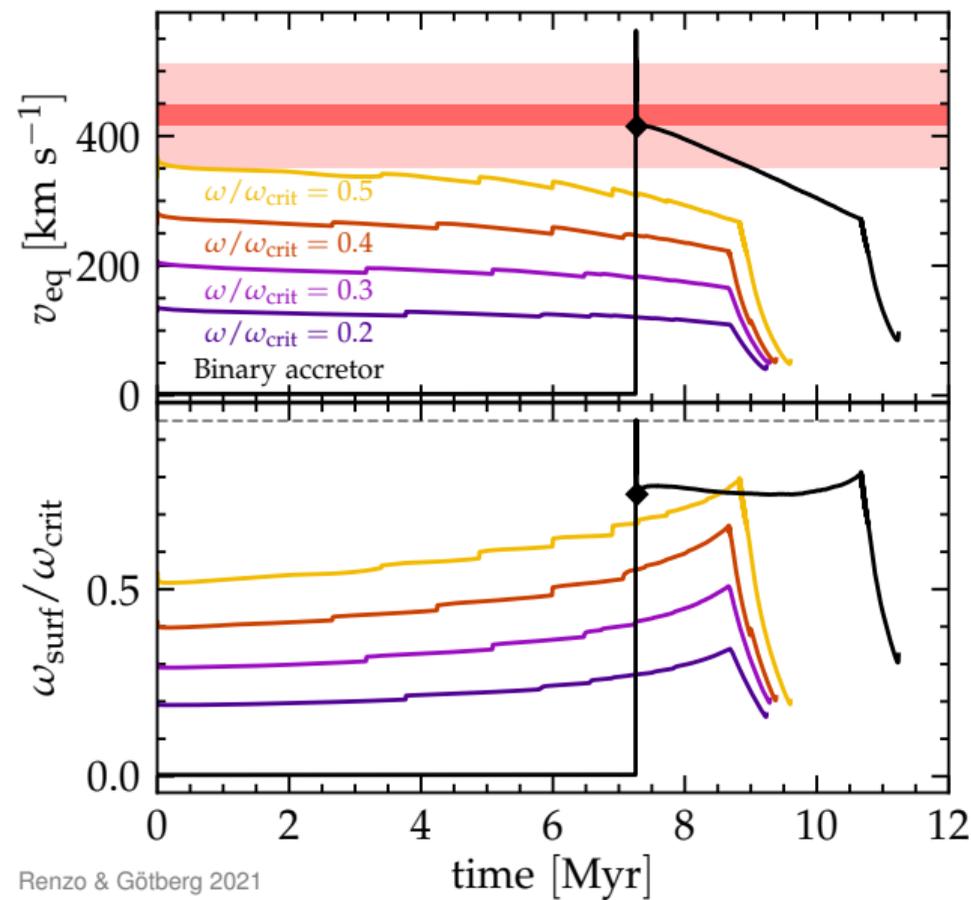


- but overestimating by  $\sim 100\times$  wind mass loss!

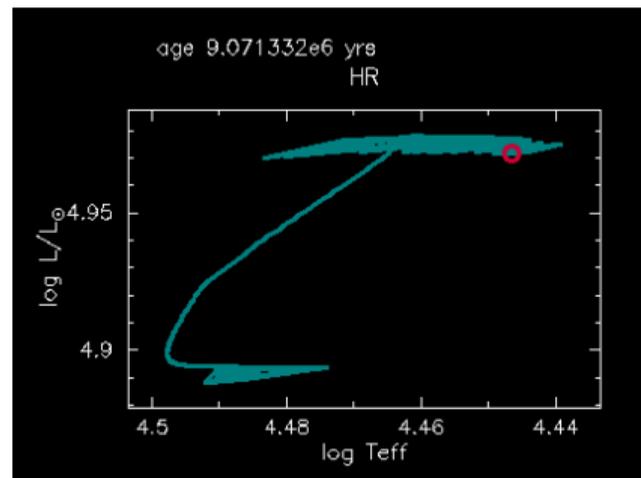
$$\omega_{\text{crit}} = \sqrt{\left(1 - \frac{L}{L_{\text{Edd}}}\right) \frac{GM}{R_{\text{eq}}^3}}$$

Gravity = Centrifugal forces  
at equator

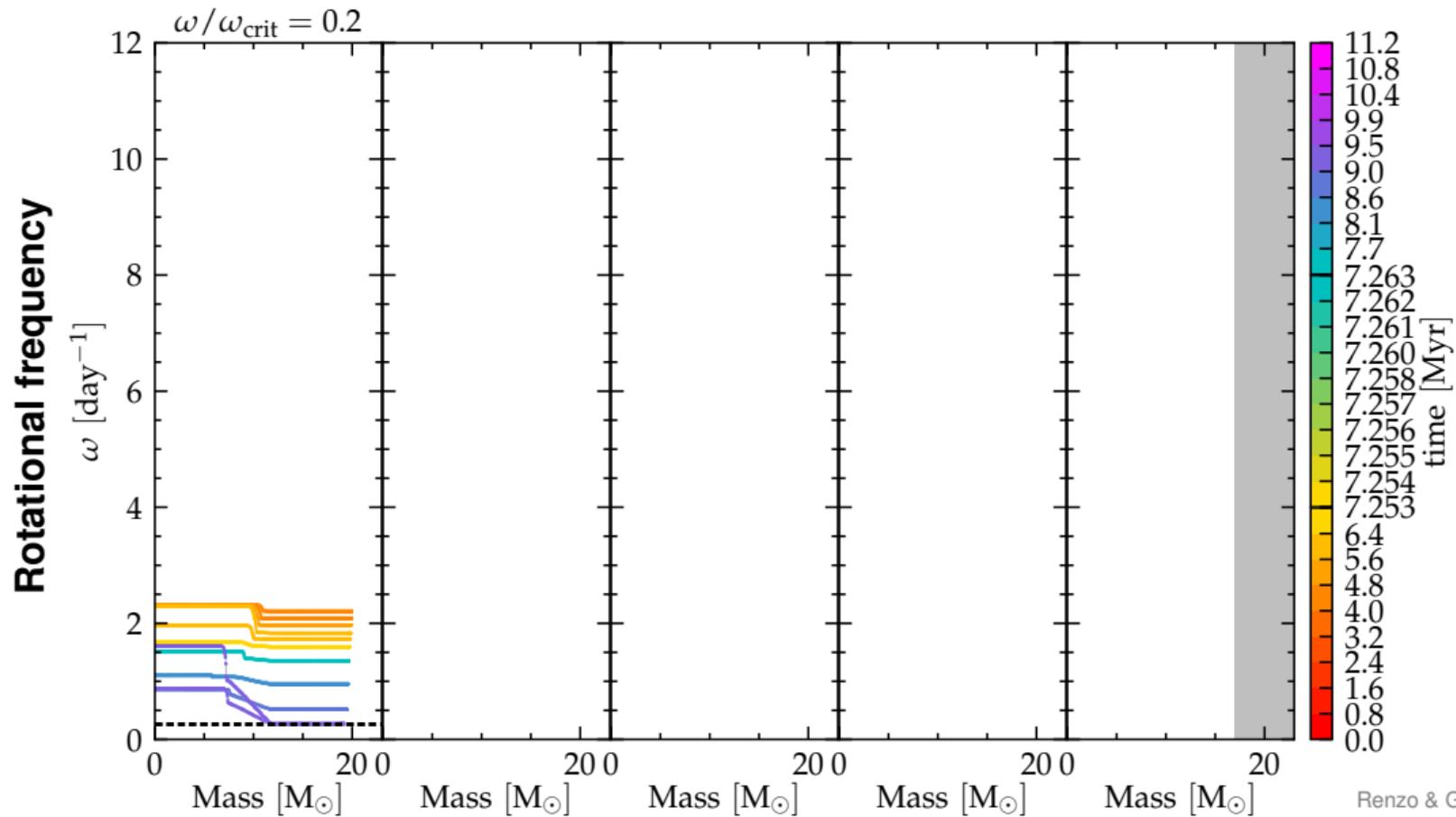
# Surface rotation rate



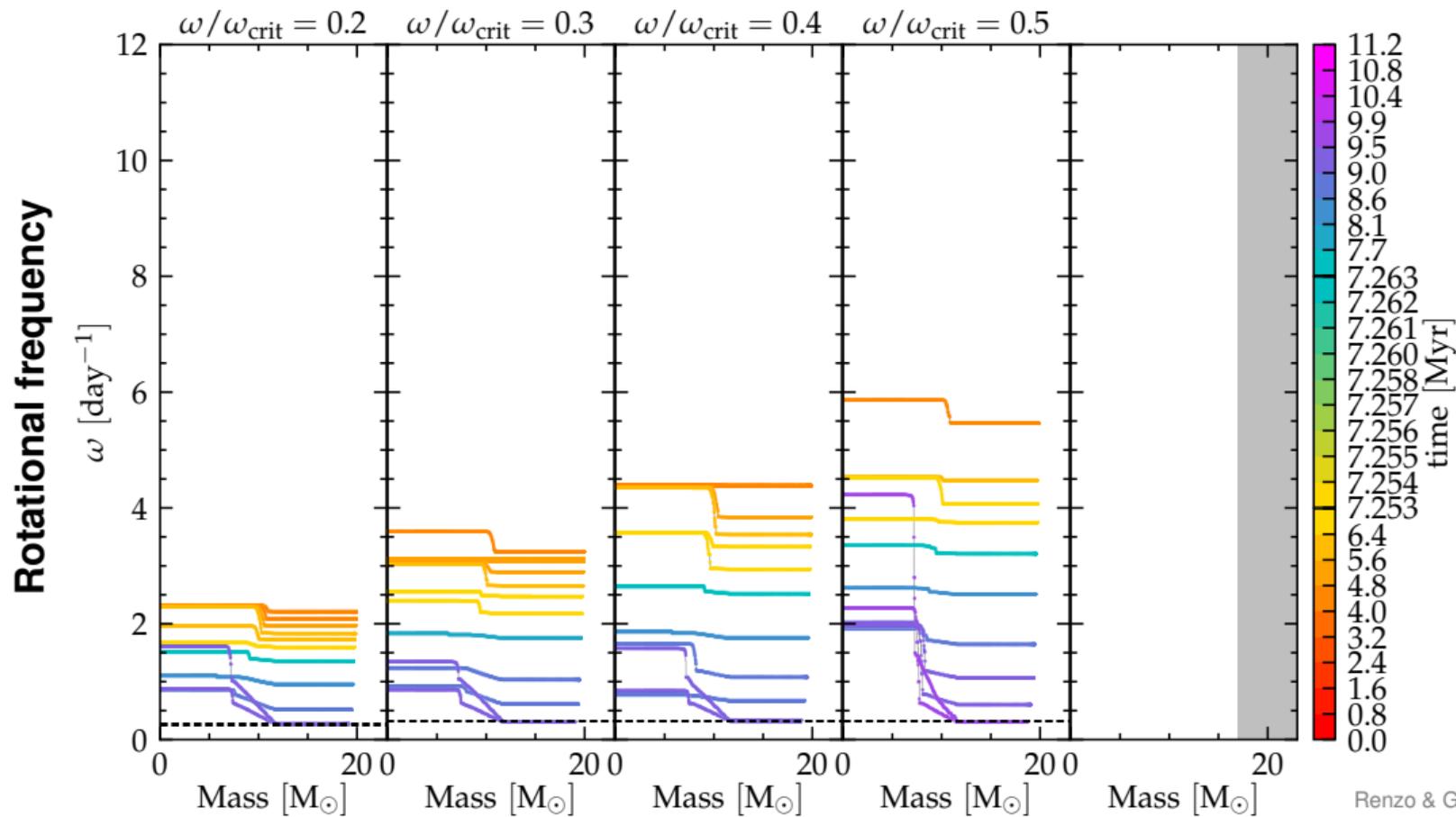
- but overestimating by  $\sim 100\times$  wind mass loss!
- Decreasing the wind yields  
 $\omega/\omega_{\text{crit}} > 1$



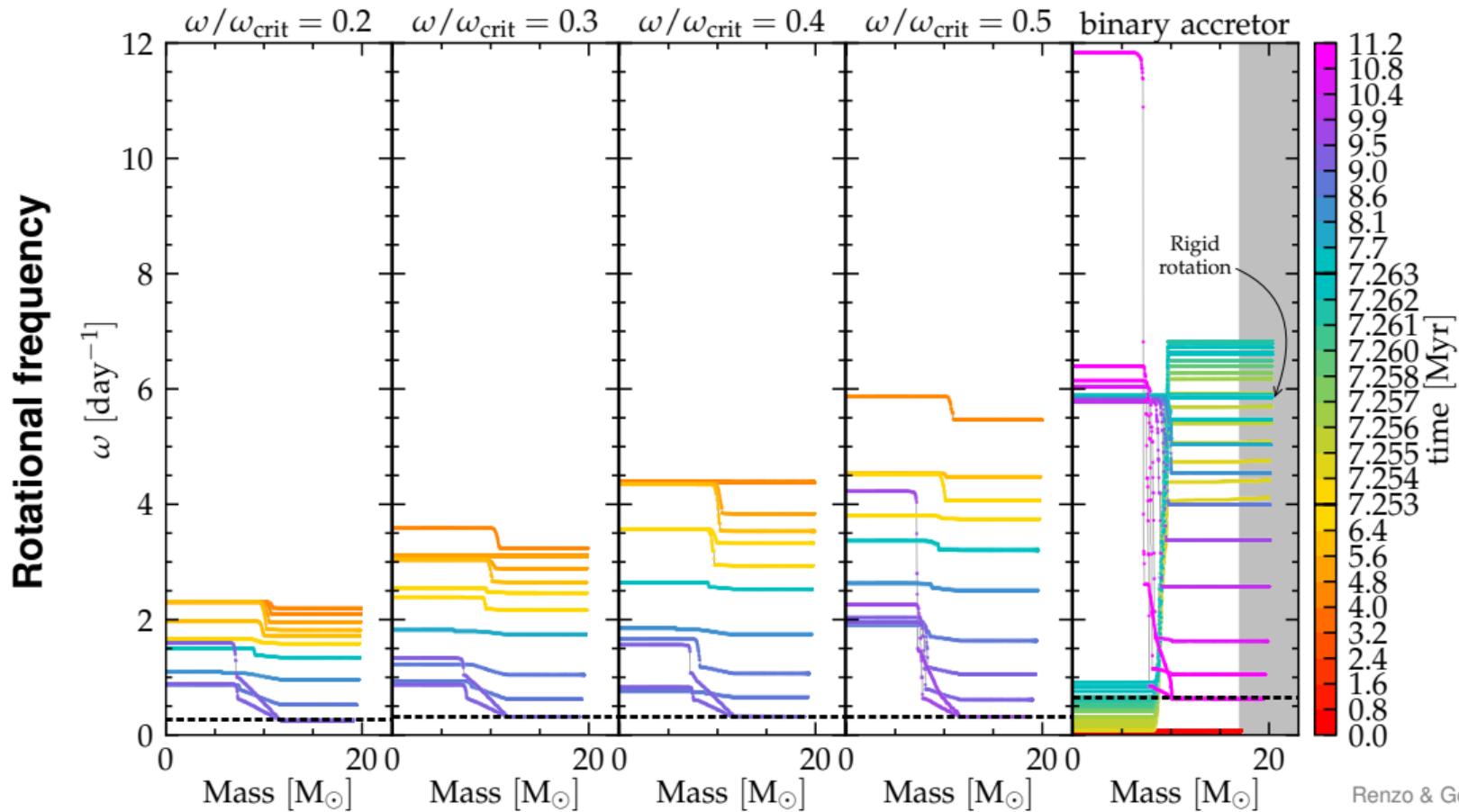
# Internal rotational profile: single stars



# Internal rotational profile: single stars



# Internal rotational profile: accretor

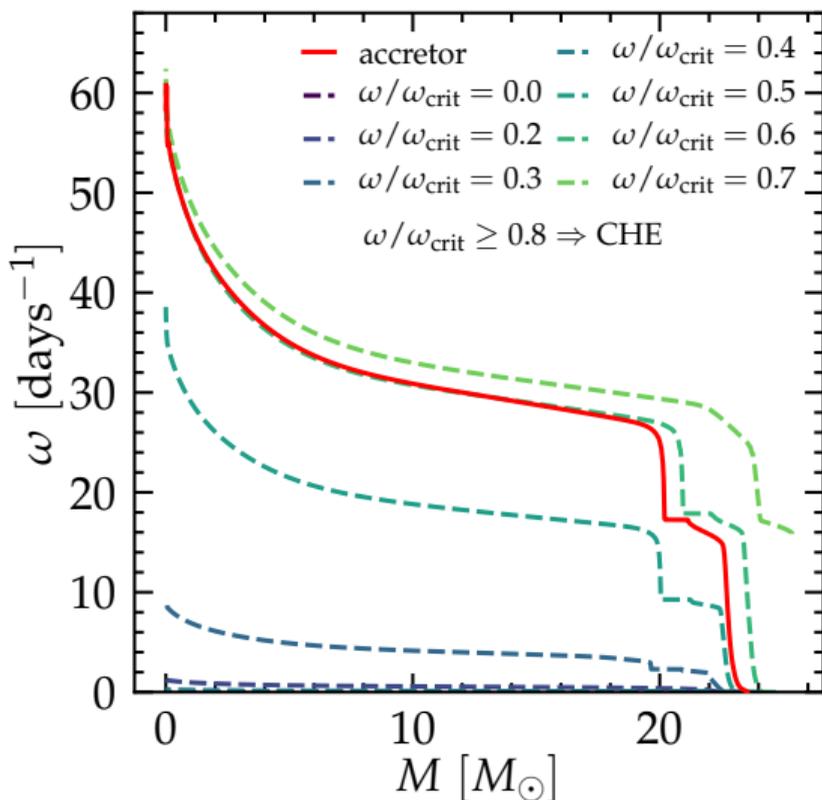


## **Generalization to BH progenitors**

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## Preliminary: accretion spin-up for BH progenitors

$40 M_{\odot}$  stars at  $Z \simeq Z_{\odot}/10$  evolved until carbon depletion with Spruit-Tayler dynamo



$50 M_{\odot} + 40 M_{\odot}$ , initial separation  $200 R_{\odot}$

The 2<sup>nd</sup> BH might be fast spinning even without tidal interactions

## **Conclusions**

---

**Stable mass transfer**

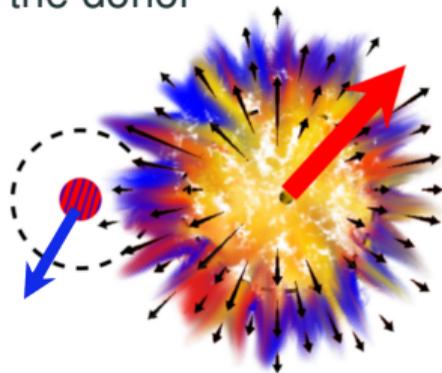
## Take home points

Renzo *et al.* 2019b

- Most massive binaries disrupted at 1<sup>st</sup> core-collapse
- Masses of widowed stars can inform BH kicks
- Most are slow moving walkaway, some are runaway

Renzo & Götzberg 2021

- Widowed stars are modified by binary interactions
  - Surface abundance influenced by matter from the donor
  - Rotation profile of widowed stars unlike single rotating stars
- ⇒ implications for long GRB and BH spins



## Mass transfer in binary systems can be

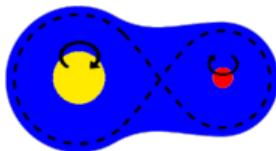
dynamically stable:

testing models of massive accretor stars with  $\zeta$  Ophiuchi



**dynamically unstable:**

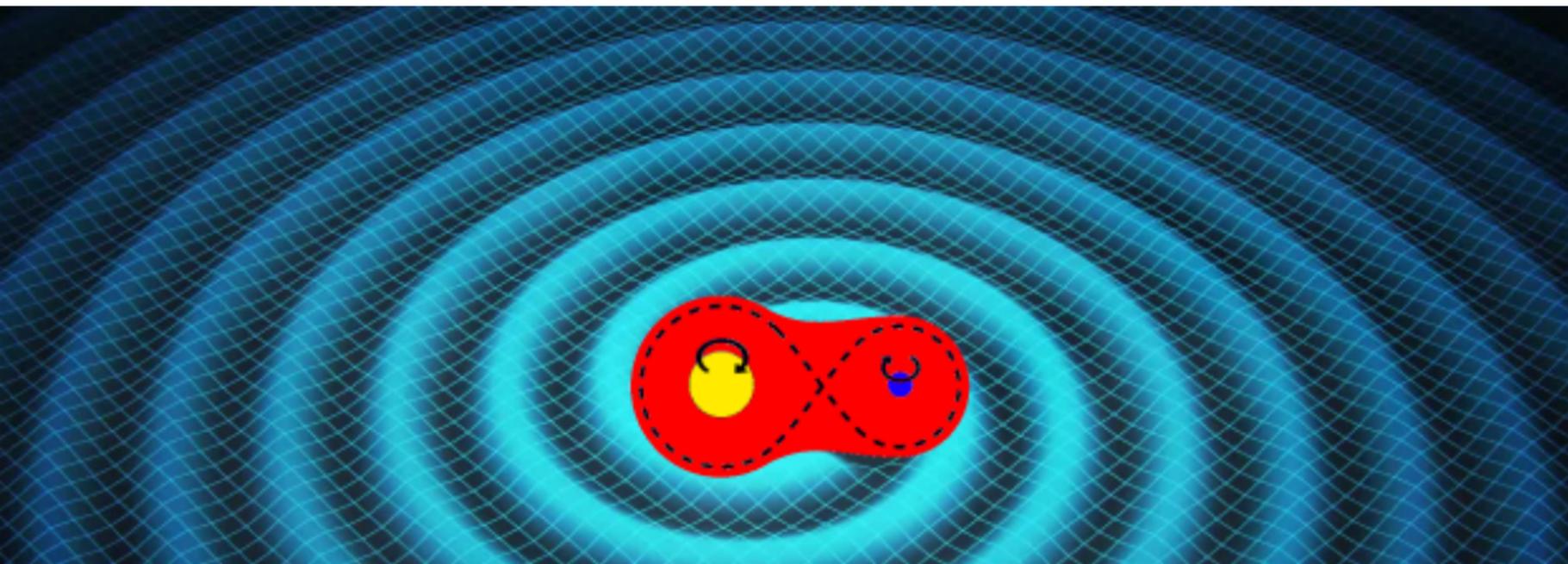
can LISA detect GW from common envelope evolution?



# Prospects of gravitational-waves detections from common-envelope evolution with LISA

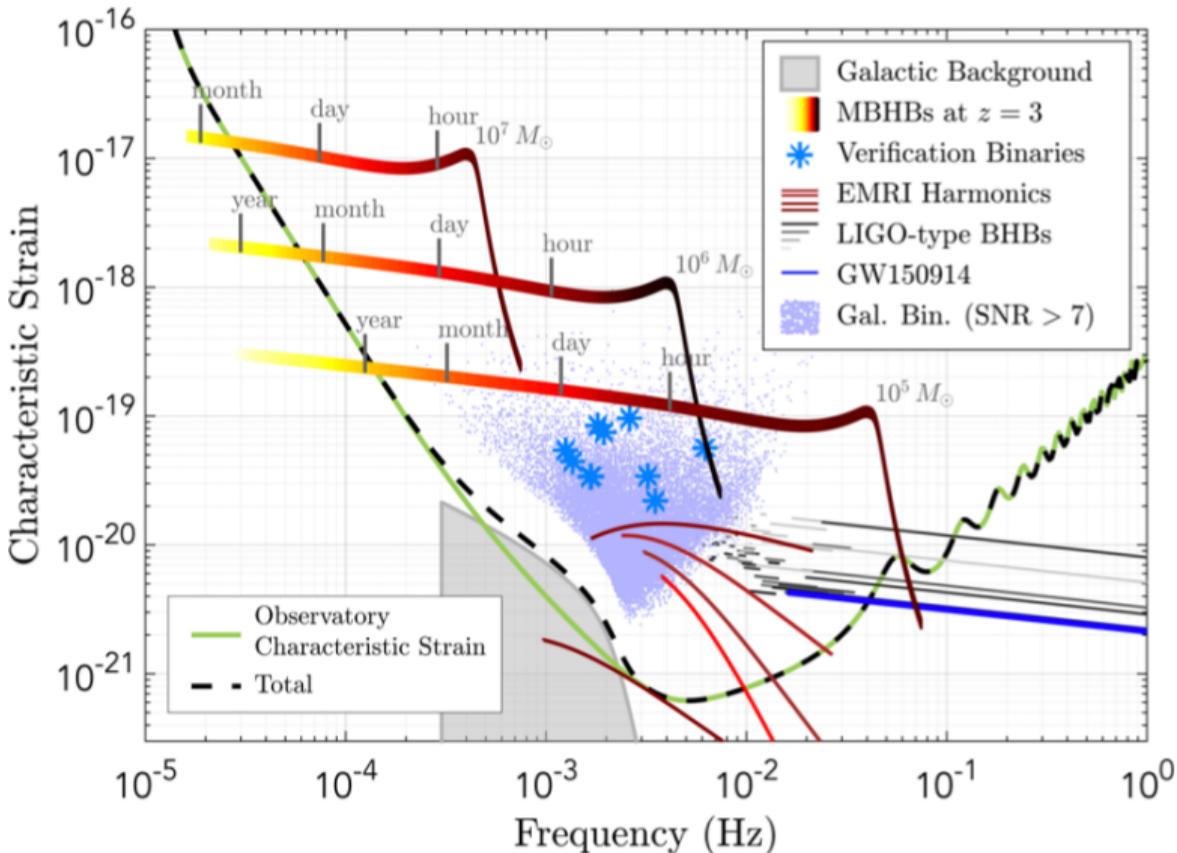
**Mathieu Renzo**, T. Callister, K. Chatziioannou, L. van Son, C. M. F. Mingarelli, M. Cantiello,  
K. E. S. Ford, B. McKernan, and G. Ashton

arXiv:2102.00078, accepted ApJ



# LISA can see Galactic double white dwarfs formed via common envelope

⇐ PTA



LIGO/Virgo ⇒

## Common Envelope Evolution

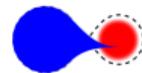
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**Is *not* GW-driven!**

**But GW passively trace the dynamics**

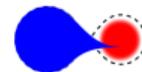
# Common envelope evolution in one slide

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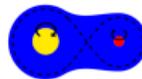


a. Mass transfer becomes dynamically unstable

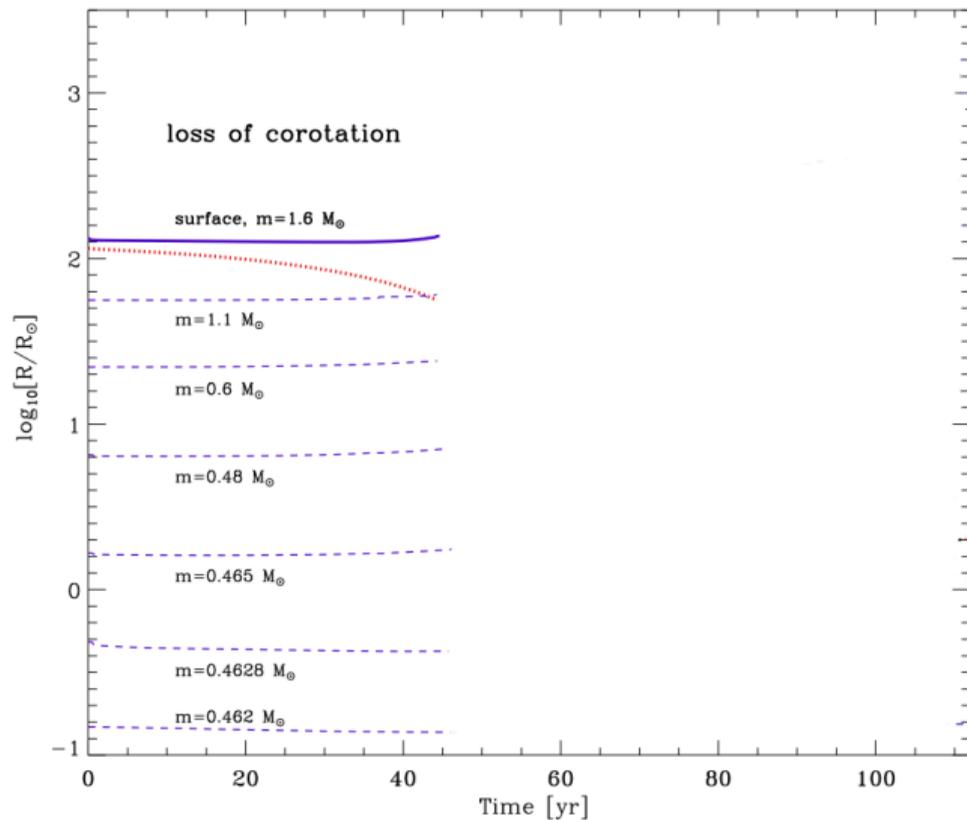
# Common envelope evolution in one slide



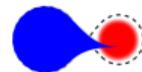
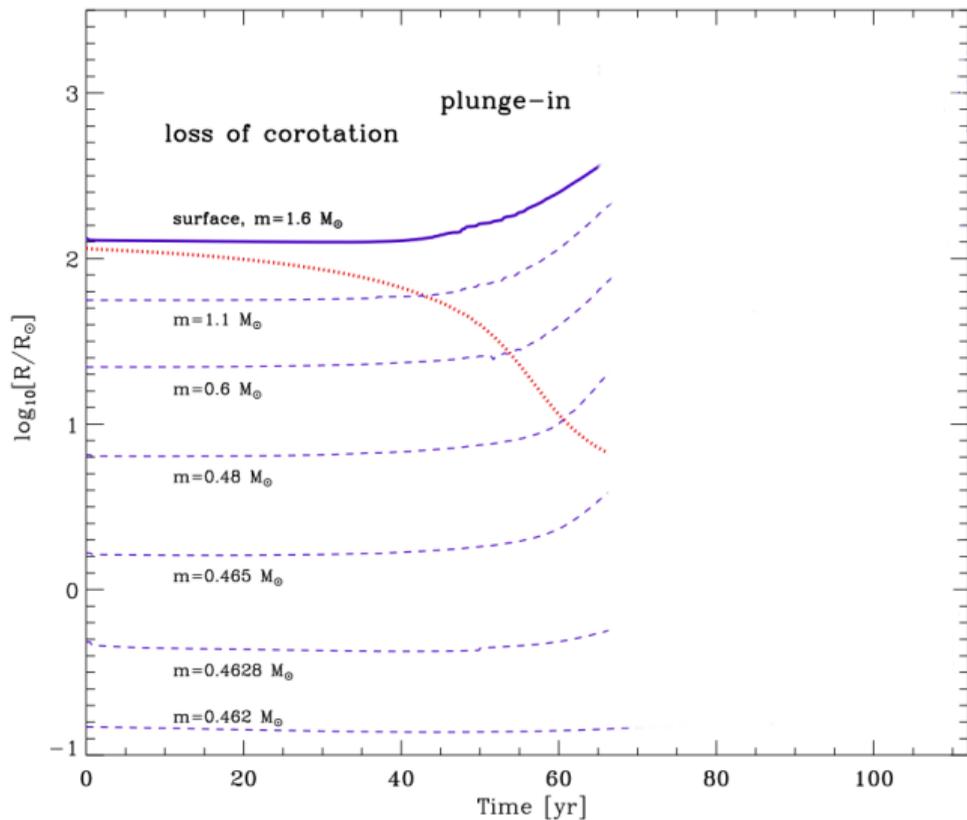
a. Mass transfer becomes dynamically unstable



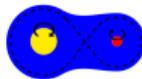
b. Loss of corotation between the cores and the envelope



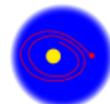
# Common envelope evolution in one slide



a. Mass transfer becomes dynamically unstable

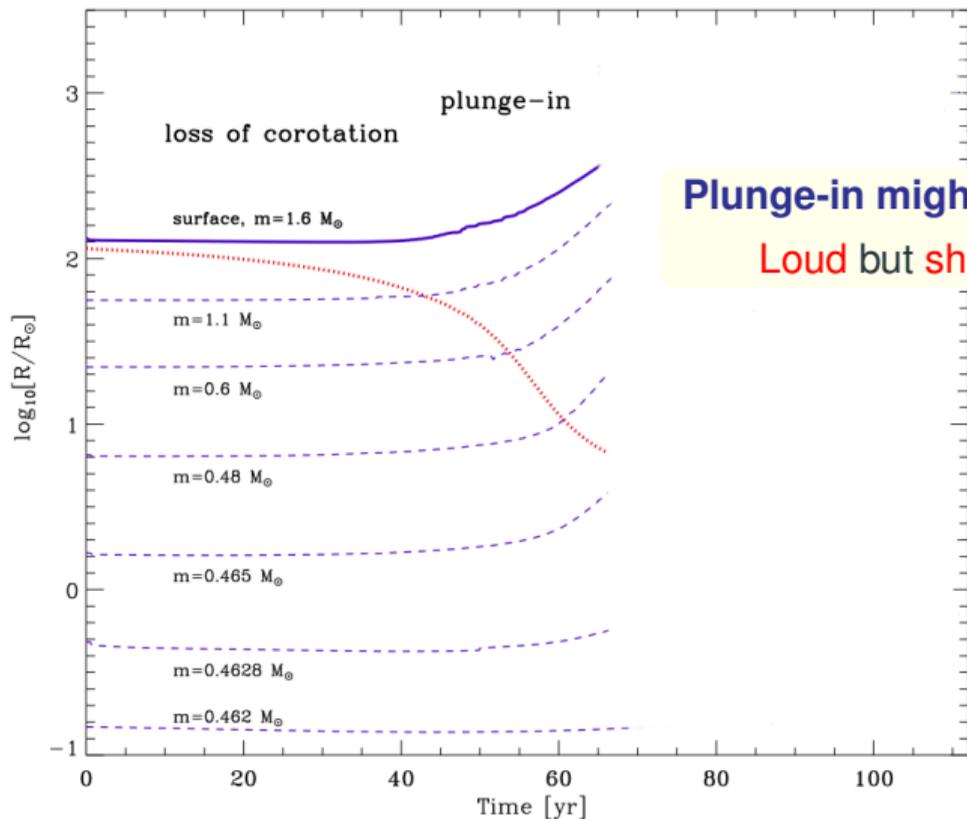


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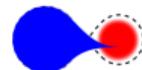
c. Dynamical plunge-in

# Common envelope evolution in one slide

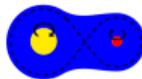


Plunge-in might be detectable

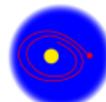
Loud but short and rare



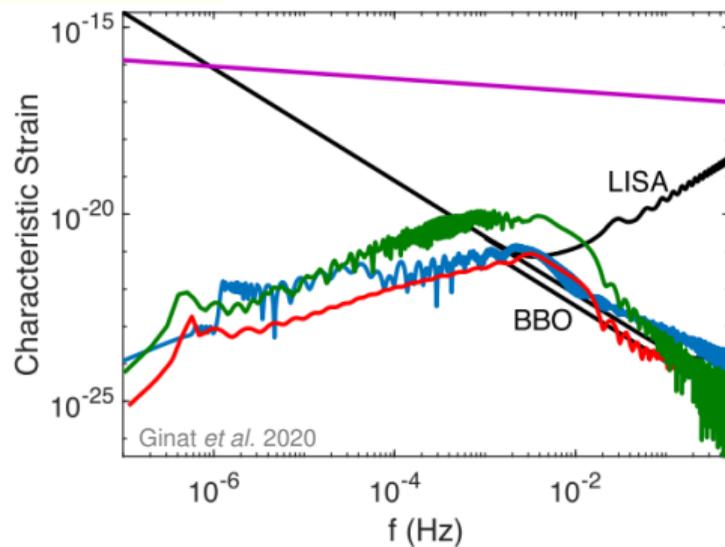
a. Mass transfer becomes dynamically unstable



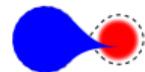
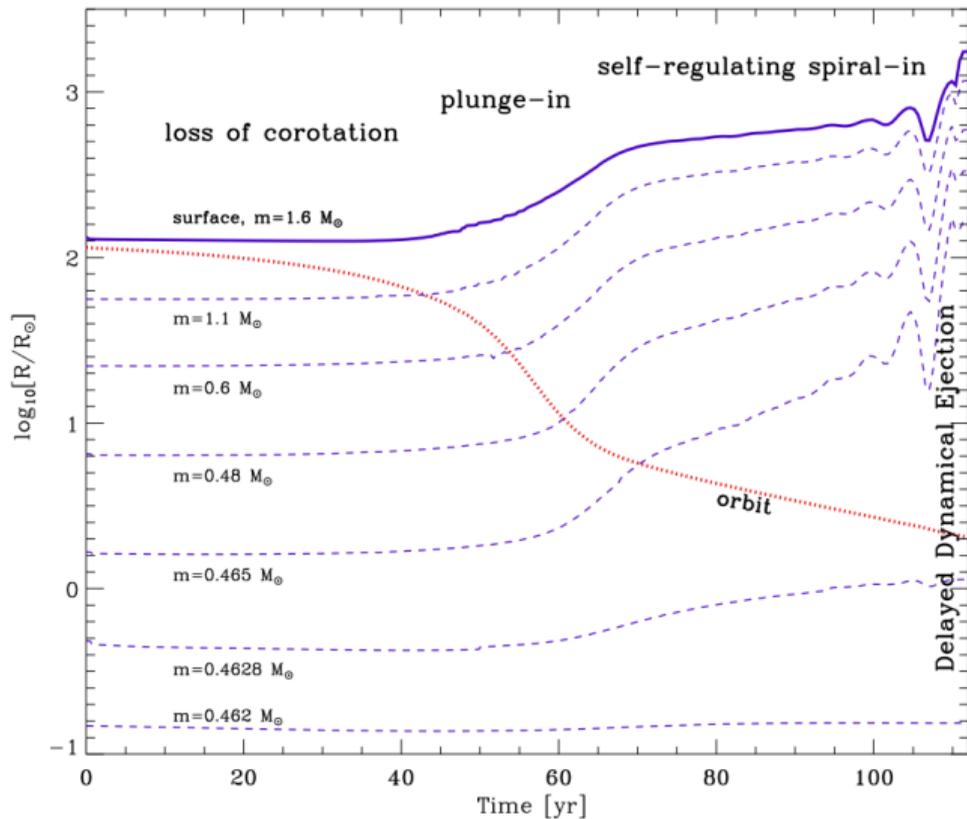
b. Loss of corotation between the cores and the envelope



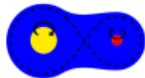
c. Dynamical plunge-in



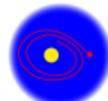
# Common envelope evolution in one slide



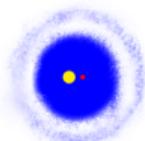
a. Mass transfer becomes dynamically unstable



b. Loss of corotation between the cores and the envelope

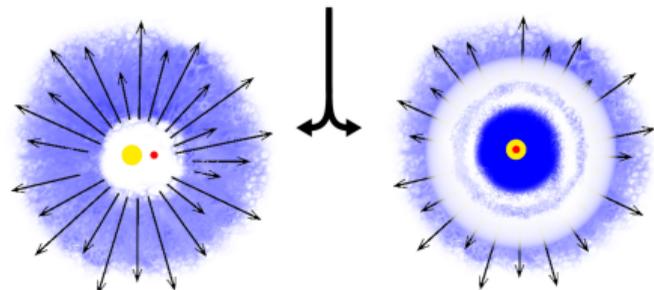
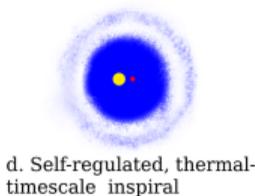
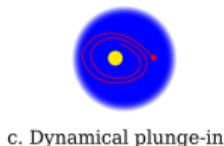
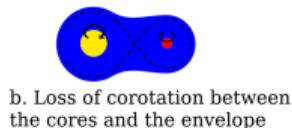
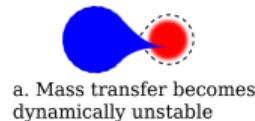
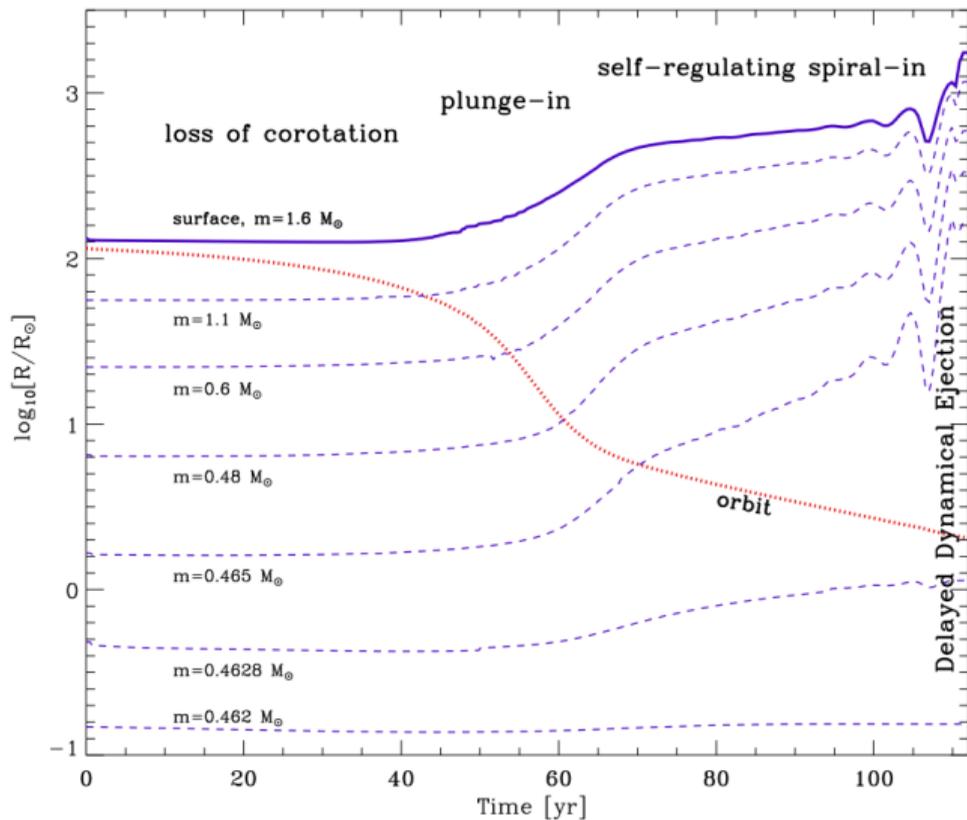


c. Dynamical plunge-in

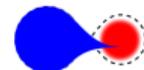
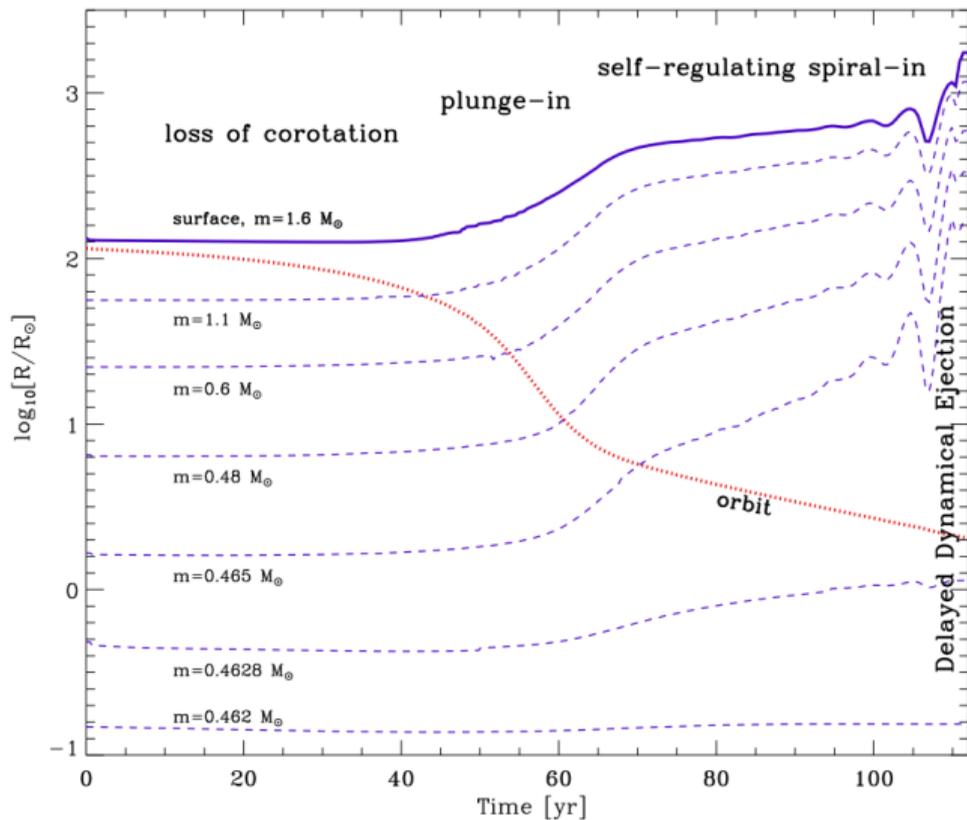


d. Self-regulated, thermal-timescale inspiral

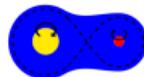
# Common envelope evolution in one slide



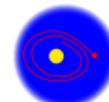
# Common envelope evolution in one slide



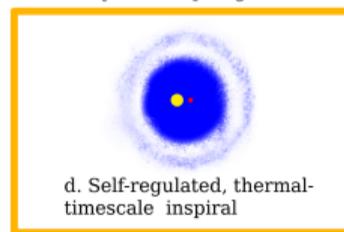
a. Mass transfer becomes dynamically unstable



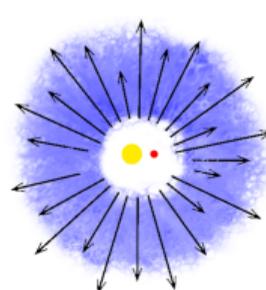
b. Loss of corotation between the cores and the envelope



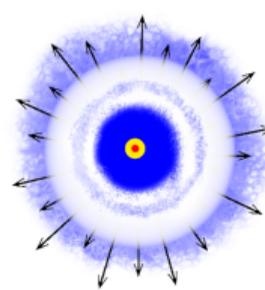
c. Dynamical plunge-in



d. Self-regulated, thermal-timescale inspiral



Common envelope ejection and formation of a short period binary



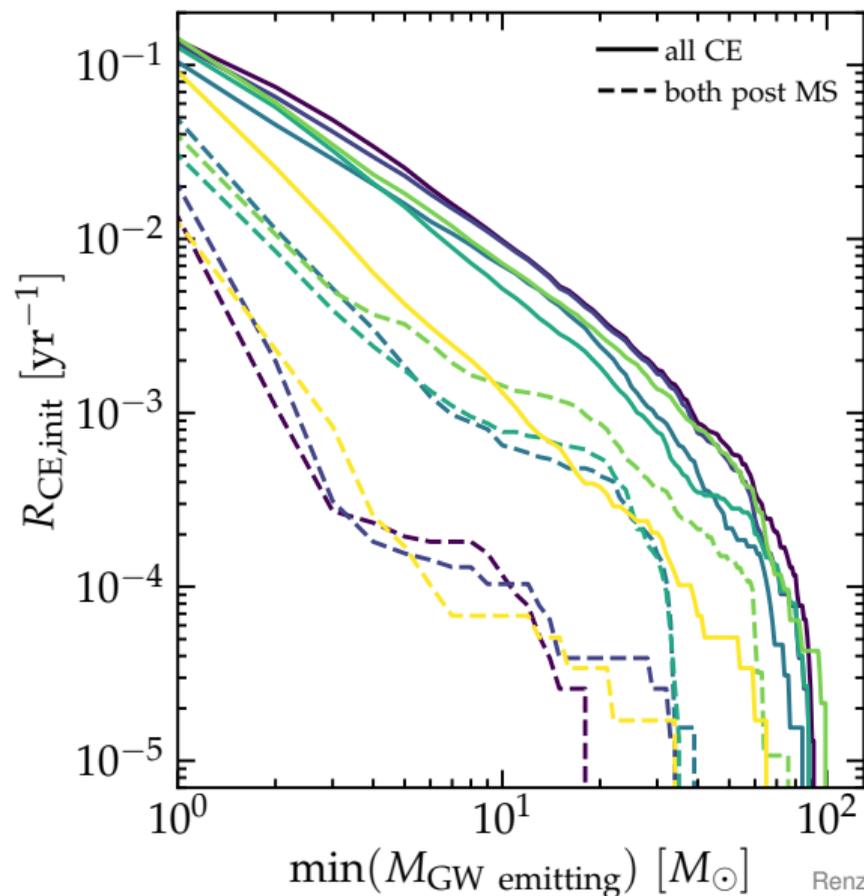
Stellar merger

## How many sources do we expect?

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$$N_{\text{CE}} = R_{\text{CE,init}} \times \Delta t_{\text{CE}}$$

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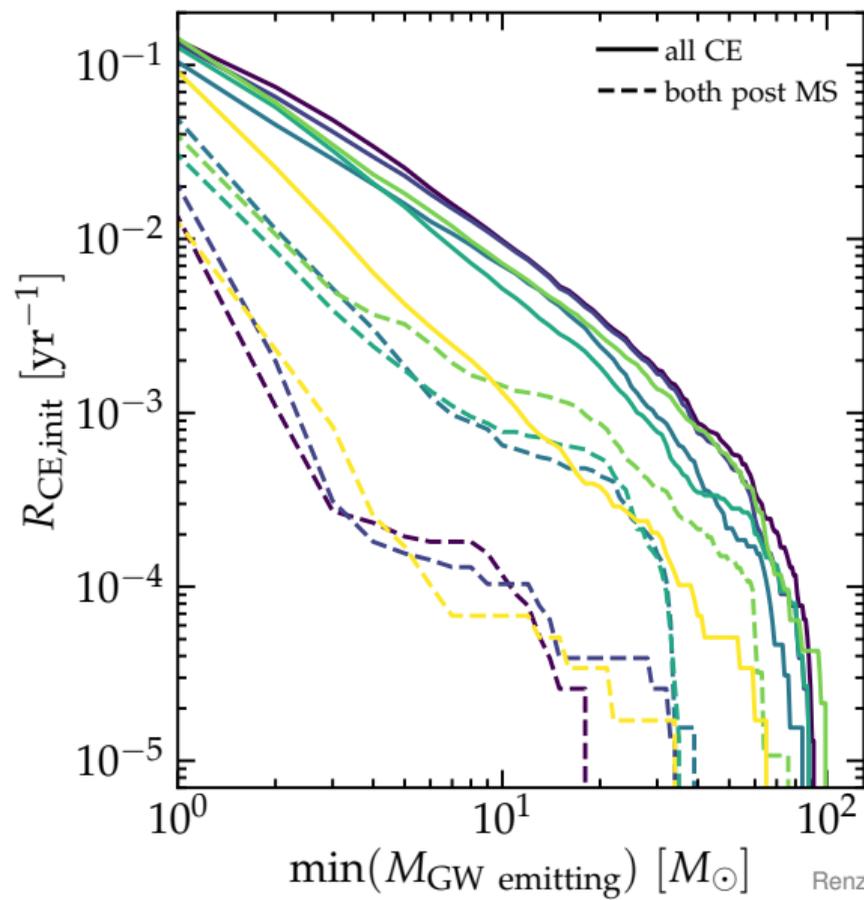
Renzo *et al.* 2021

$$R_{\text{CE,init}} = 0.18_{-0.09}^{+0.02} \quad (0.06_{-0.02}^{+0.03})$$

c.f. LRN rate  $\sim 0.3 \text{ yr}^{-1}$

Kochanek *et al.* 2014, see also Howitt *et al.* 2020

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Renzo *et al.* 2021

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**Duration (in band) is very uncertain**

$$\Delta t_{\text{CE}} \simeq 10^{-2} - 10^5 \text{ years}$$

(e.g., Meyer & Meyer-Hofmeister 1979, Fragos *et al.* 2019, Igoshev *et al.* 2020, Chamandy *et al.* 2020, Law-Smith *et al.* 2020)



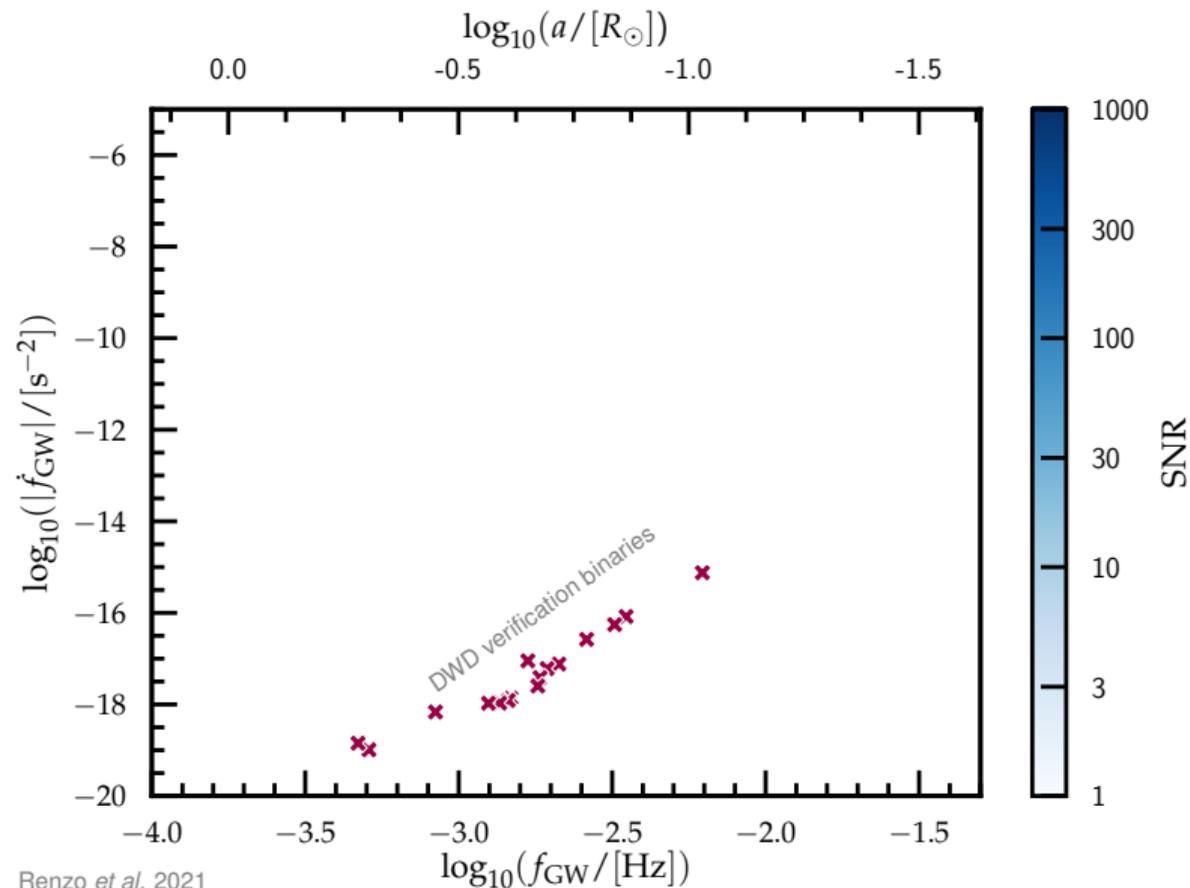
$$0 \lesssim N_{\text{CE}} \lesssim 1000$$



**Could we detect something?**

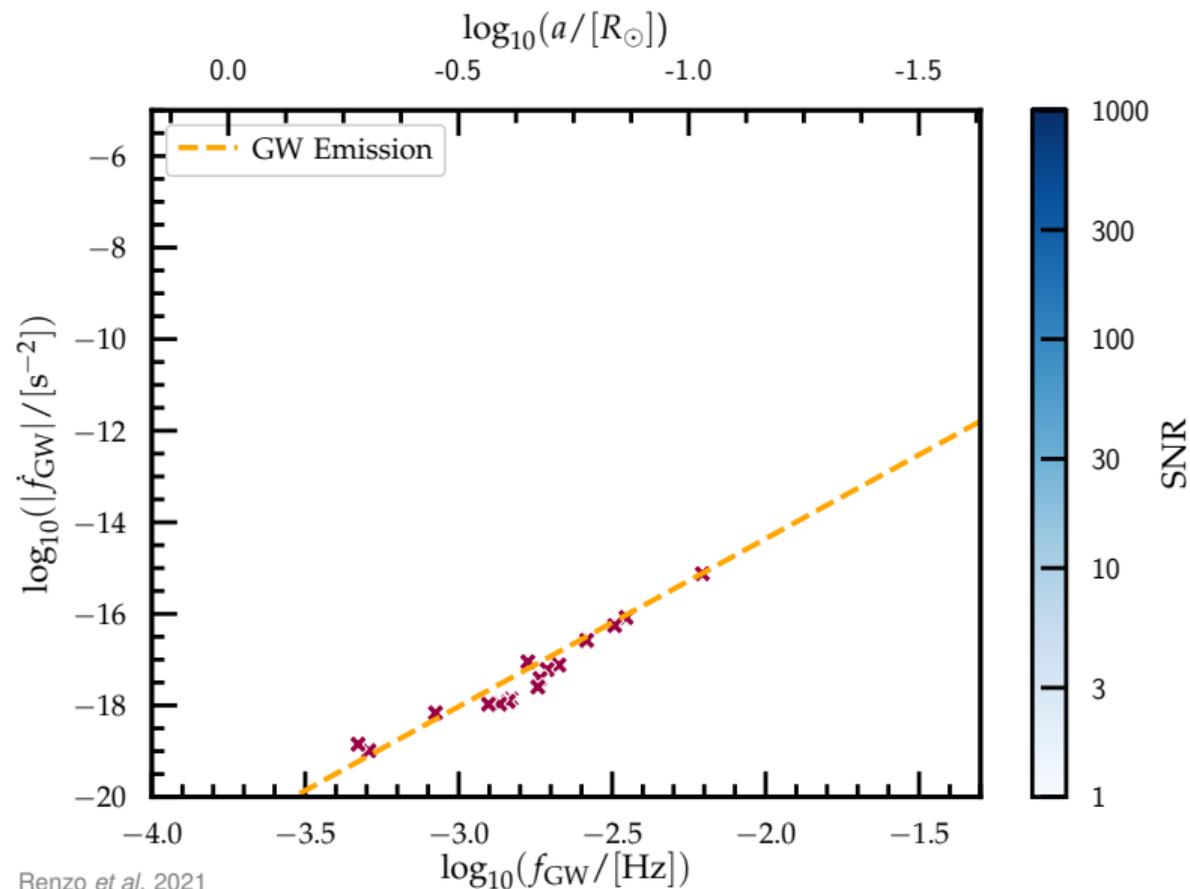
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## Could we see it? An answer not relying on a specific model



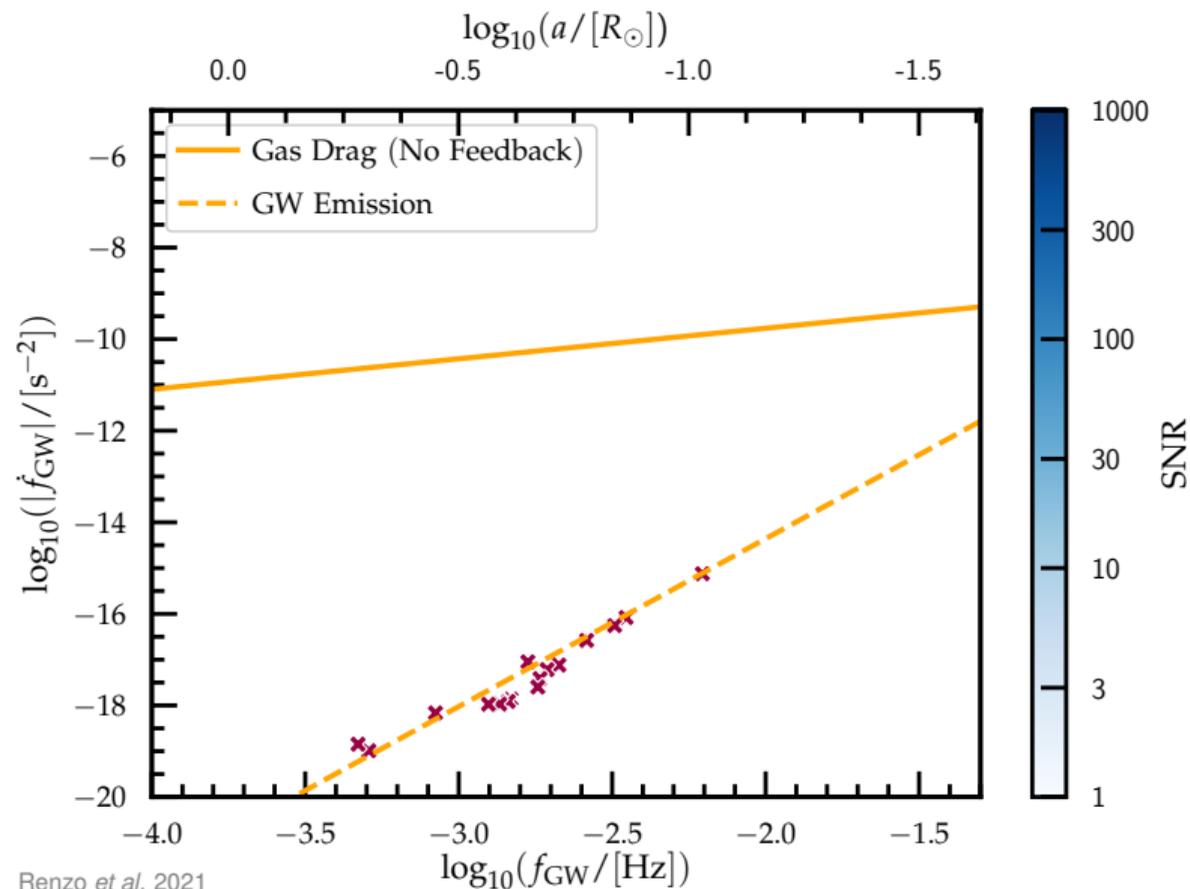
$M_{\text{core}} = 0.5 M_{\odot}, M_2 = 0.3 M_{\odot},$   
 $D = 3 \text{ kpc}, T = 5 \text{ years},$   
averaged over  
orientation and sky location

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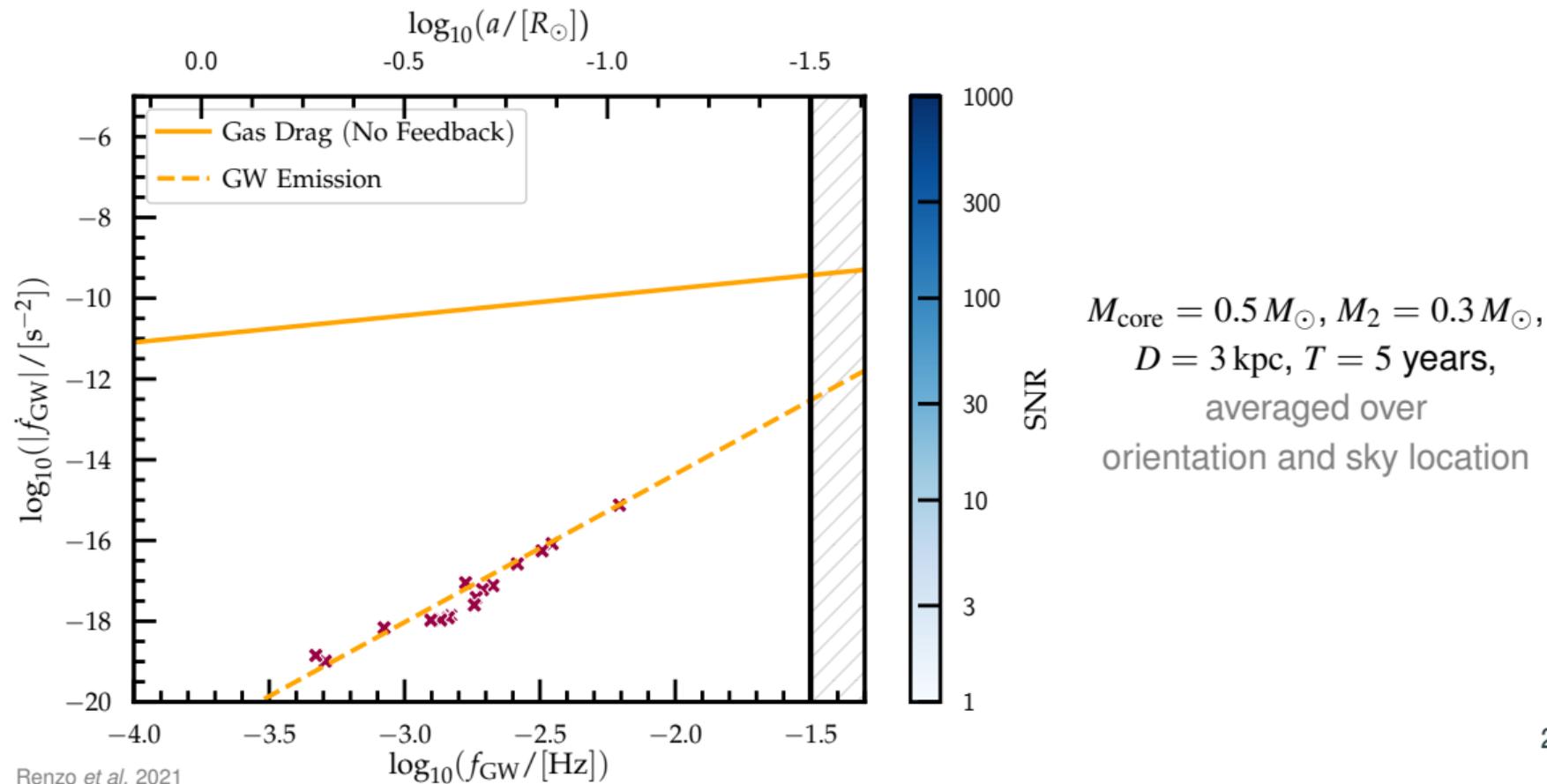
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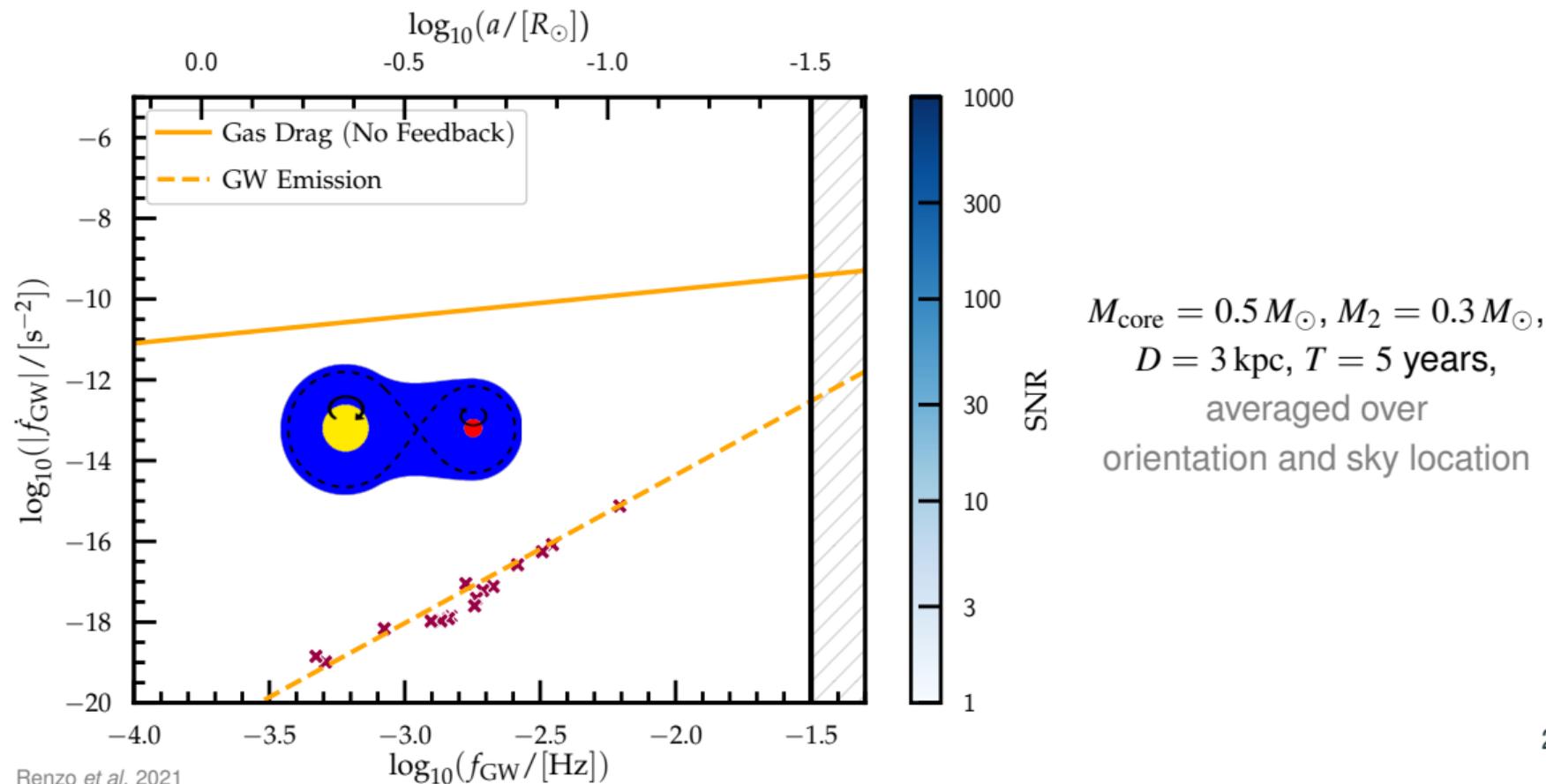


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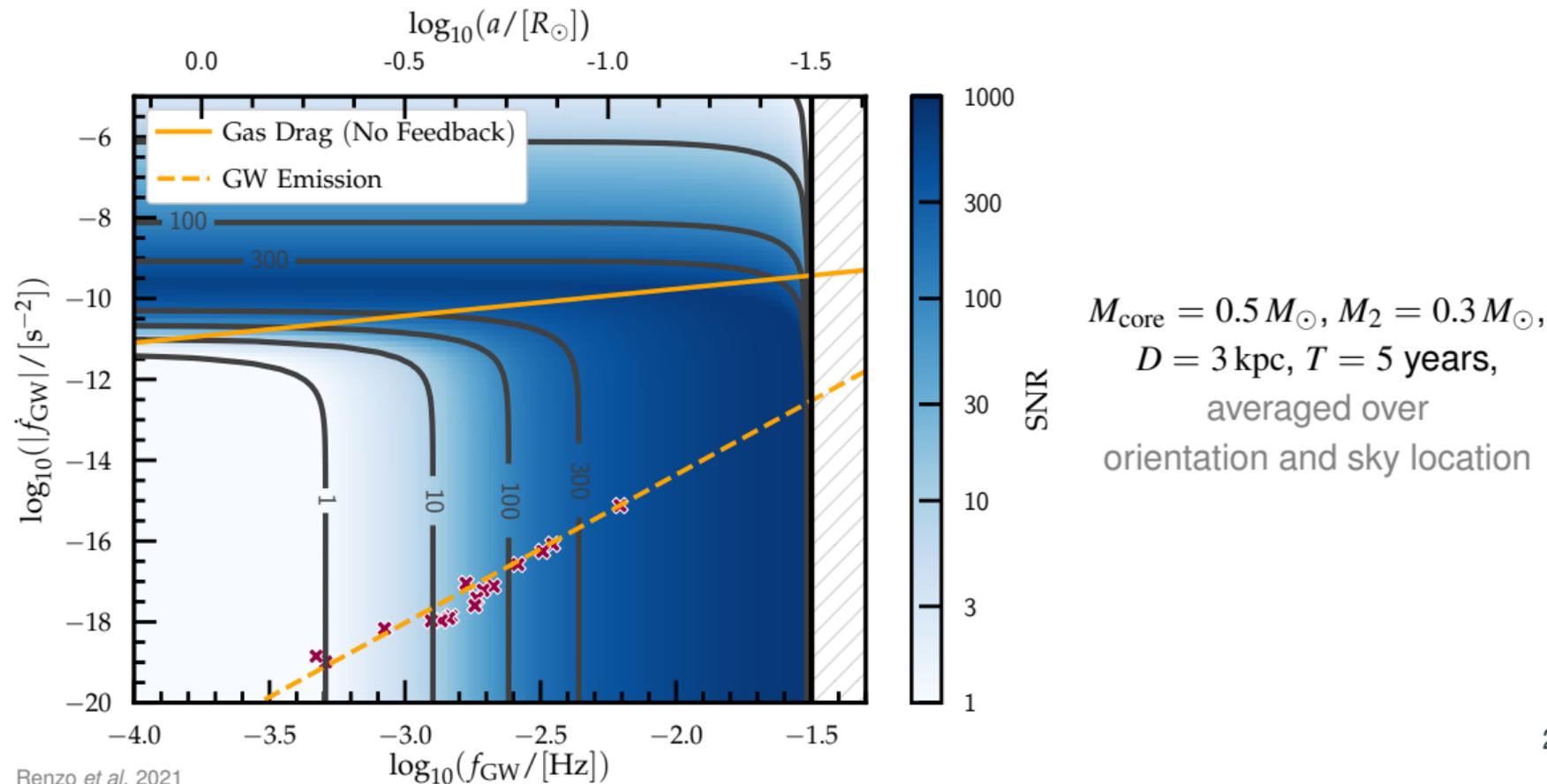
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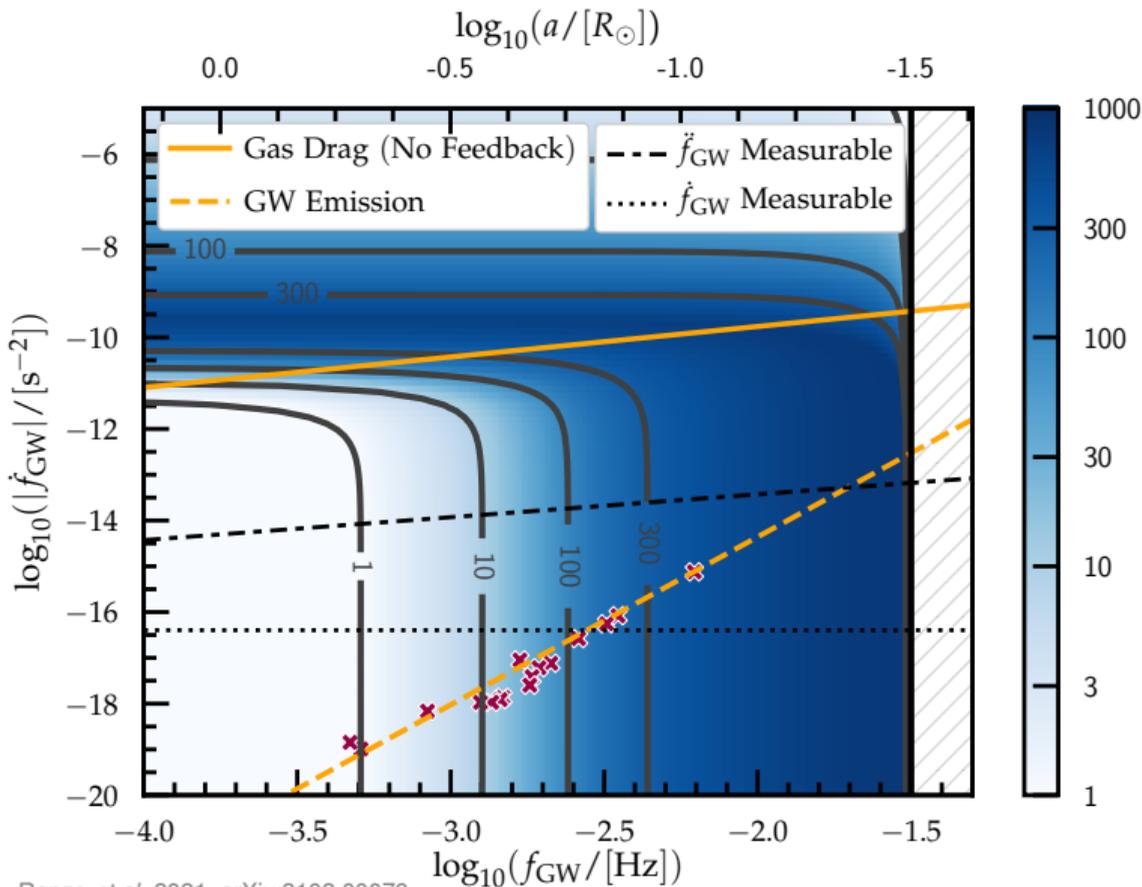


## **Conclusions**

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**Unstable mass transfer**

# Can LISA see common-envelope events? **Maybe!**



- $\sim$  One CE-begin per 10 yr
- $0 \lesssim N_{\text{CE}} \lesssim 1000$
- if stalls at short separation they might be detectable



Direct window on the inside

**If non-detection**

- stalls at large separation
- and/or
- stalling phase is short

## Conclusions

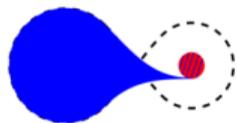
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for real!

## Mass transfer makes binaries different than single stars

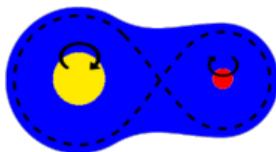
### dynamically stable:

common, widens the orbit, changes *significantly* both stars



### dynamically unstable:

rare, shrinks the orbit, necessary for *most* compact binary formations



**Backup slides**

# Summary of ejection mechanisms

## Binary SN disruption

- Ejects initially less massive star
- Requires SN kick
- Final  $v \simeq v_2^{\text{orb}}$
- Most binaries are disrupted
- Leaves **binary signature**

fast rotation, He/N enrichment, lower apparent age

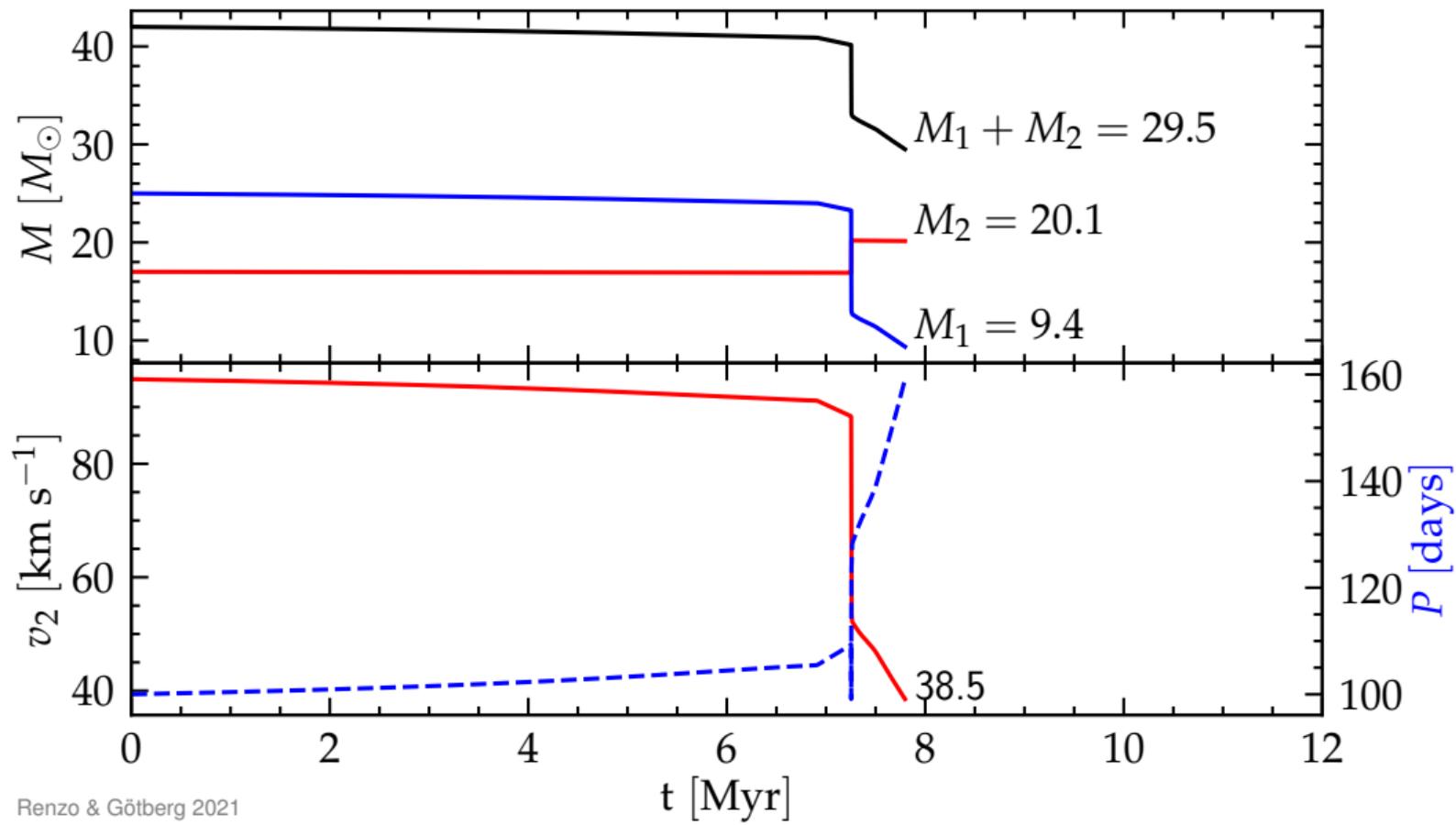
## Cluster ejections

- Happen early on, before SNe
- Can produce faster stars
- Least massive thrown out
- *Gaia* hint: high efficiency dynamical ejection

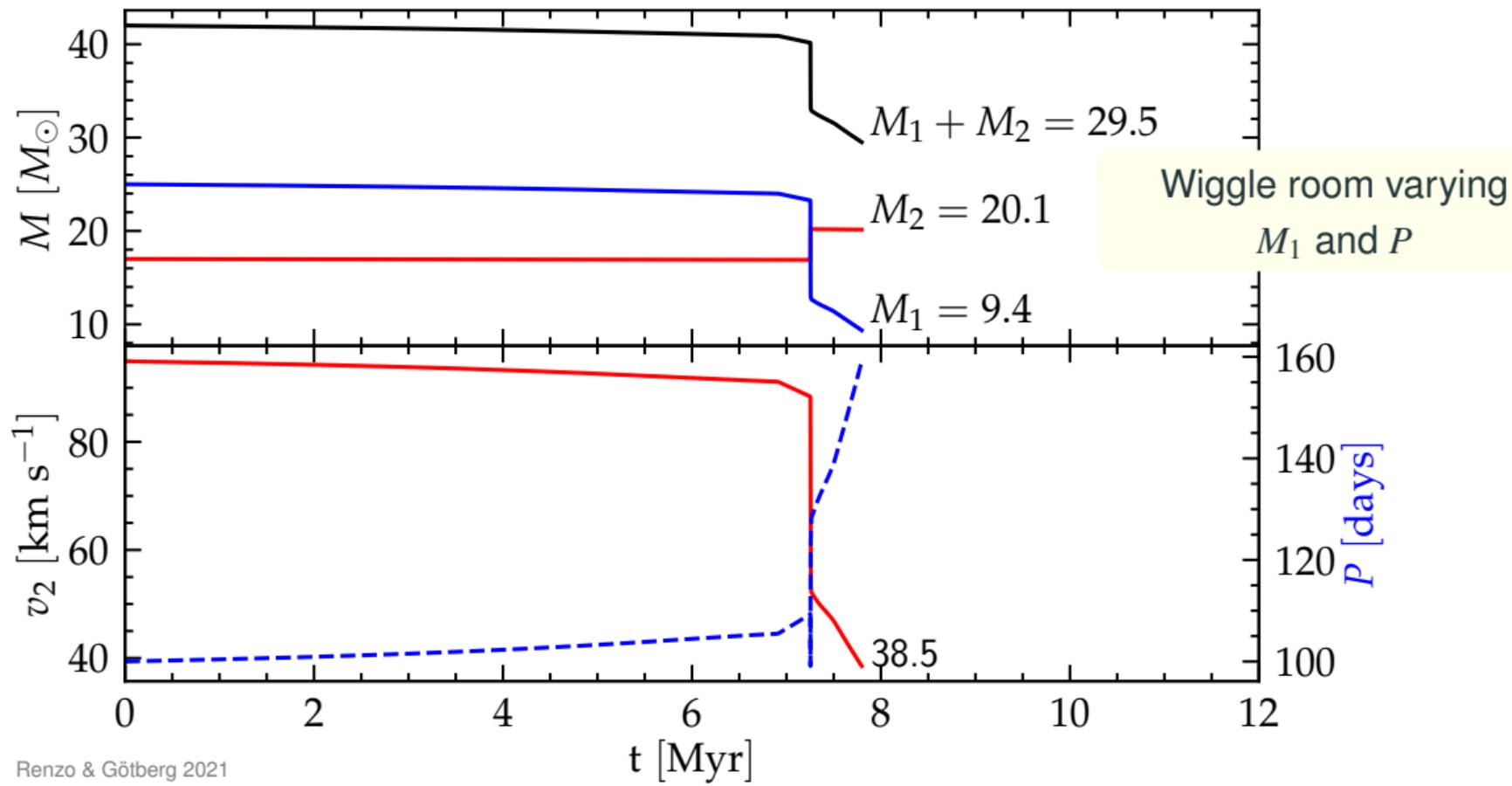
...Binaries are still important! but might not leave signature



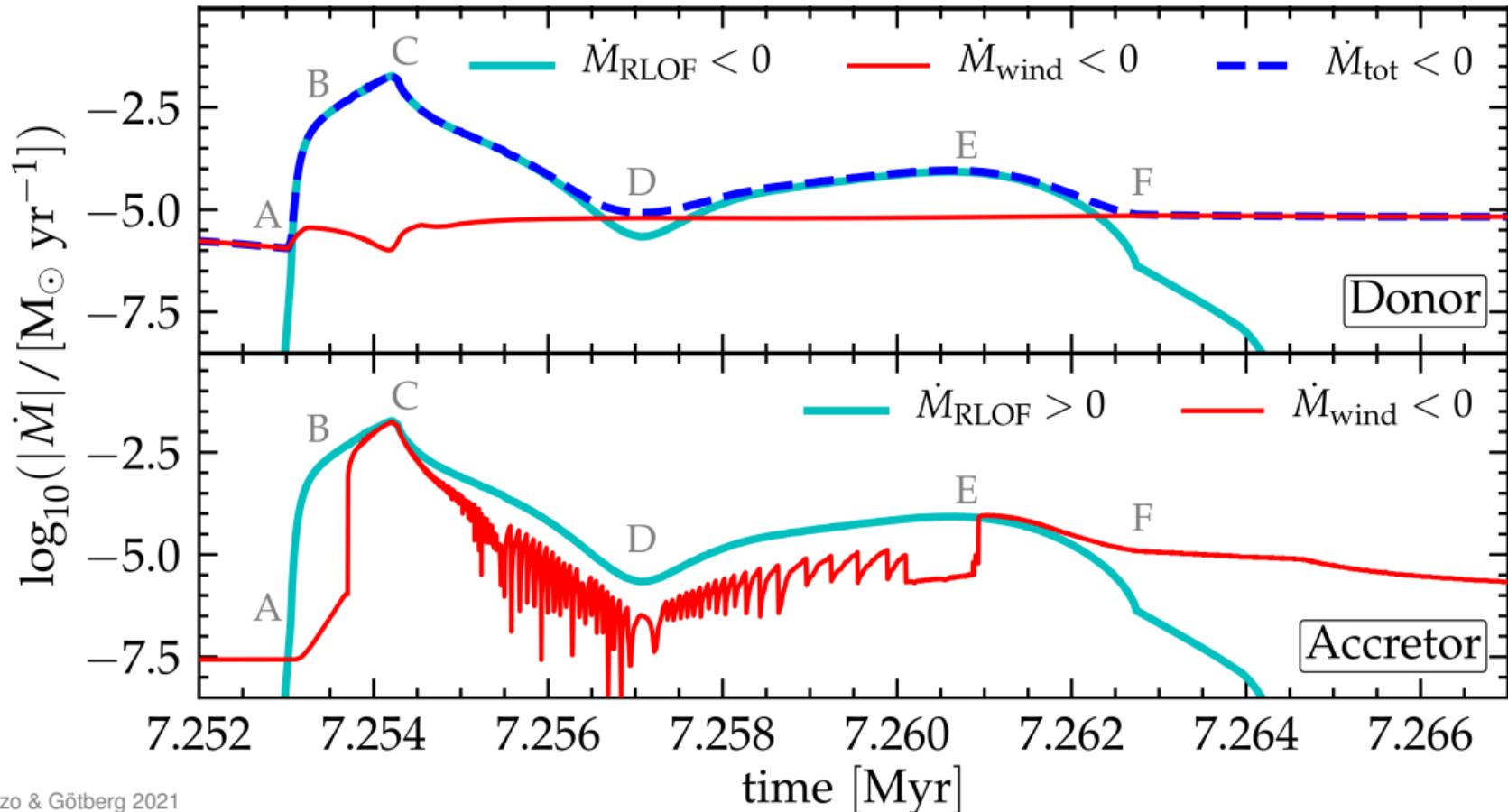
## Spatial velocity & mass



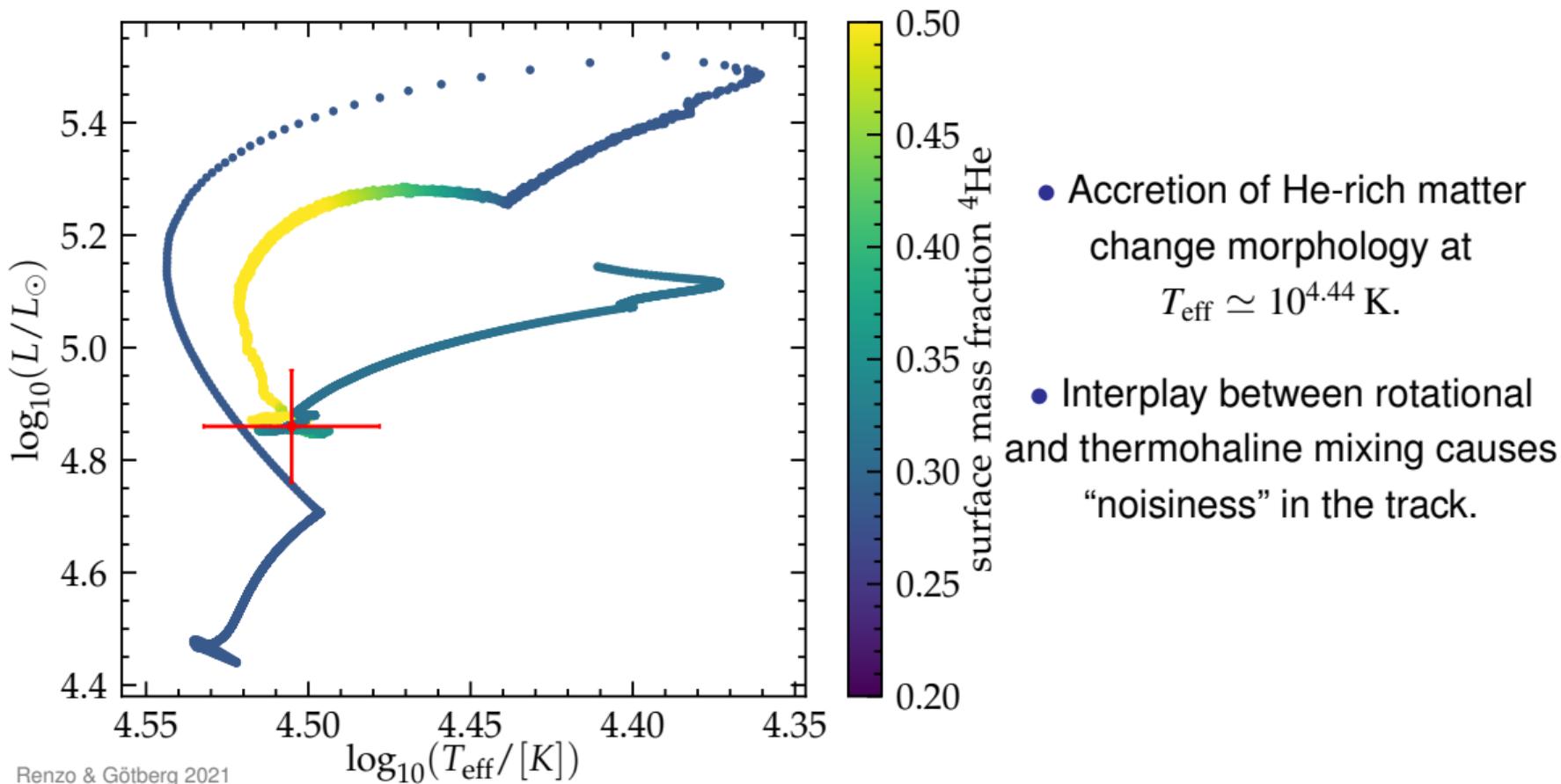
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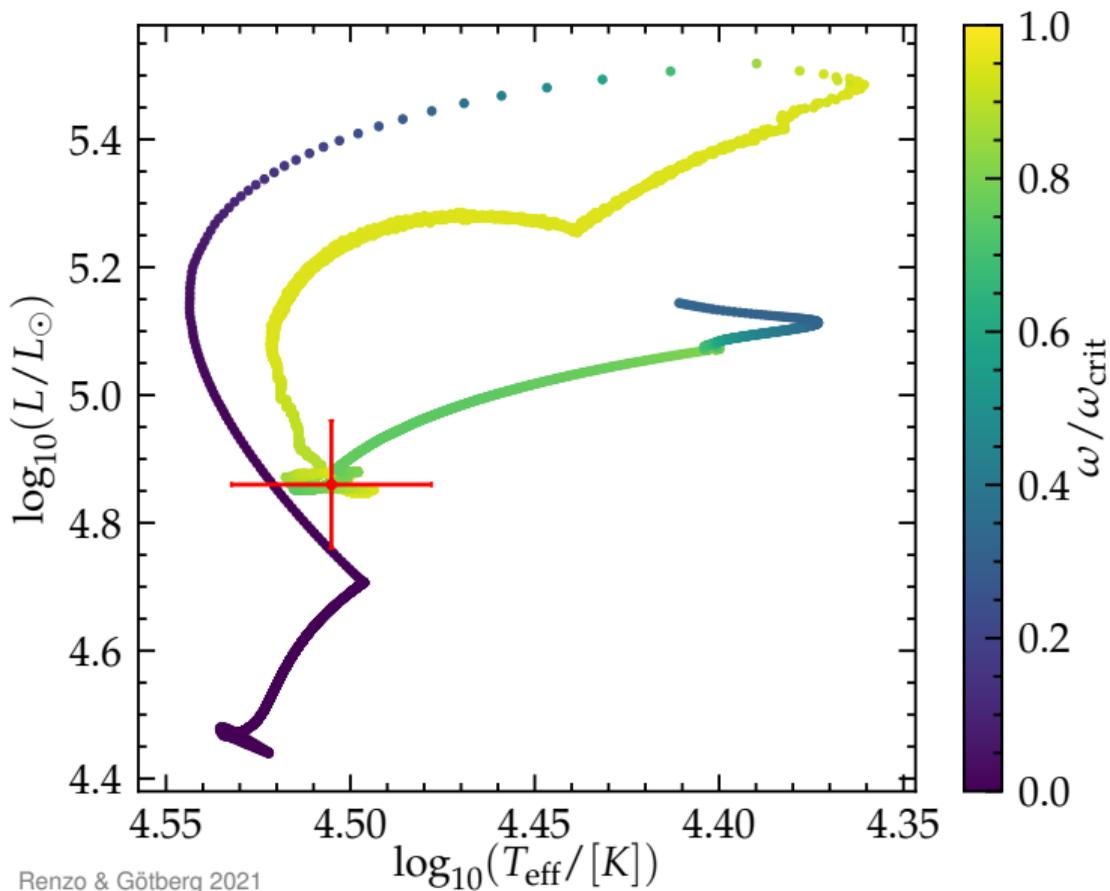
## Mass transfer history: $\Delta t_{\text{RLOF}} \simeq 2 \times 10^4$ years



## Hertzsprung-Russell diagram: helium surface abundance

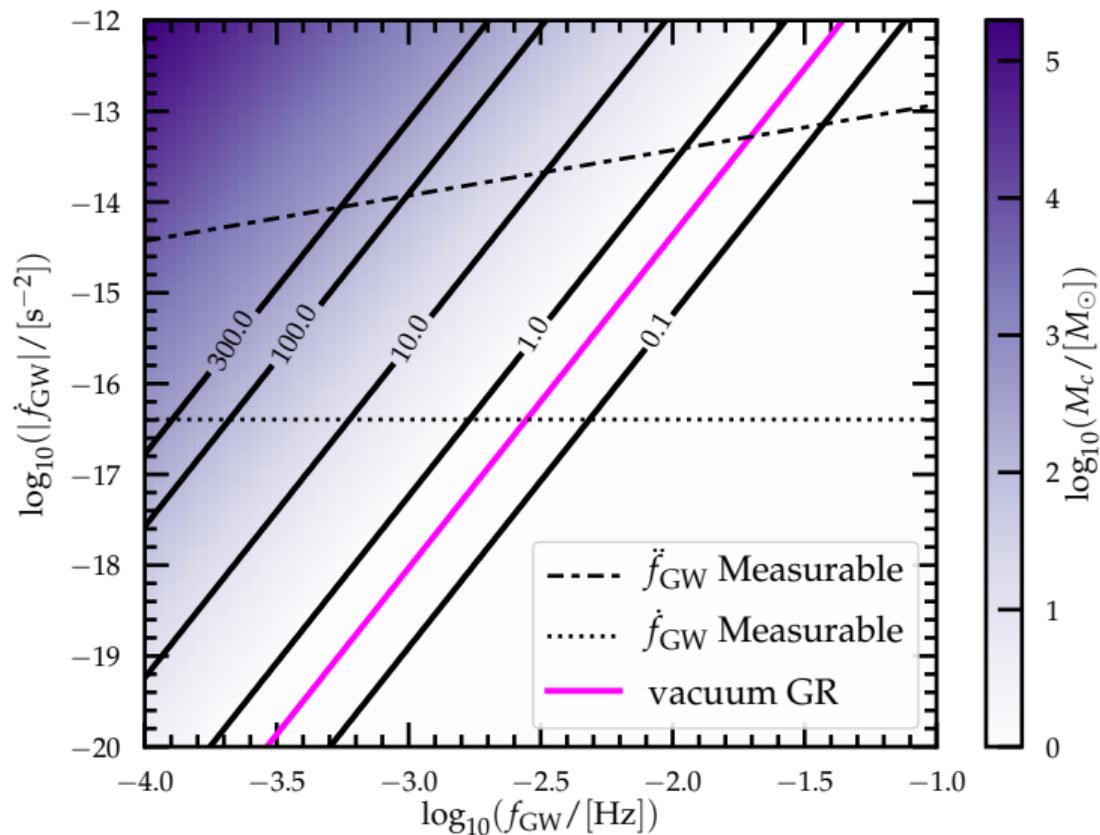


## Hertzsprung-Russell diagram: accretor rotation



- Minimum  $T_{\text{eff}}$  during RLOF reached at onset of critical rotation.
- Rotation close to critical for large part of the main sequence.

## “Stealth bias” assuming GR in vacuum: chirp mass



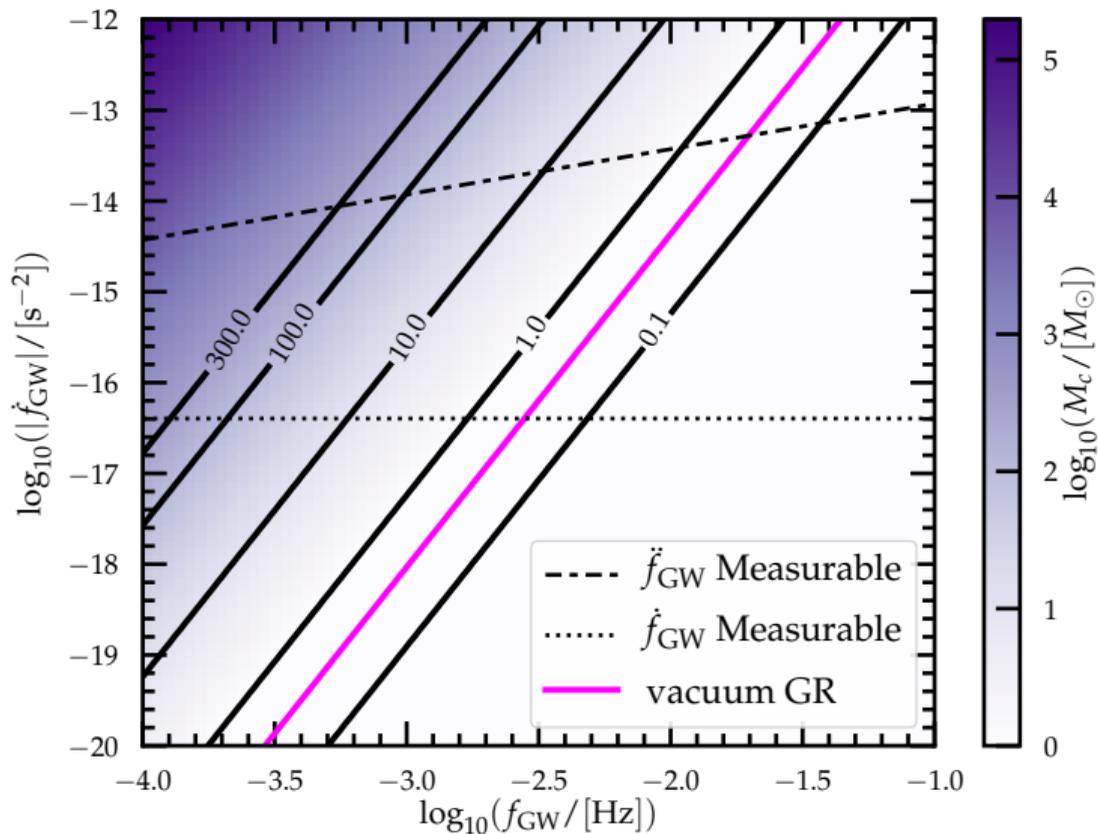
# “Stealth bias” assuming GR in vacuum: chirp mass

“Braking index”

$$n = \frac{f\ddot{f}}{\dot{f}^2}$$

↓

$$n_{\text{GR}} = \frac{11}{3}$$



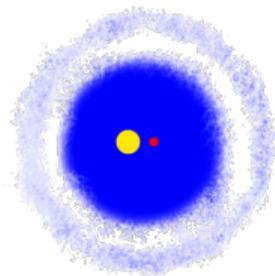
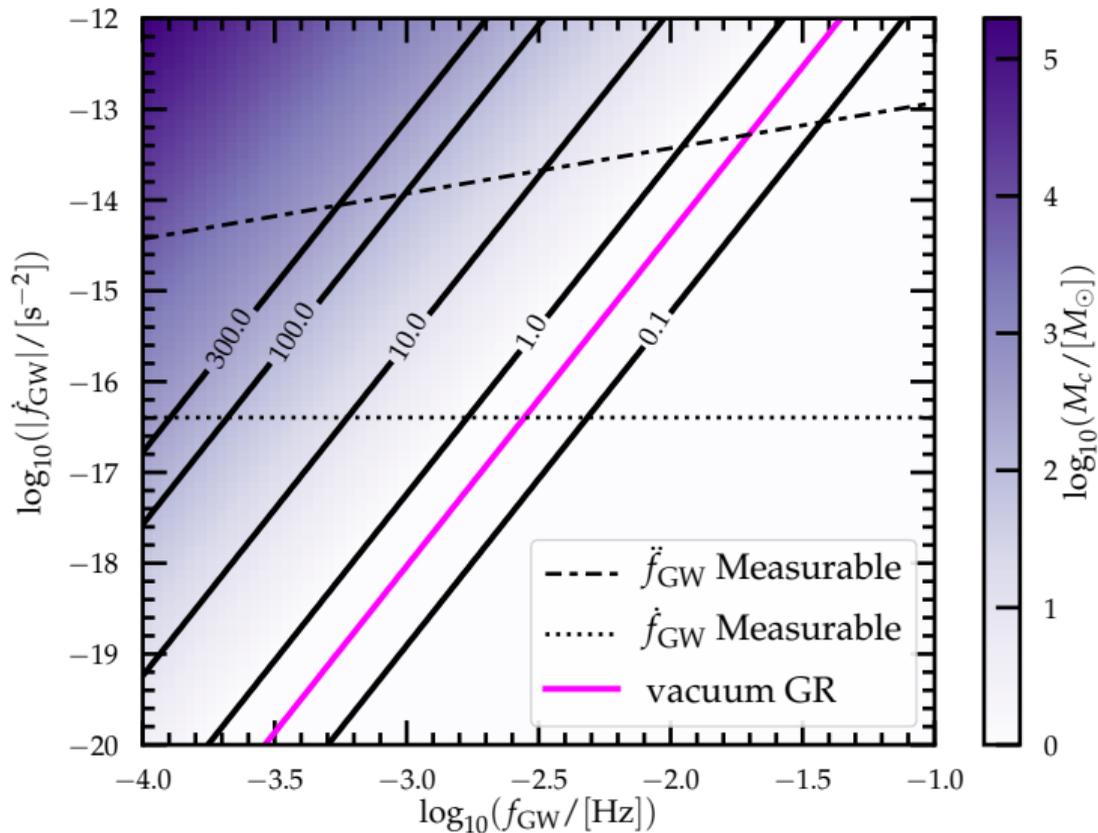
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EM counterparts:

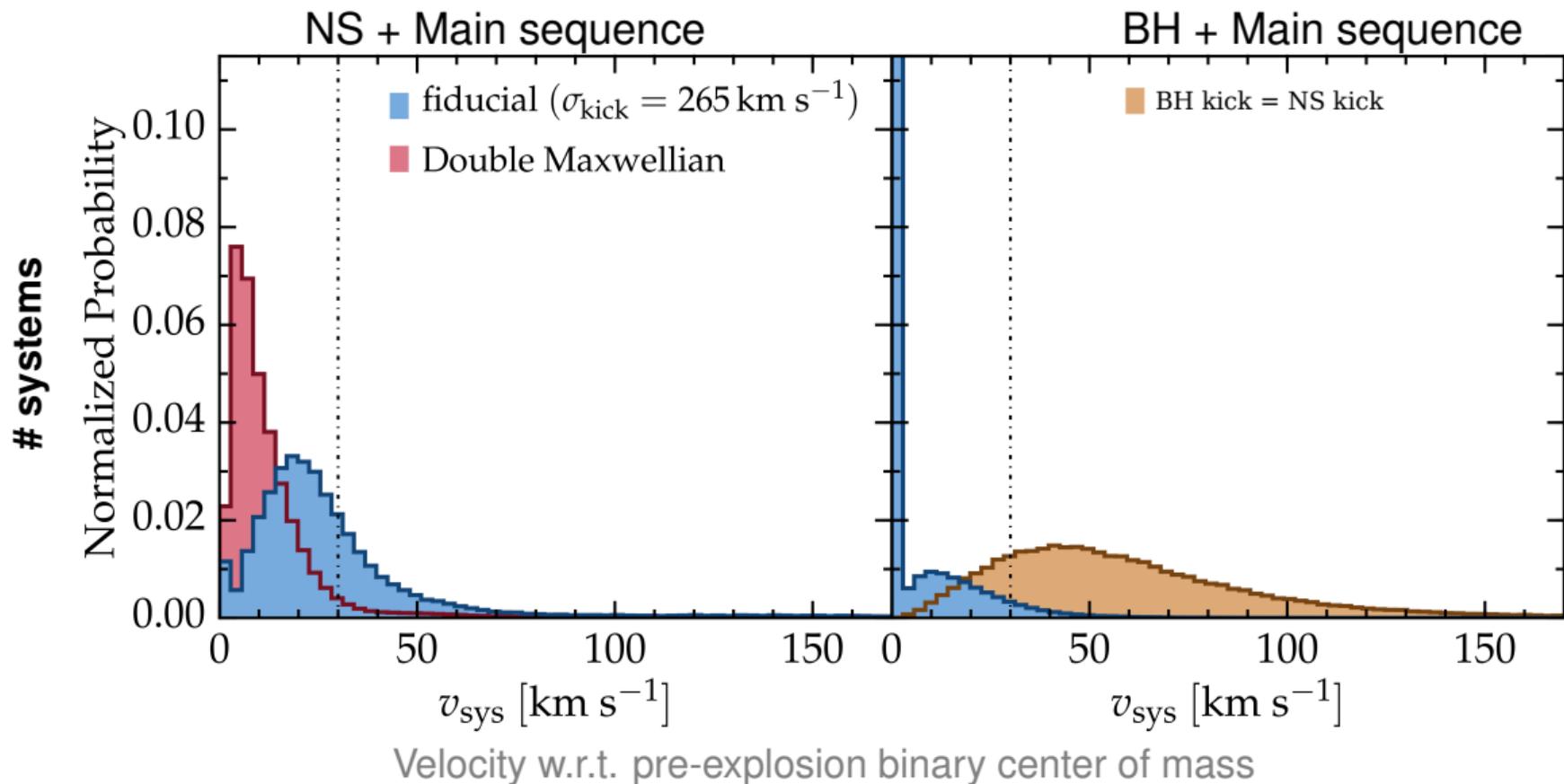
- Optical/IR transients

(Blagorodnova *et al.* 2020)

- “weird” red giant star

(Clayton *et al.* 2017)

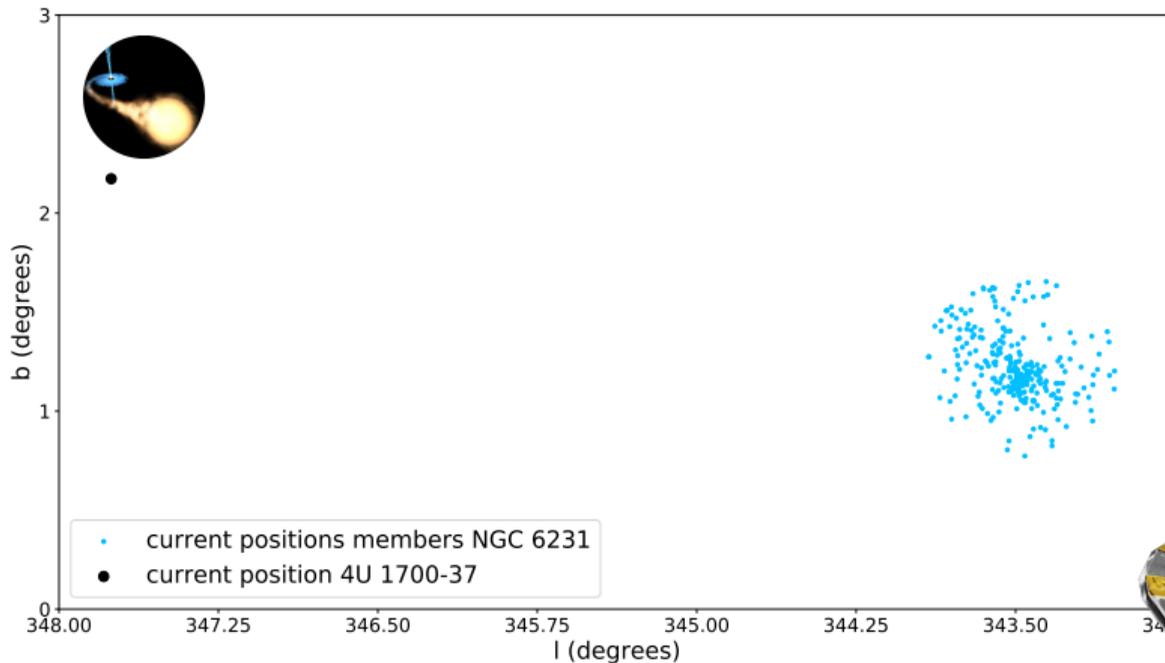
## Post-SN velocity of surviving binaries



## Preliminary: The case of 4U1700-37

$$M \simeq 2.5 M_{\odot}, M_{*} \simeq 60 \pm 10 M_{\odot}, P \simeq 3.4 \text{ days}, e \simeq 0.22, v \simeq 60 \text{ km s}^{-1}$$

Galactic longitude



Galactic latitude

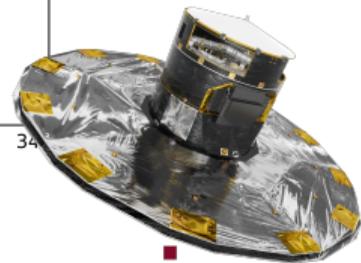
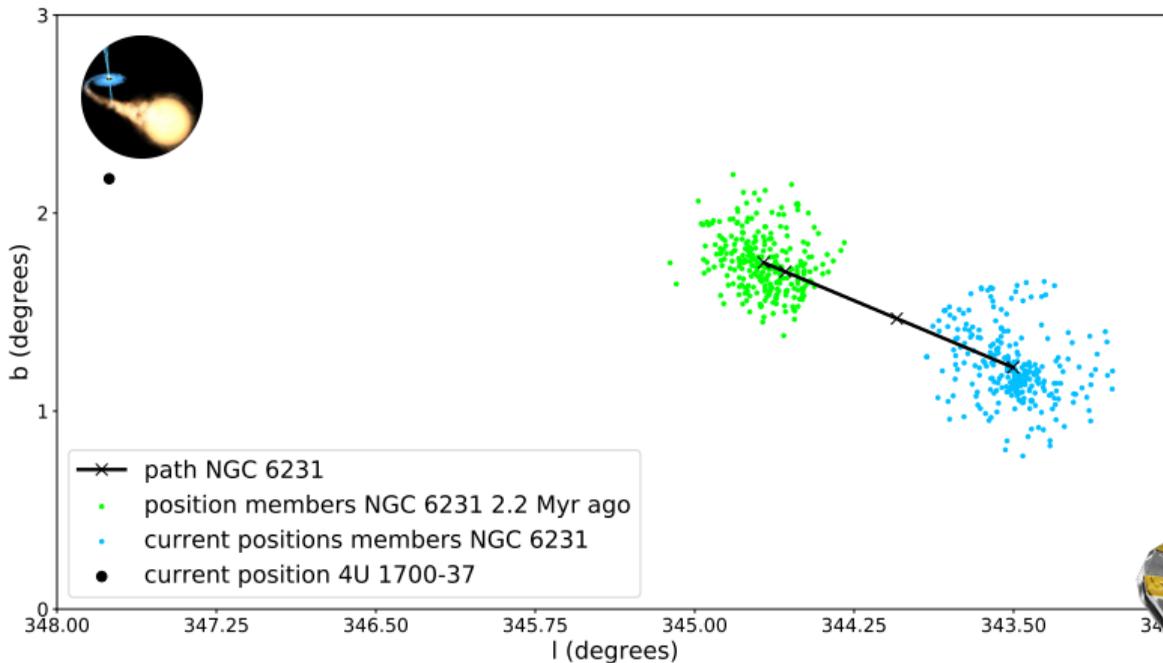


gaia

# Preliminary: *Gaia* corroborates cluster of origin

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Galactic longitude



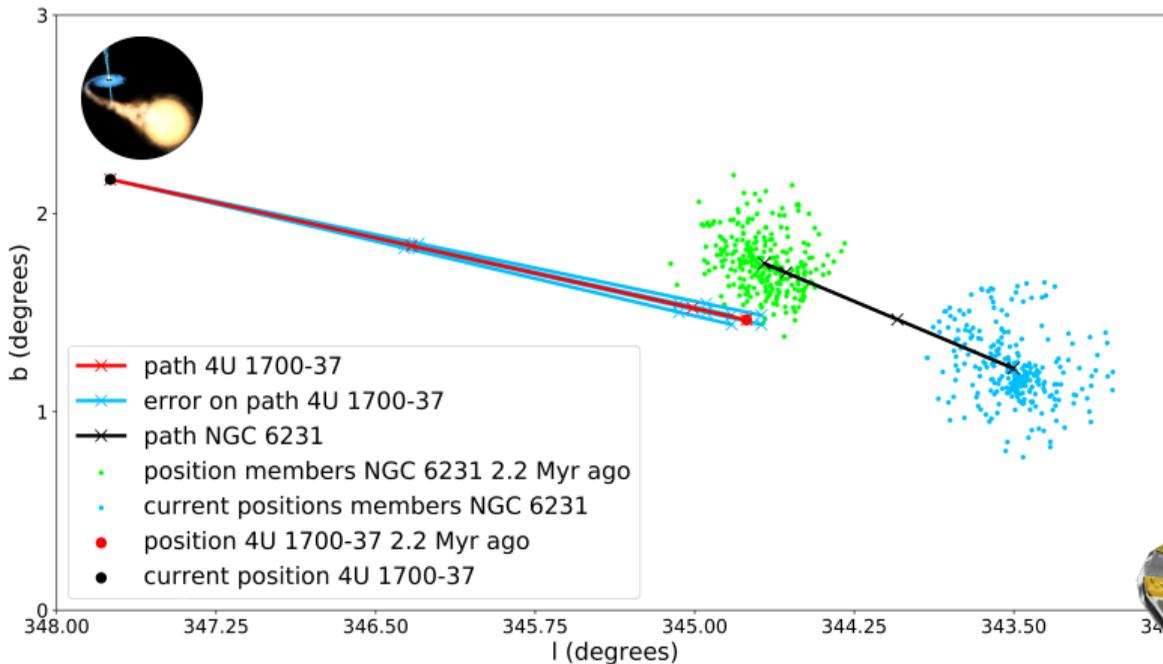
gaia

Galactic latitude

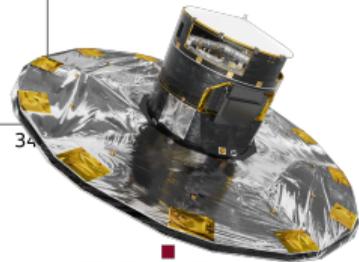
# Preliminary: Cluster of origin constrains past evolution

$$M \simeq 2.5 M_{\odot}, M_* \simeq 60 \pm 10 M_{\odot}, P \simeq 3.4 \text{ days}, e \simeq 0.22, v \simeq 60 \text{ km s}^{-1}$$

Galactic longitude



Galactic latitude



gaia

# Period evolution depends on uncertain free parameters

