

Explosions in massive binaries:

“widowed” stars and consequences for GW astronomy



Mathieu Renzo
PhD in Amsterdam

S. E. de Mink, E. Zapartas, Y. Götberg, E. Laplace,
Collaborators: R. J. Farmer, S. Toonen, S. Justham, R. G. Izzard,
P. Marchant, D. J. Lennon, H. Sana, S. N. Shore, ...

Why are massive stars important?

Nucleosynthesis &
Chemical Evolution

Ionizing Radiation

Star Formation

Supernovae

GW Astronomy



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Nucleosynthesis &
Chemical Evolution

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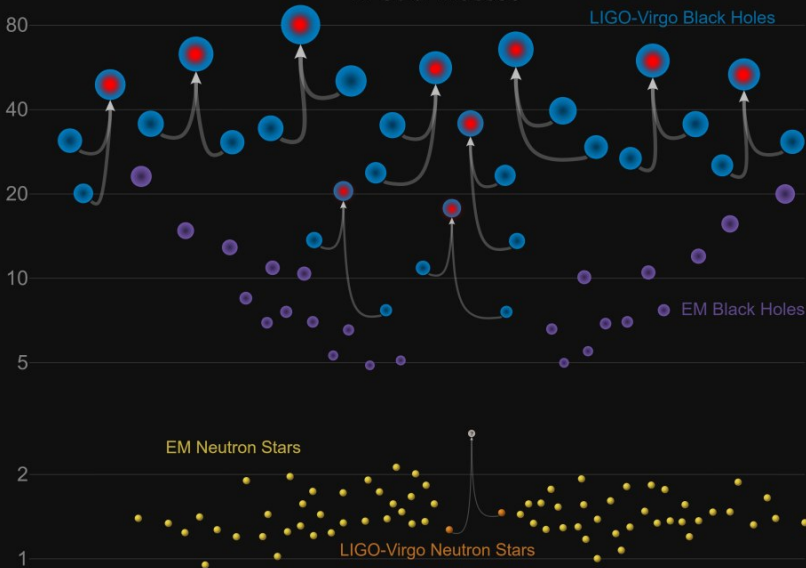
GW Astronomy

**~70% of O type stars will
interact with a companion**

(e.g., Mason *et al.* '09, Sana & Evans '11,
Sana *et al.* '12, Kiminki & Kobulnicky '12,
Kobulnicky *et al.* '14, Almeida *et al.* '17)

Masses in the Stellar Graveyard

in Solar Masses



BH or NS?

- Single stars winds impact on the core structure

Keep the stars together

- The most common evolution for massive binaries
- Constraints on BH kicks using runaway “widow”

The most massive (stellar) BHs

- (Pulsational) pair instability
- The BH mass distribution
 - Induced eccentricity

Conclusions

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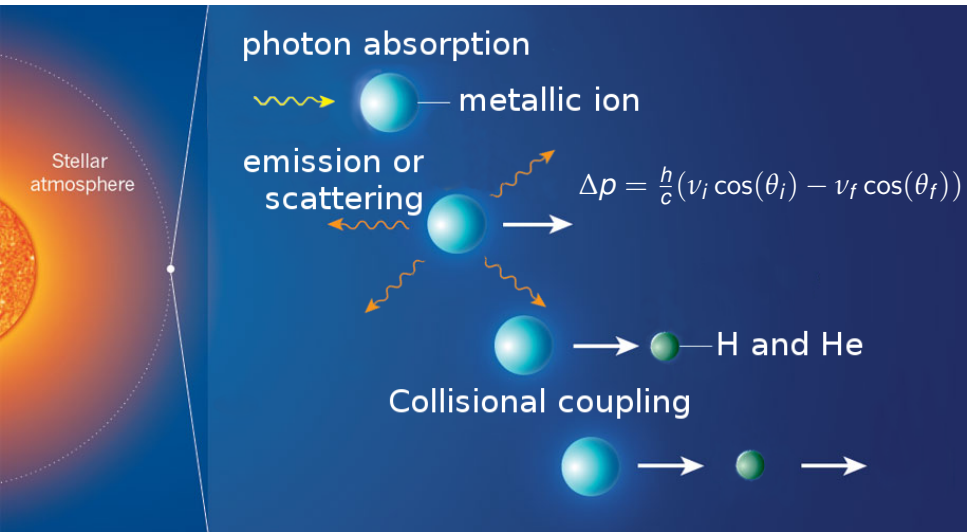
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Problems: High Non-Linearity and Clumpiness

Inhomogeneities:

$$f_{\text{cl}} \stackrel{\text{def}}{=} \frac{\langle \rho^2 \rangle}{\langle \rho \rangle^2} \neq 1 \Rightarrow \dot{M} \neq 4\pi r^2 \rho v(r)$$

Risk:

Possible overestimation of the
wind mass loss rate

Inhomogeneities:

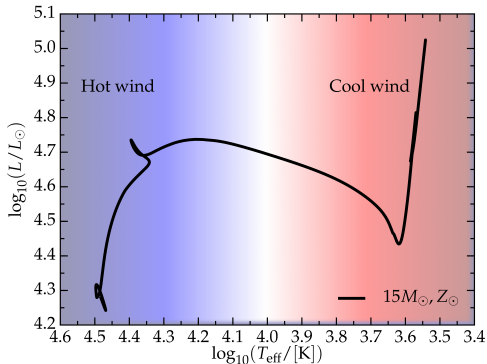
$$f_{\text{cl}} \stackrel{\text{def}}{=} \frac{\langle \rho^2 \rangle}{\langle \rho \rangle^2} \neq 1 \Rightarrow \dot{M} = \eta 4\pi r^2 \rho v(r)$$

Aim: quantify systematic uncertainty

Grid of Z_{\odot} non-rotating models:

$$M_{\text{ZAMS}} = 15, 20, 25, 30, 35 M_{\odot}$$

$$\eta \equiv \sqrt{f_{\text{cl}}} = 1, \frac{1}{3}, \frac{1}{10}$$

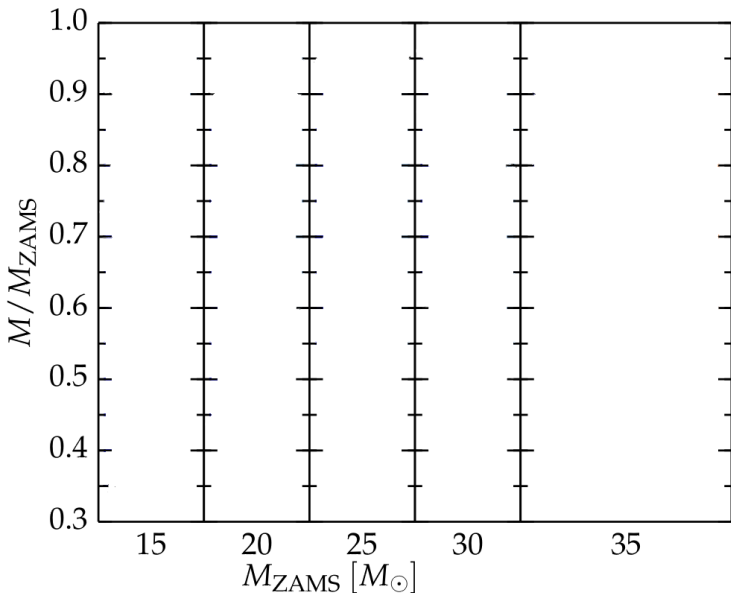


Combinations of wind mass loss rates for “hot” ($T_{\text{eff}} \geq 15$ [kK]), “cool” ($T_{\text{eff}} < 15$ [kK]) and WR (H deficient):

Kudritzki *et al.* '89; Vink *et al.* '00, '01;

Van Loon *et al.* '05; Nieuwenhuijzen *et al.* '90; De Jager *et al.* '88;

Nugis & Lamers '00; Hamann *et al.* '98.

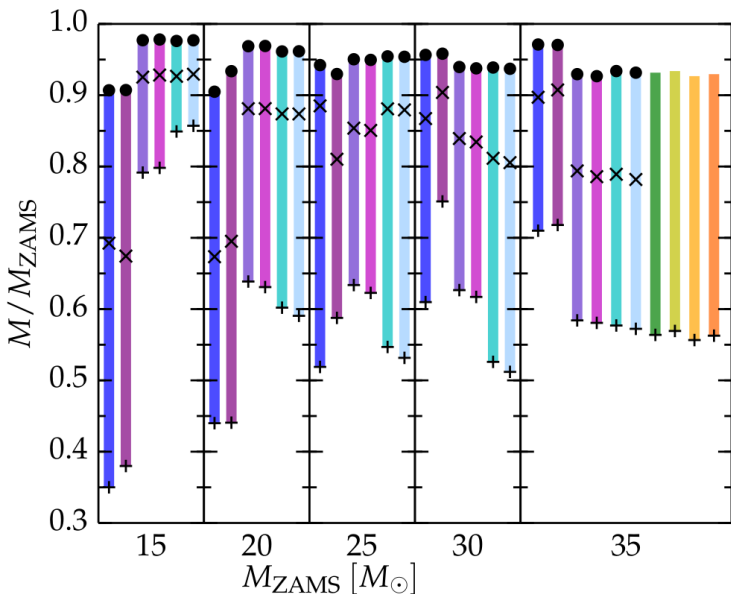


MESA

Legend:

- $\eta = 0.1$
- x $\eta = 0.33$
- + $\eta = 1.0$

$\eta \rightarrow$ largest
uncertainty

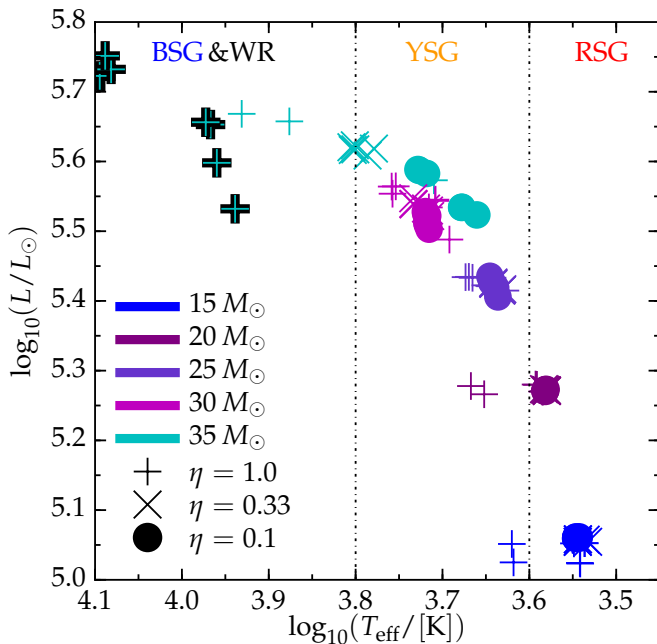


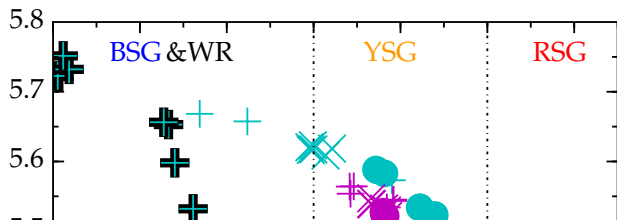
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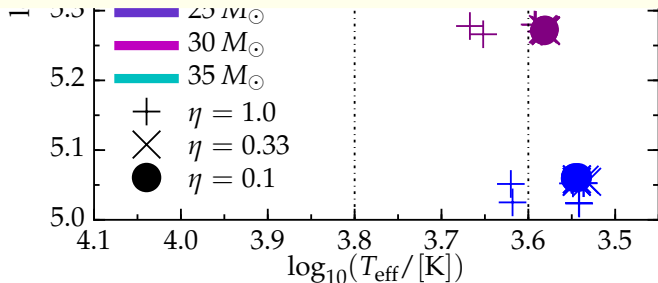
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Same appearance with ΔM up to 50%

⇒ The internal structure re-adjusts to the wind mass loss



$$\tilde{\zeta}_{\mathcal{M}}(t) \stackrel{\text{def}}{=} \frac{\mathcal{M}/M_{\odot}}{R(\mathcal{M})/1000 \text{ km}}$$

Single parameter to describe the core structure

e.g., O'Connor & Ott '11,

Ugliano *et al.* '12,

Sukhbold & Woosley '14,

but see (for 3D explosions):

Ott *et al.* '18,

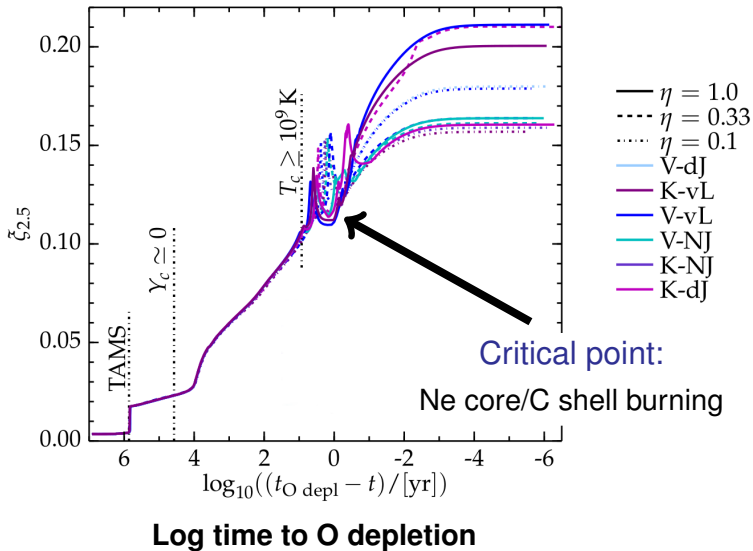
Kuroda *et al.* '18


$$\mathcal{M} = 2.5 M_{\odot}$$

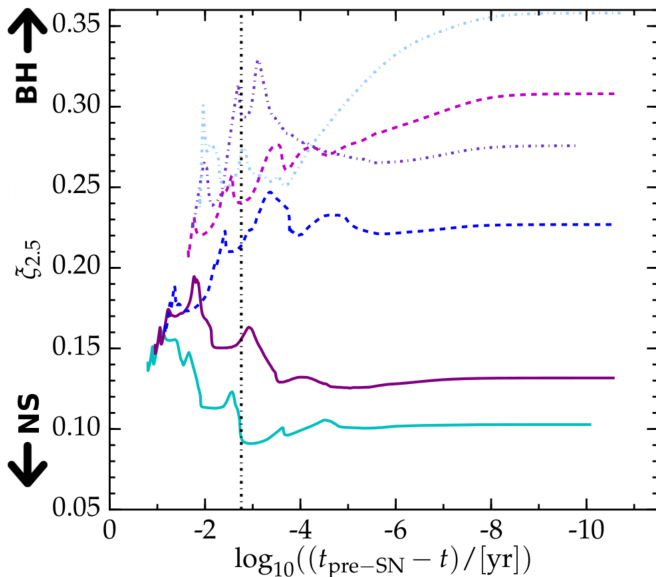
not to scale!

$R(\mathcal{M})$

$M_{ZAMS} = 25 M_{\odot}$ MESA models



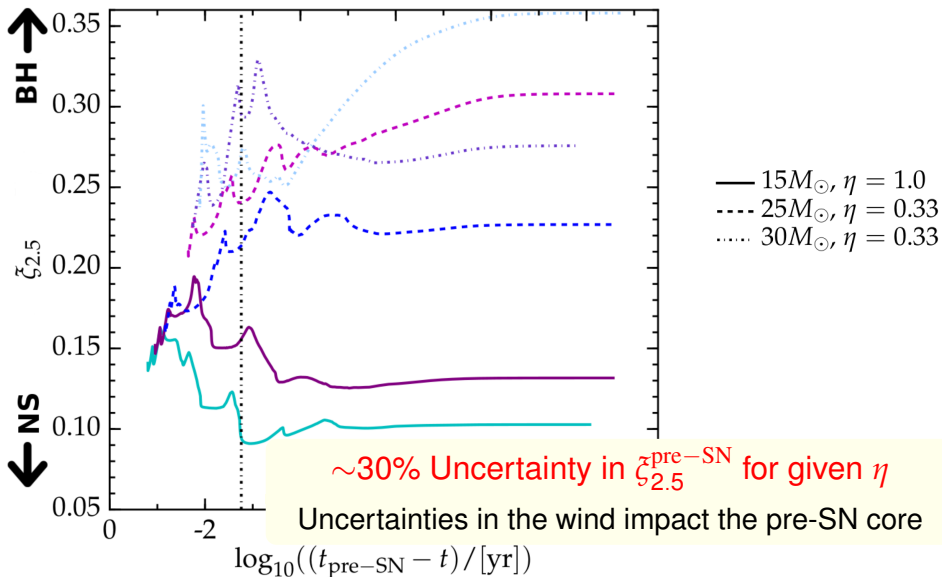
Si shell burning →



- $15M_{\odot}, \eta = 1.0$
- - - $25M_{\odot}, \eta = 0.33$
- ⋯ $30M_{\odot}, \eta = 0.33$

Log time to core-collapse

Si shell burning →



Log time to core-collapse

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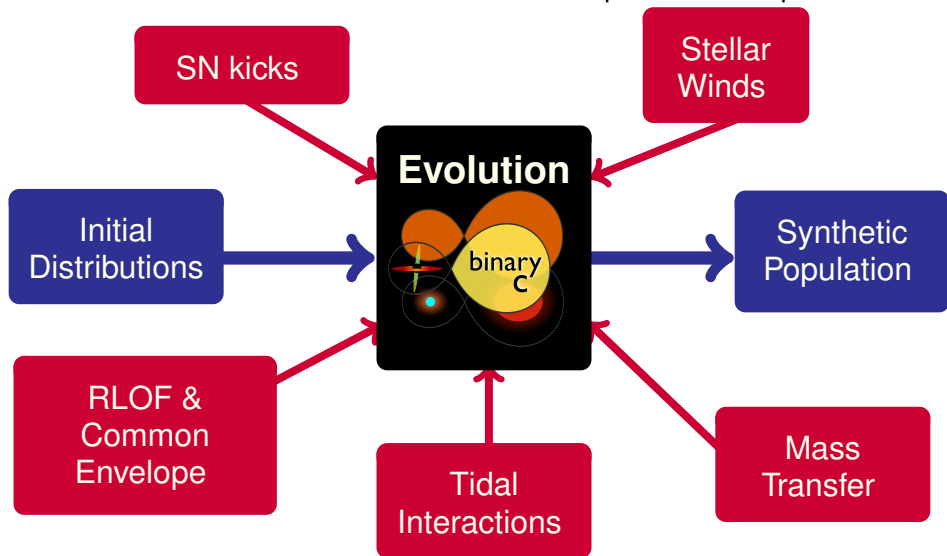
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Methods: Population Synthesis



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Fast \Rightarrow Allows statistical tests of the inputs & assumptions



Binary disruption



Credits: ESO, L. Calçada, M. Kornmesser, S.E. de Mink

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The binary disruption shoots out the accretor

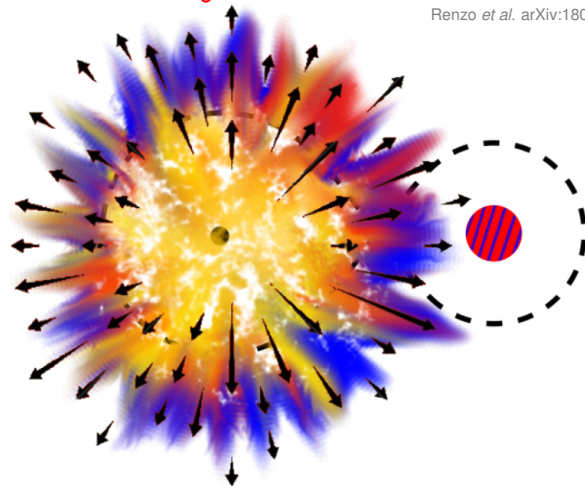
Spin up: Packet '81, Cantiello *et al.* '07, de Mink *et al.* '13

Pollution: Blaauw '93

Rejuvenation: Hellings '83, Schneider *et al.* '16

86_{-9}^{+11} % of massive binaries are disrupted

Renzo *et al.* arXiv:1804.09164, Eldridge *et al.* 11, De Donder *et al.* 97



- Unbinding Matter

(e.g., Blaauw '61)

- Ejecta Impact

(e.g., Wheeler *et al.* '75,
Tauris & Takens '98, Liu *et al.* '15)

- SN Natal Kick

(e.g., Shklovskii '70, Janka '16)

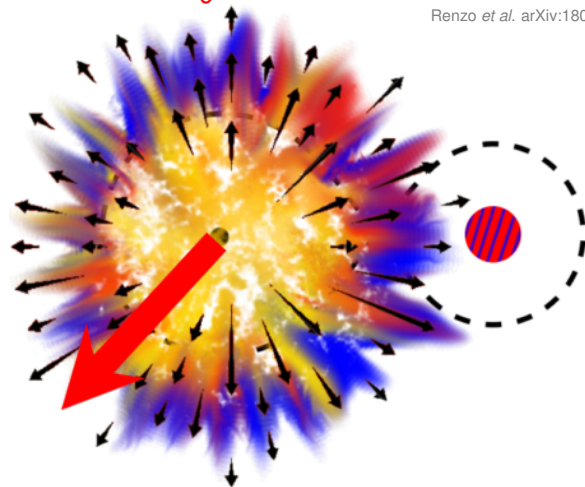
What exactly disrupts the binary?



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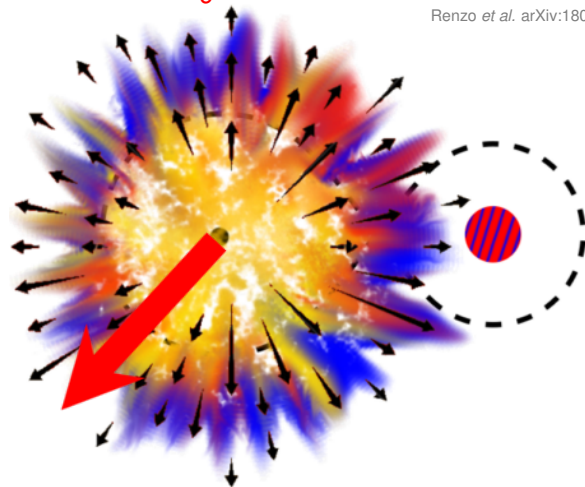
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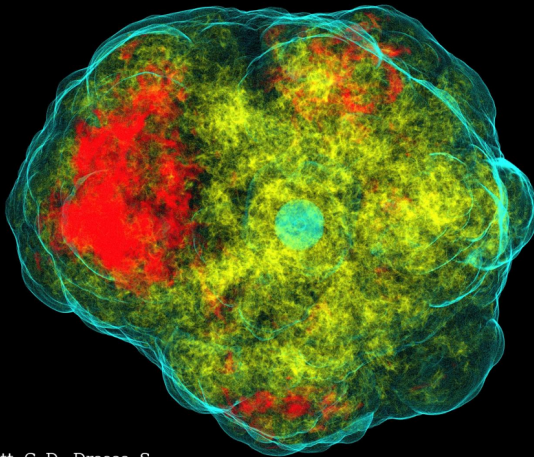
$$v_{\text{dis}} \simeq v_{2,\text{orb}}^{\text{pre-SN}} = \frac{M_1}{M_1 + M_2} \sqrt{\frac{G(M_1 + M_2)}{a}}$$

Most binaries produce a slow “walkaway” star

SN natal kick

Observationally: $v_{\text{pulsar}} \gg v_{\text{OB-stars}}$

Physically: ν emission and/or ejecta anisotropies

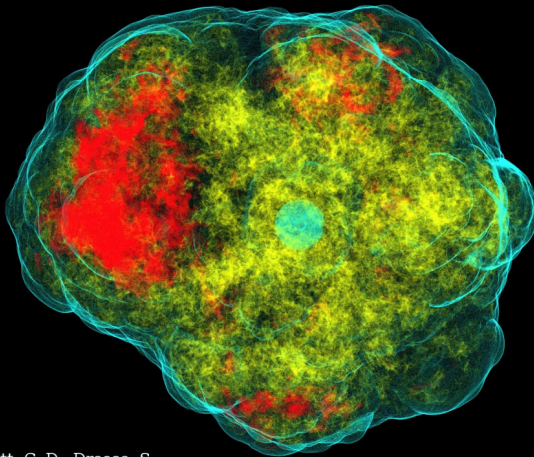


Credits: Ott, C. D., Drasco, S.

SN natal kick

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BH kicks?

BH or NS?

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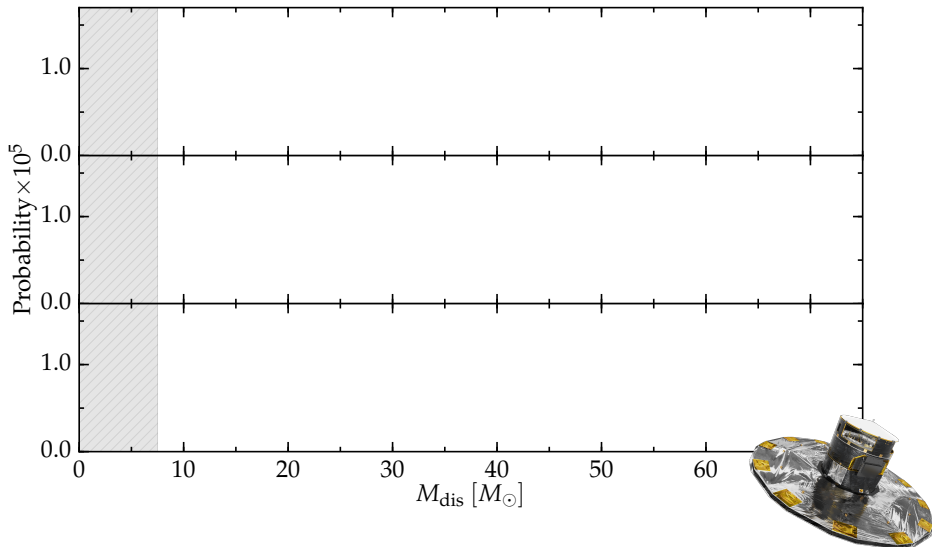
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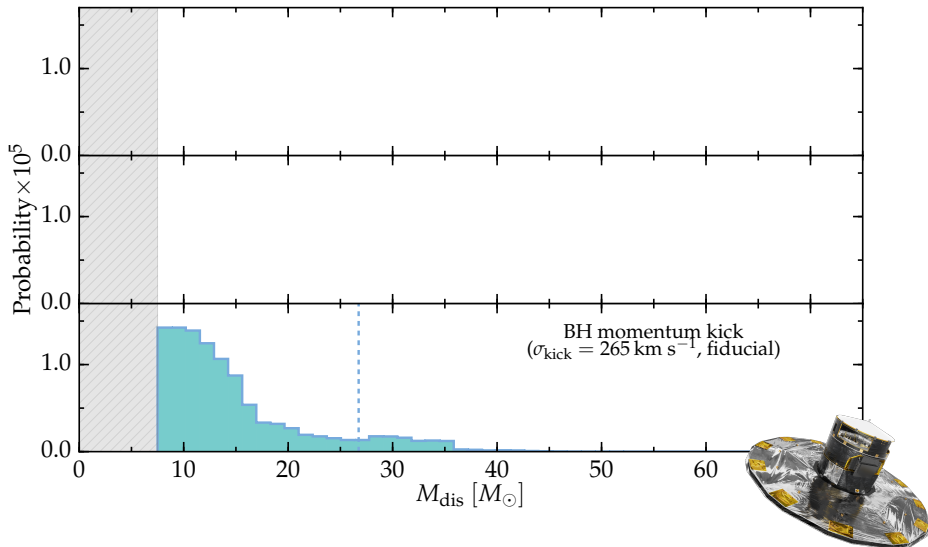
The most massive (stellar) BHs

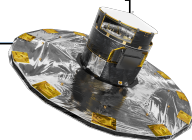
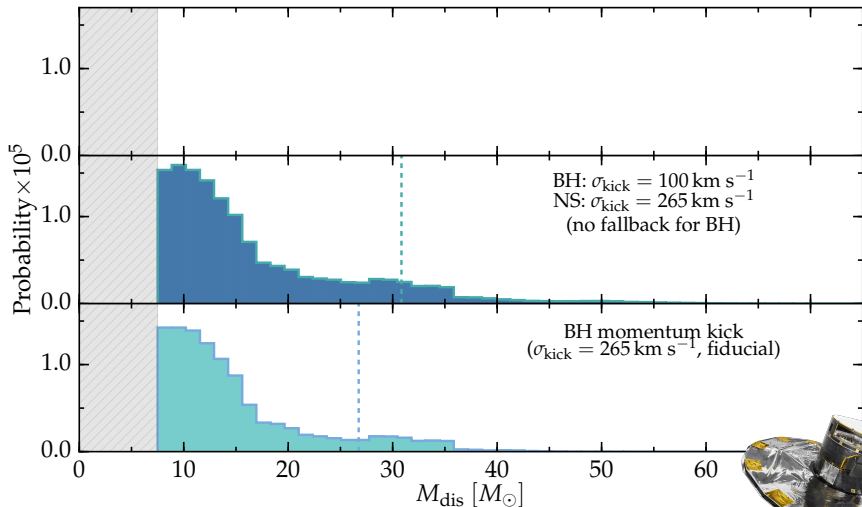
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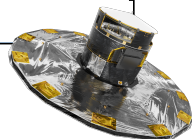
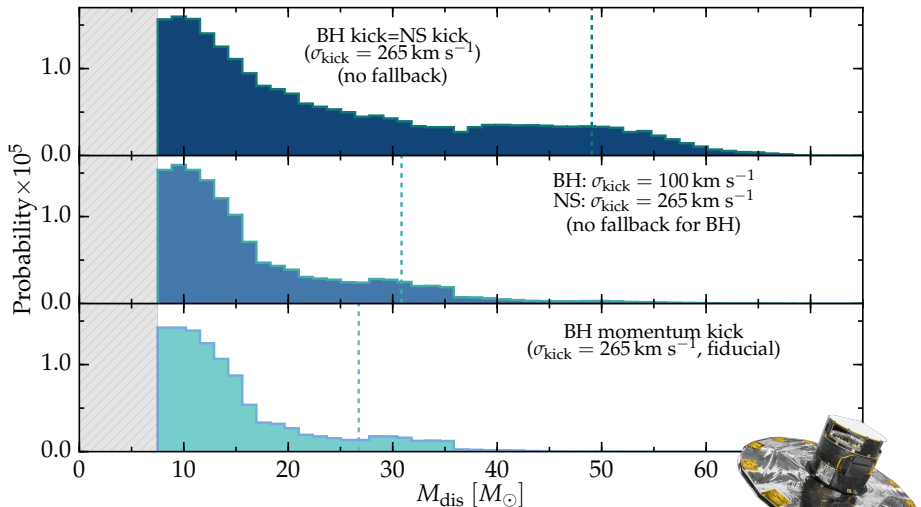
Massive runaways mass function ($v \geq 30 \text{ km s}^{-1}$, $M \geq 7.5 M_{\odot}$)

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Radiation dominated:

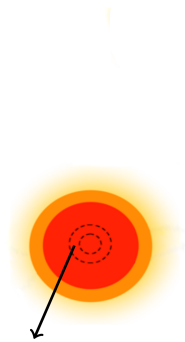
$$P_{\text{tot}} \simeq P_{\text{rad}}$$

$$M_{\text{He}} \gtrsim 32 M_{\odot}$$

Woosley 2017,

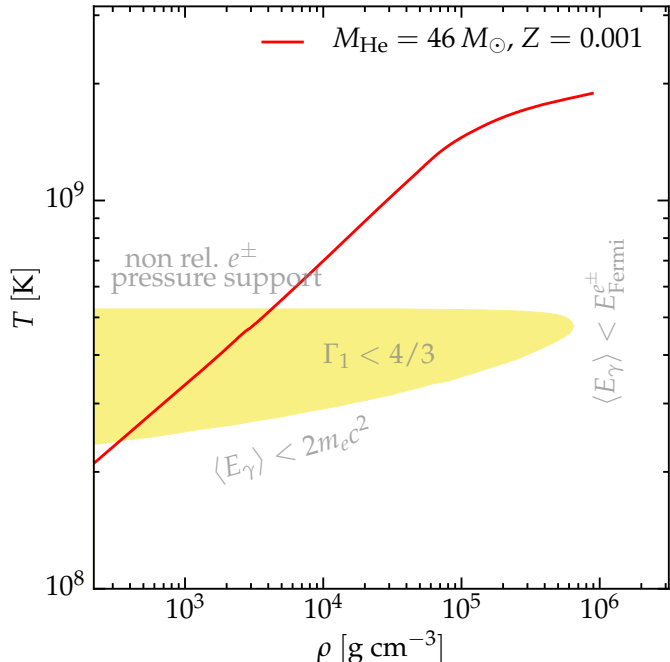
Marchant, Renzo *et al.* arXiv:1810.13412,

Renzo, Farmer *et al.*, to be submitted



1. Pair production
 $\gamma\gamma \rightarrow e^+e^-$

$$\Gamma_1 \stackrel{\text{def}}{=} \left(\frac{\partial \ln P}{\partial \ln \rho} \right)_s$$



He core computed with **MESA**

2. Softening of EOS
triggers collapse

$$\Gamma_1 < \frac{4}{3}$$



1. Pair production



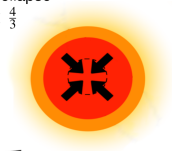
Thermal timescale

$$\tau \propto \frac{GM_{\text{He}}^2}{RL_\nu}, \quad L_\nu \gg L$$

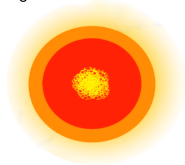
(Fraleigh 68)

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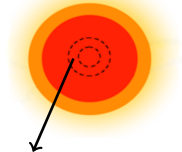
$$\Gamma_1 < \frac{4}{3}$$



3. Explosive
(oxygen)
ignition

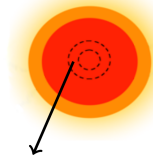


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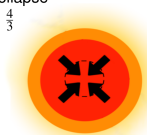


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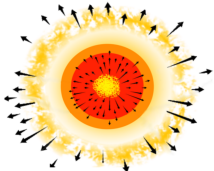
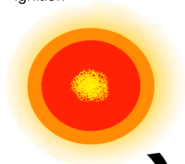
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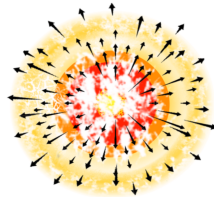
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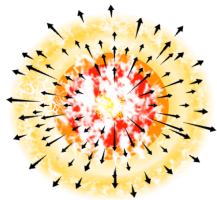
4a. Pulse with mass ejection



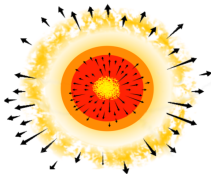
4b. PISN: complete disruption

2. Softening of EOS
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4b. PISN: complete disruption



4a. Pulse with mass ejection

7. BH



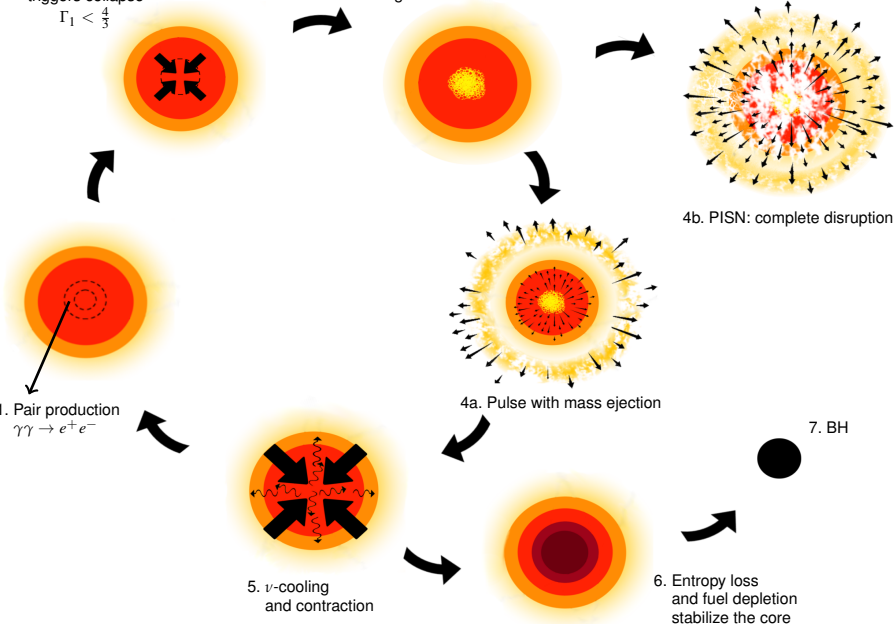
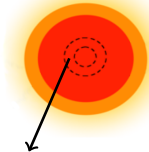
6. Entropy loss
and fuel depletion
stabilize the core

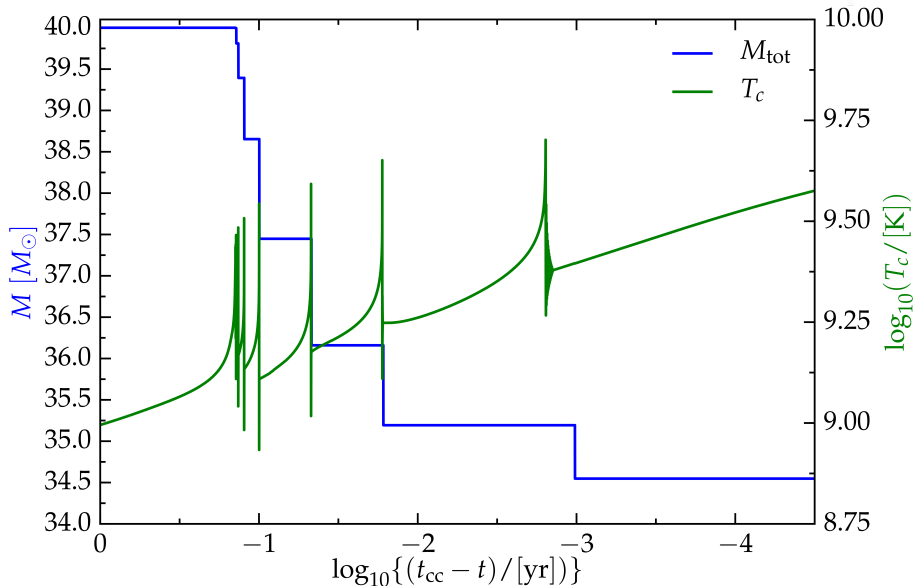


5. ν -cooling
and contraction

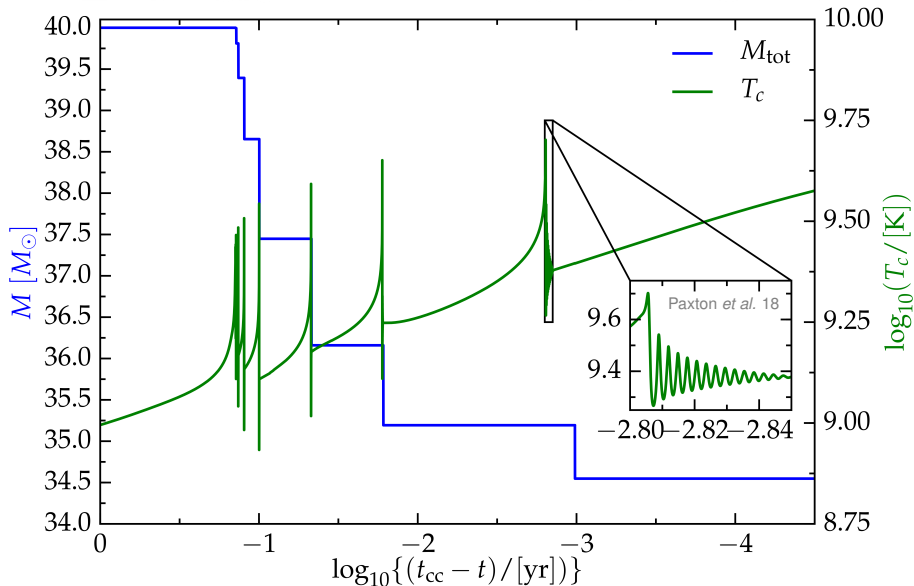


1. Pair production
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Log Time to core-collapse



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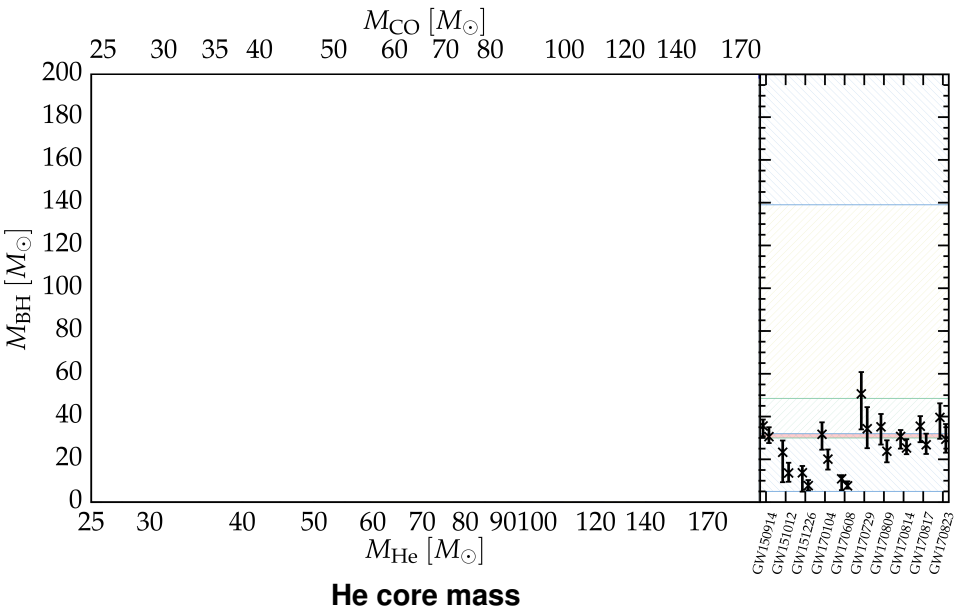
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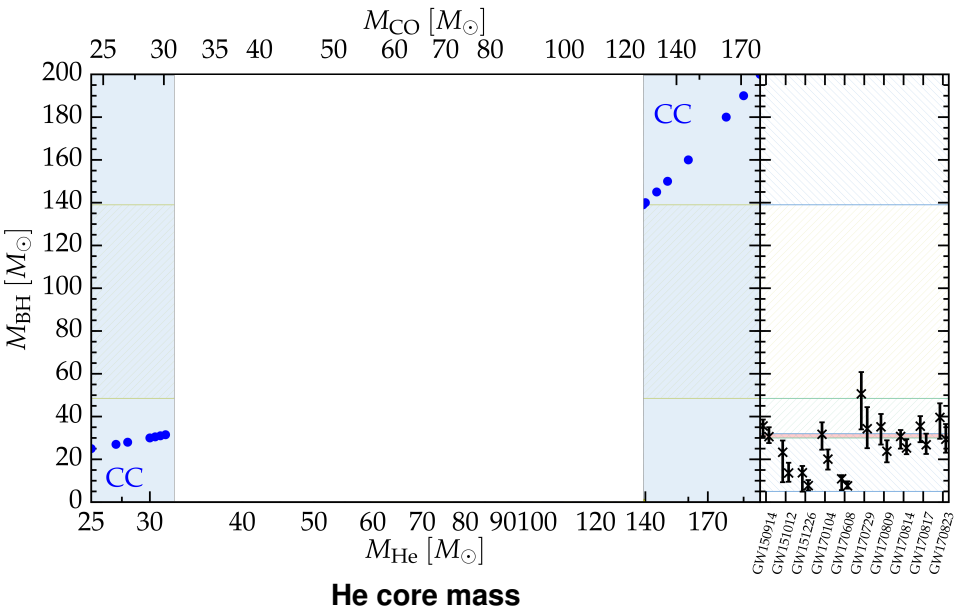
Conclusions



The origin of very massive BHs



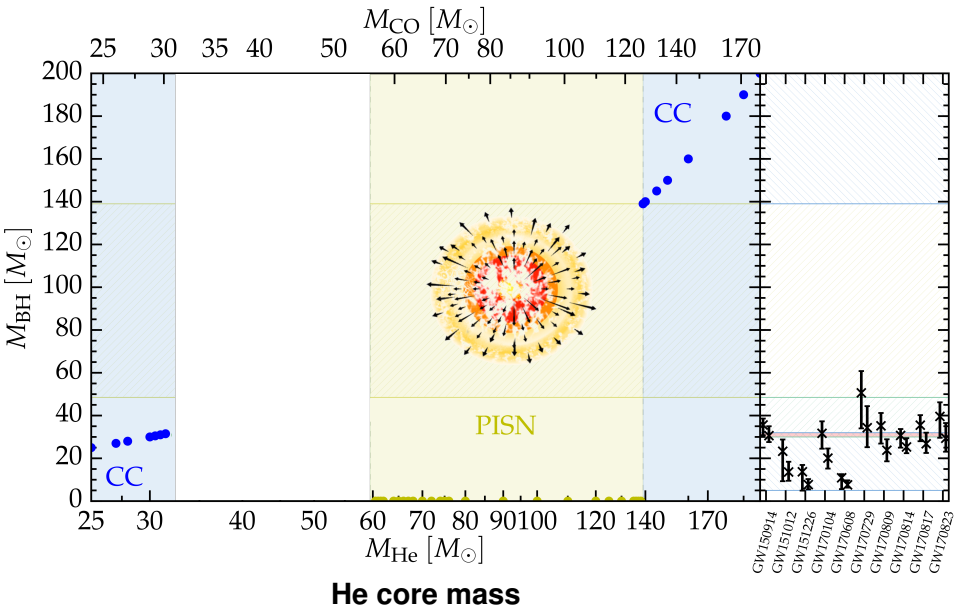
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The origin of very massive BHs



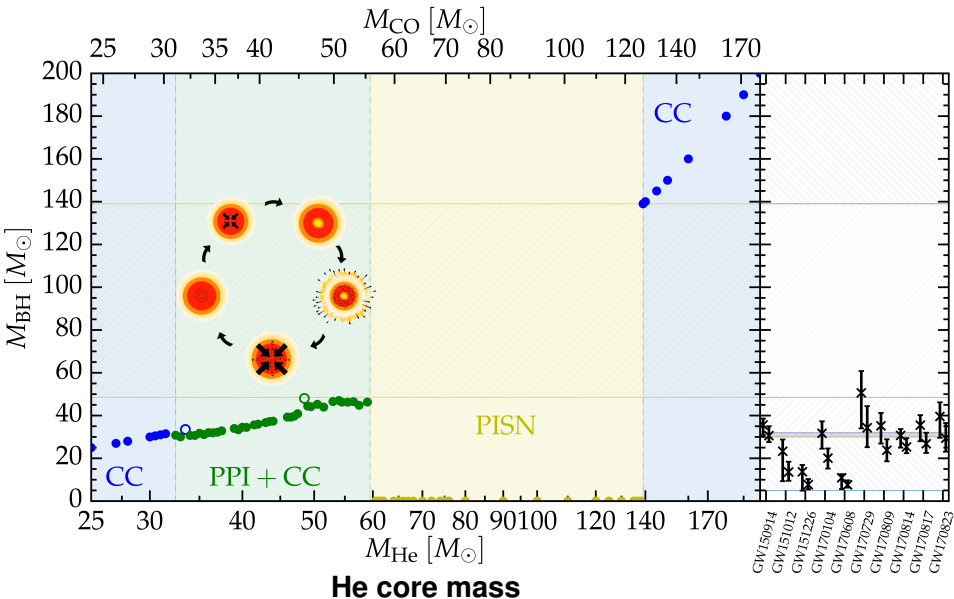
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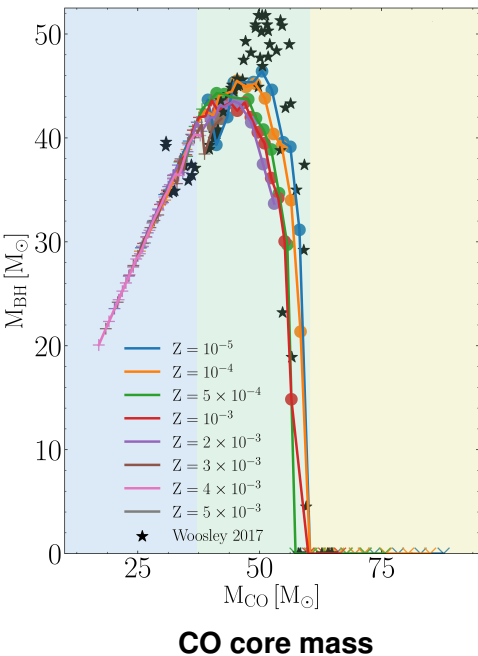


The origin of very massive BHs



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Other robustness tests:

- Spatial & temporal resolution
- Wind mass loss rate
- $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$ rate

Farmer, Renzo, *et al.* (in prep.)

Takahashi 18

Woosley 17, 19



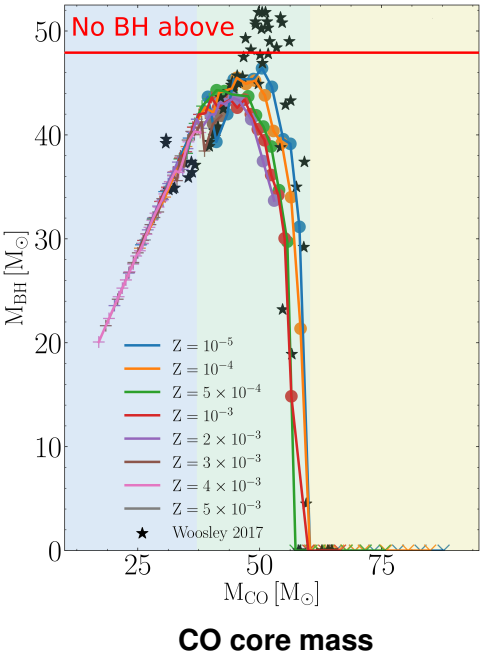
**max{BH mass} robust as
function of M_{CO}**
(rate will vary with Z)



Metallicity variations?



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Other robustness tests:

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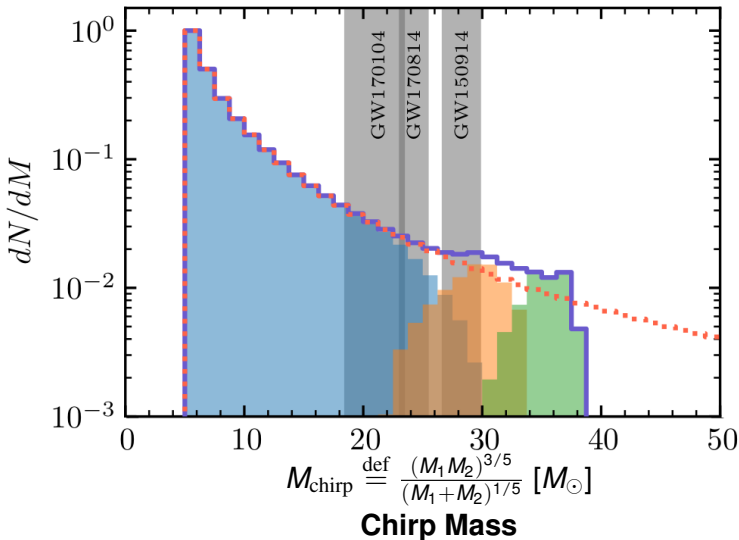
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**max{BH mass} robust as
function of M_{CO}**
(rate will vary with Z)

— with PPISNe ■ 0-PPISN ■ 2-PPISN
⋯ no PPISNe ■ 1-PPISN

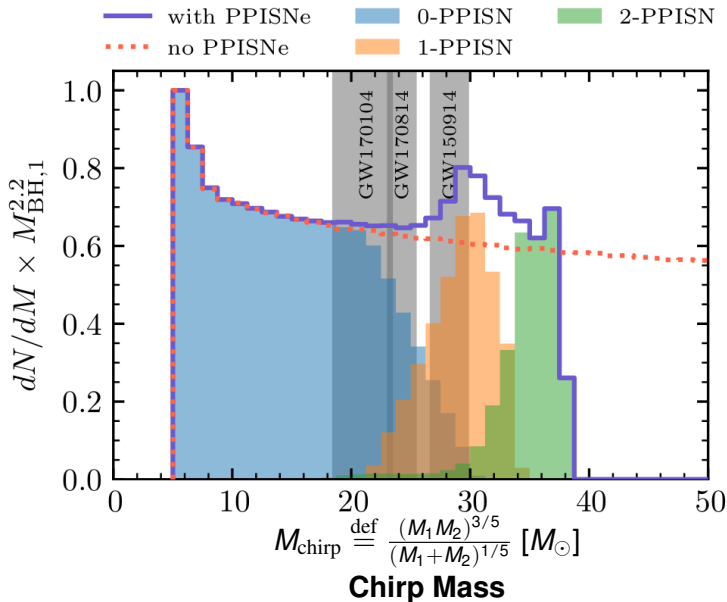


$$\frac{dN}{dM_{\text{He}}} \propto M_{\text{He}}^{-2.35}$$

$$q \stackrel{\text{def}}{=} \frac{M_2}{M_1} \geq 0.5$$

(motivated by LVC 2016)

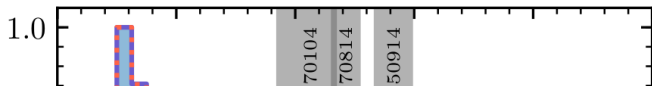
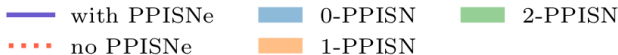
(Fishbach & Holtz 2017)



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LIGO/Virgo O3 will answer!

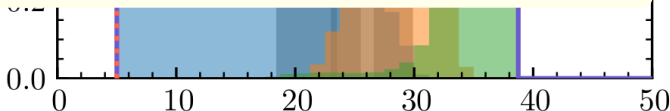
- Is there a gap?
 $\Rightarrow \mathcal{O}(10)$ binary BH detection
- Where is the lower edge of the gap?
 $\Rightarrow \mathcal{O}(100)$ binary BH detection

$$\frac{dN}{dM_{\text{He}}} \propto M_{\text{He}}^{-2.35}$$

$$q \stackrel{\text{def}}{=} \frac{M_2}{M_1} \geq 0.5$$

(motivated by LVC 2016)

(Fishbach & Holtz 2017)



$$M_{\text{chirp}} \stackrel{\text{def}}{=} \frac{(M_1 M_2)^{3/5}}{(M_1 + M_2)^{1/5}} [M_{\odot}]$$

Chirp Mass

BH or NS?

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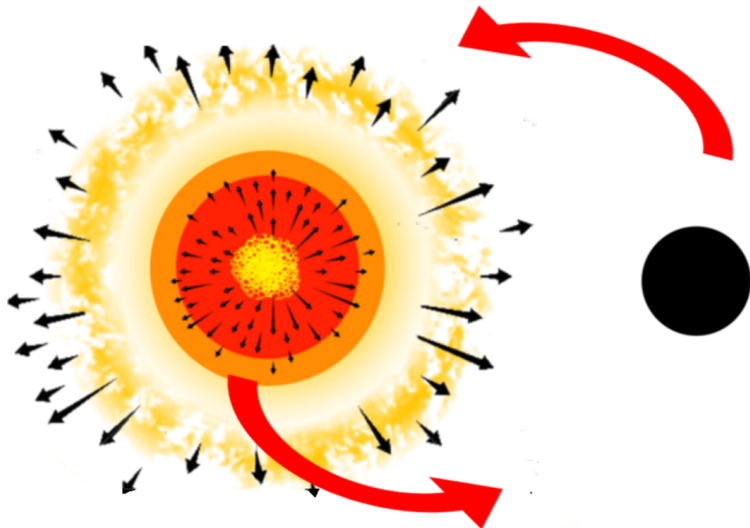
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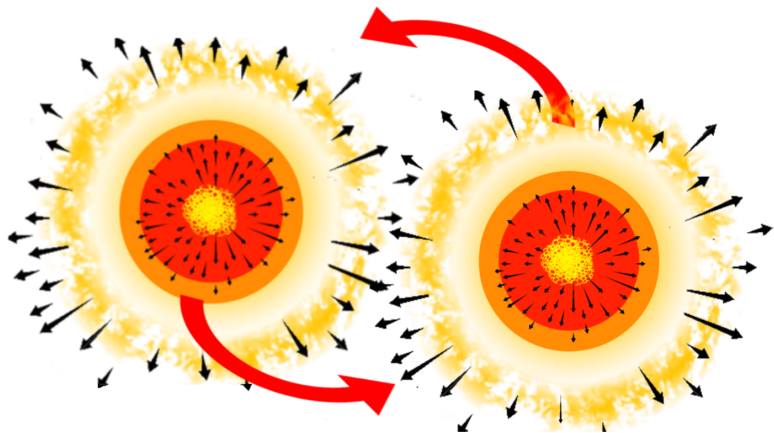
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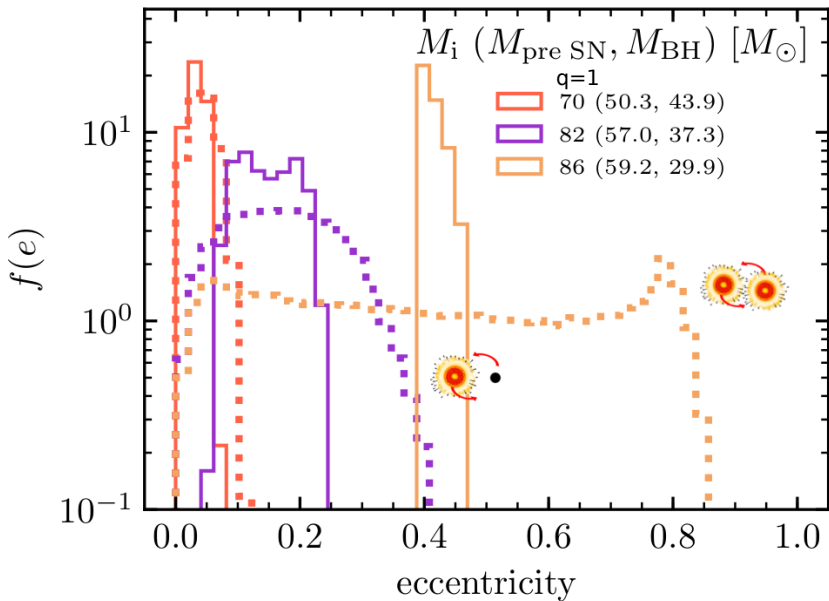


$$\Delta e = \frac{\Delta M}{M_1 + M_2 - \Delta M}$$

Two PPI in a binary



$$\Delta e = \frac{\Delta M}{M_1 + M_2 - \Delta M}$$



BH or NS?

- Single stars winds impact on the core structure

Keep the stars together

- The most common evolution for massive binaries
- Constraints on BH kicks using runaway “widow”

The most massive (stellar) BHs

- (Pulsational) pair instability
- The BH mass distribution
 - Induced eccentricity

Conclusions



- **Uncertain wind mass loss rates influence the pre-SN core**

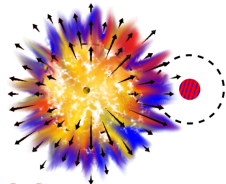
⇒ systematic bias in SN initial conditions and outcome?

- **The vast majority of binaries are disrupted**

⇒ X-ray binaries and GW sources are exceptions

- **Binarity leaves imprint on the ejected star**

- **“Widow” companions ejected constrain BH kicks**

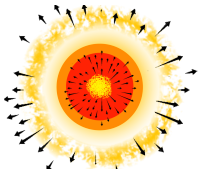


Simulations of Pulsational Pair Instability possible with **MESA**
including self-consistently dynamical evolution

- **can modify binary orbit and remnant spin**

⇒ Signature on gravitational wave signals?

- **determines BH masses below PISN gap**





- **Uncertain wind mass loss rates influence the pre-SN core**

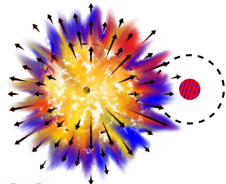
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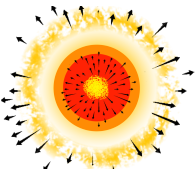


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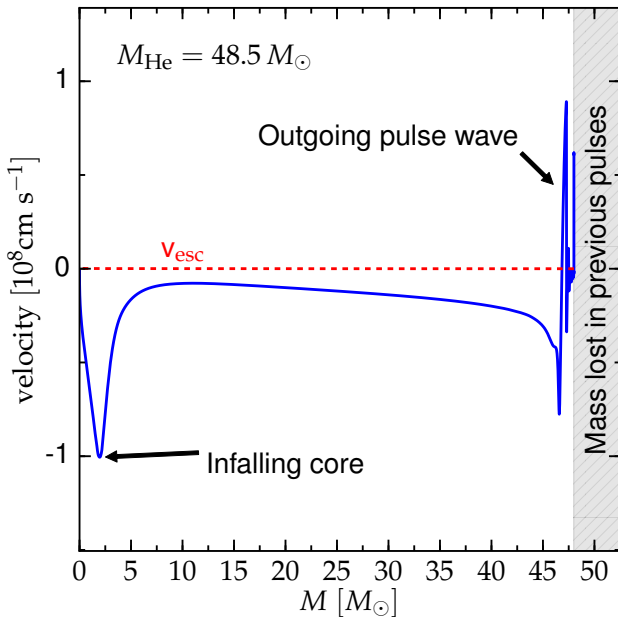
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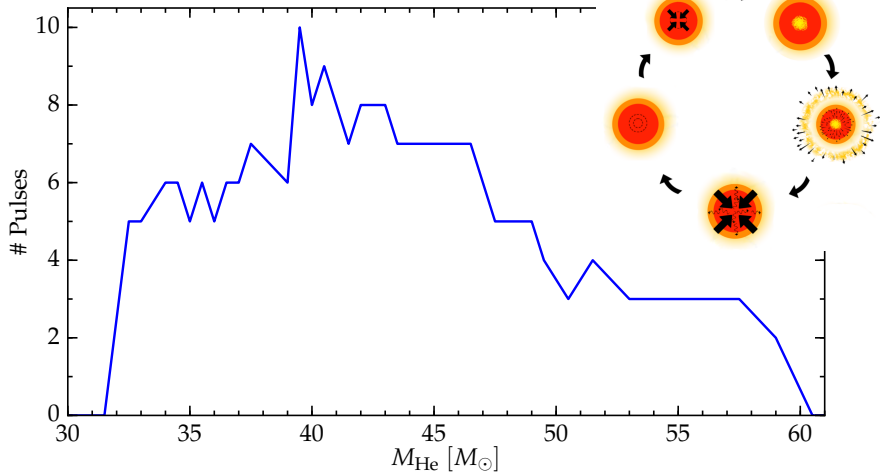
Thank you!

Backup slides



How many pulses?

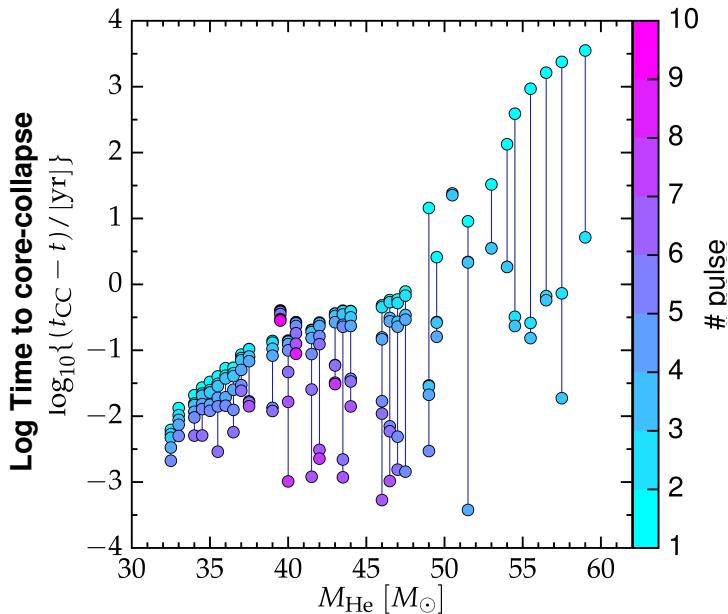
- as a function of He core mass

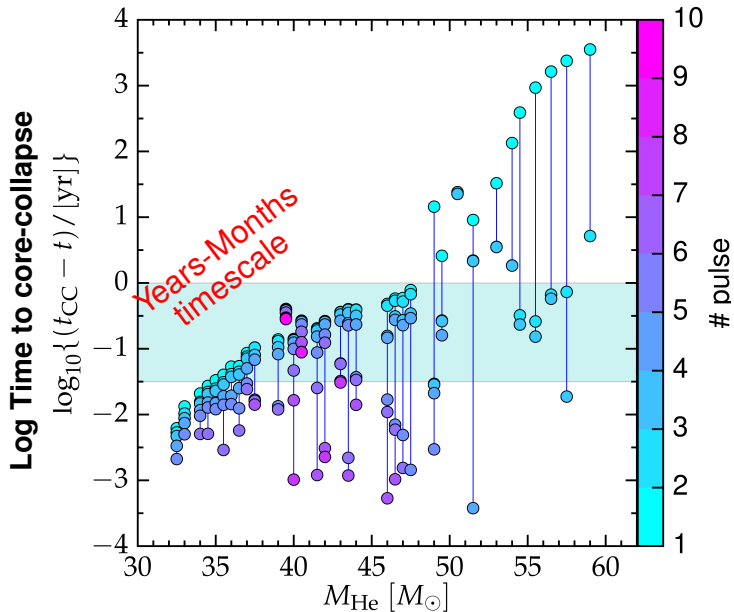


One pulse = One mass ejection

When do the pulsate?

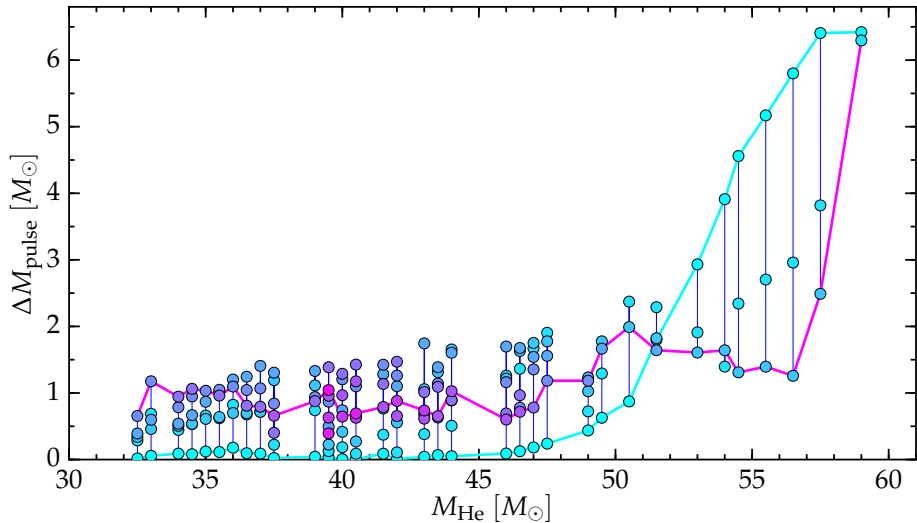
- as a function of He core mass

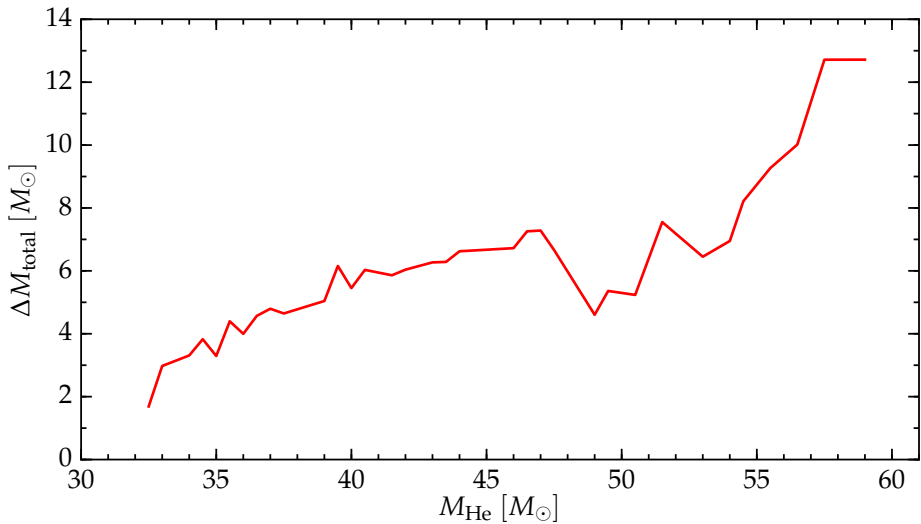




How much mass is ejected per pulse?
How much mass is ejected in total?

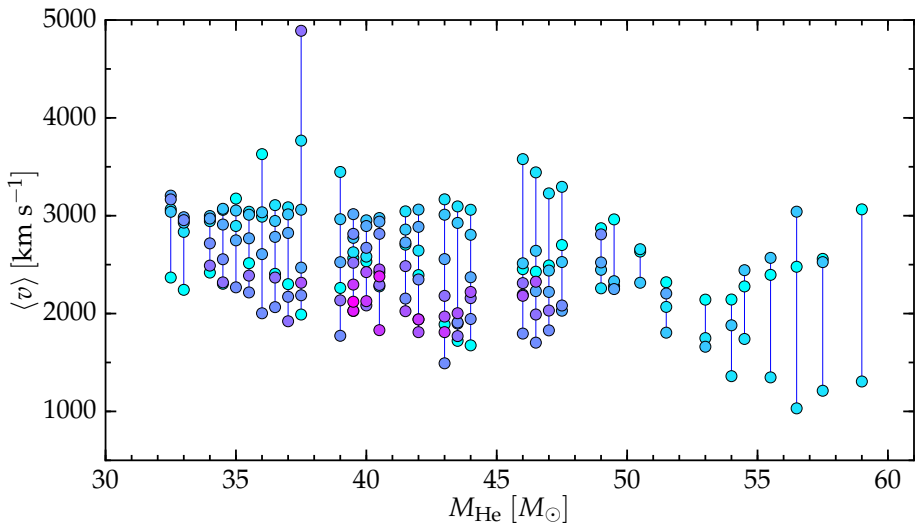
- as a function of He core mass

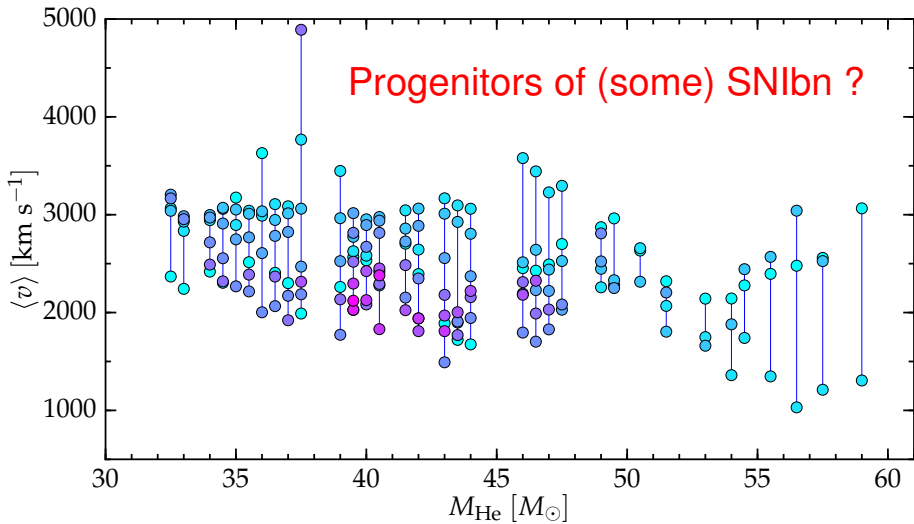




How fast are the ejected shells?

- as a function of He core mass

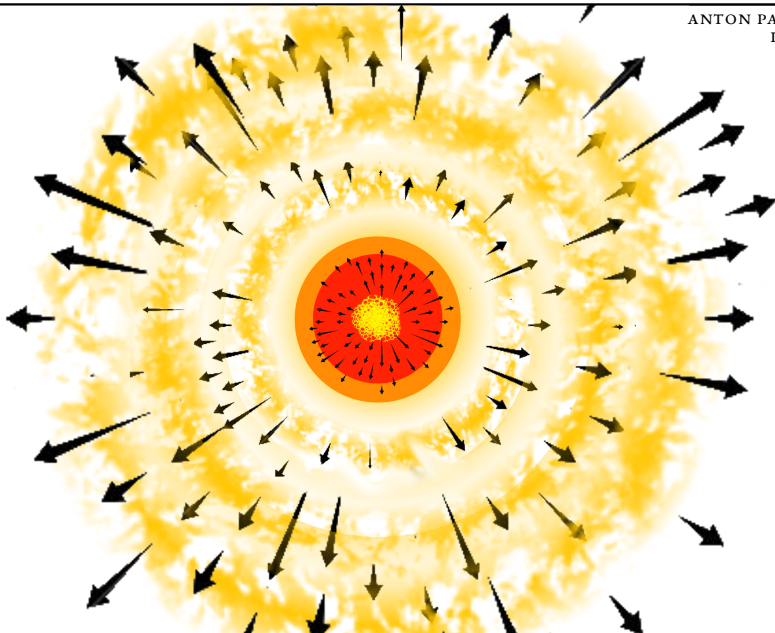




Can the mass shell collide?



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Can the mass shells collide?



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Distance to the star

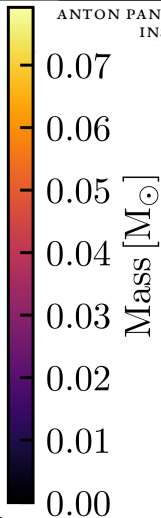
$\log_{10} R$ [cm]

16
15
14
13
12

-0.15 -0.10 -0.05 0.00

$\tau - \tau_{CC}$ [yr]

Time to core-collapse

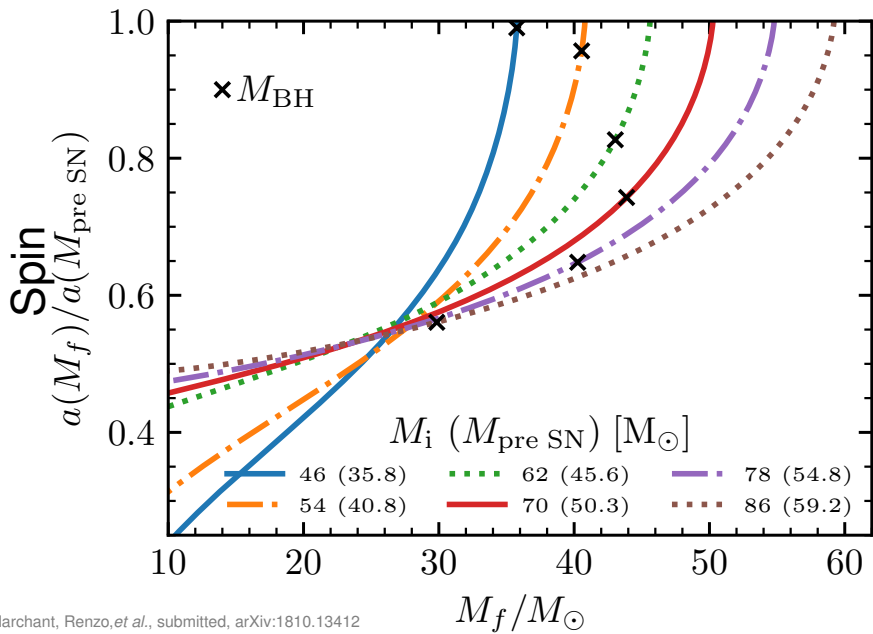


No self-interaction
or potential well
 $M_{\text{He}} = 40 M_{\odot}$

Spin down due to PPI ejecta



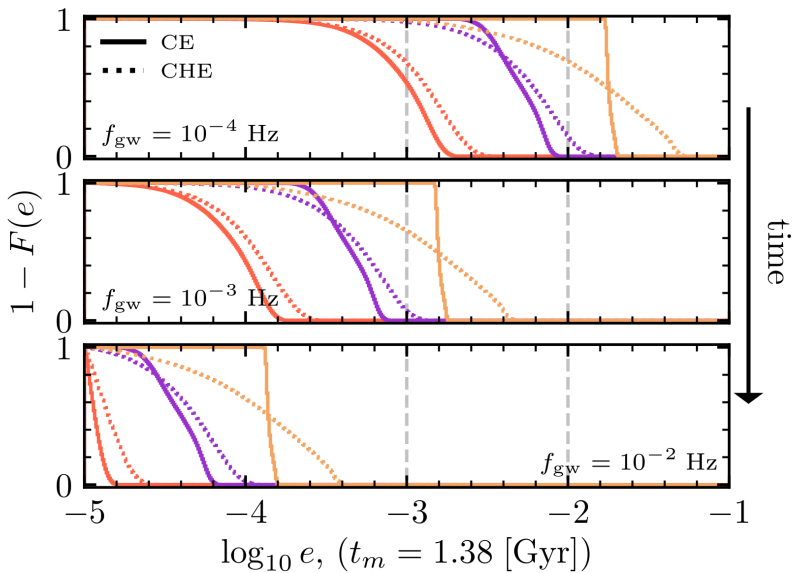
ANTON PANNEKOEK
TUE

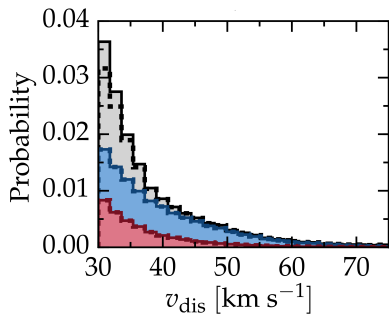


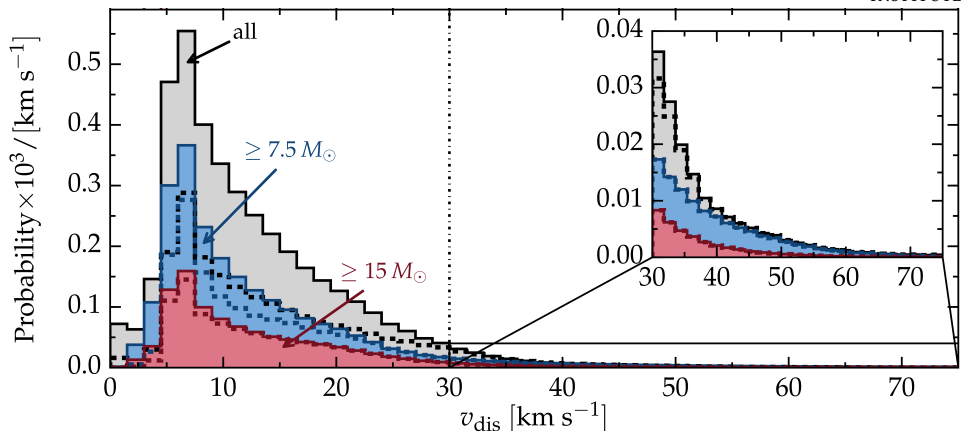
GW circularization

$M_i (M_{\text{pre SN}}, M_{\text{BH}}) [M_{\odot}]$

▭ 70 (50.3, 43.9)
 ▭ 82 (57.0, 37.3)
 ▭ 86 (59.2, 29.9)







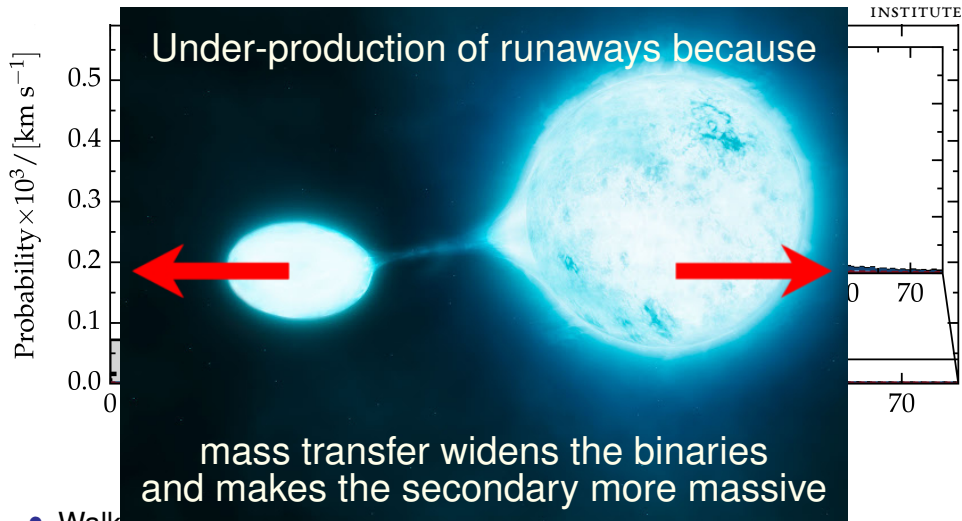
Take home points:

- Walkaways outnumber the runaways by $\sim 10\times$
- Binaries barely produce $v_{\text{dis}} \gtrsim 60 \text{ km s}^{-1}$
- All runaways from binaries are post-interaction objects

Velocity distribution: Walkaways

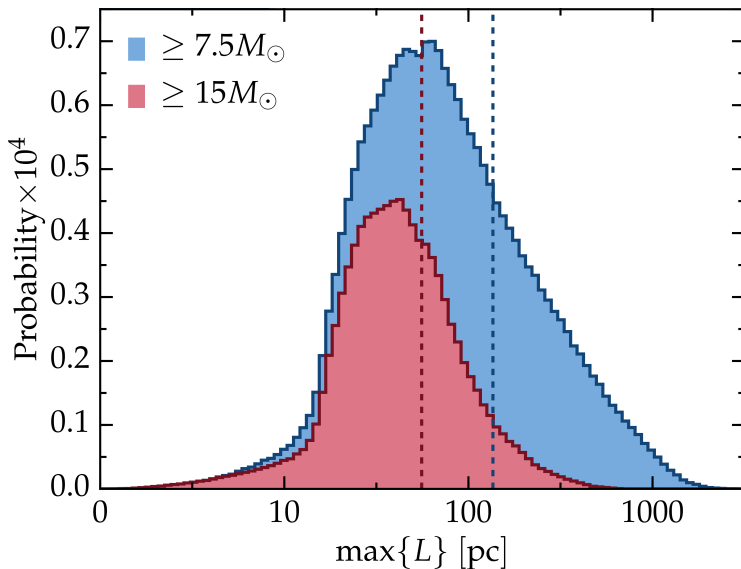


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- Walkaways outnumber the runaways by $\sim 10 \times$
- Binaries barely produce $v_{\text{dis}} \gtrsim 60 \text{ km s}^{-1}$
- All runaways from binaries are post-interaction objects

How far do they get?

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“Distance traveled”
(No potential well)