

Systematic Survey of Massive Stars Mass Loss Algorithms

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Line Driven Winds

• Highly non-linear: driving depends on $\kappa \equiv \kappa(\rho, v)$, because of Doppler shifts, and $\rho \equiv \rho(v, \dot{v})$ through hydrodynamics \Rightarrow Driving depends on outflow itself.

• Efficiency η depends on the possible presence of over-dense clumps in the atmosphere: potential overestimation of \dot{M} if interpreting the observed ρ (dominated by over-dense clumps) as the average ρ .

Too complicated to treat from first principles in Stellar Evolution Codes

Various (Semi-Empirical) Algorithmic Representations:

 $\dot{M} \equiv \dot{M} (L, T_{\text{eff}}, Z, \text{etc.}) \times \eta$

We carry out a systematic comparison of the impact of these algorithms and their efficiency scale factor on the evolution and final structure of single massive stars using the MESA stellar evolution code (Paxton et al. 2011, 2013, 2015).

Uncertainty in the Final Mass Impossible to go back in time using stellar models.

Systematic uncertainty in the final mass dominated by unknown η .



Uncertainty in the Core Structure "Explodability" depends on the chosen algorithm:

The successful explosion of a stellar model depends on the assumptions in the treatment of mass loss.

Late core structure depends on winds

- Differences between core structures arise only in late stages: core evolutionary contraction amplifies the initially small differences;
- Critical point during neon core burning (carbon shell) burning) phase for onset of significant differences;
- **•** Early (\sim main sequence) mass loss matters for the core structure.

The compactness parameter



Legend: $+ \eta = 1.0$; $\times \eta = 0.33$; • $\eta = 0.33$ Hot phase mass loss: **V** = Vink *et al.* 2000, 2001; **K** = Kudritzki *et al.* 1989; **Cool phase mass loss:** dJ = de Jager *et al.* 1988; NJ = Nieuwenhuijzen *et al.* 1990; vL = van Loon *et al.* 2005; WR mass loss: **NL** = Nugis & Lamers 2000; **H** = Hamman *et al.* 1982,1985,1998 Each algorithm is a combination of these mass loss scheme (see top axis labels)

 $\xi_{2.5} \stackrel{\text{def}}{=} \frac{2.5/M_{\odot}}{D(14)}$ $\overline{R(M_{\rm barv}=2.5M_{\odot})/1000\,\rm km}$

(O'connor & Ott 2011)

- Characterizes the density profile outside the (to-become) iron core;
- It is a function of time because of evolutionary contraction of the core;
- ► Value at onset of core-collapse can be used to predict success or failure of (neutrino driven) explosion, and nature of the compact remnant (Black Hole or Neutron Star):
 - Large values \Rightarrow hard to explode \Rightarrow probable BH; • Small values \Rightarrow easier to explode \Rightarrow probable NS;
- Value at Oxygen depletion is already a qualitative indicator of the SN outcome;
- $\xi_{2,5}$ is sensitive to mass loss algorithm after critical point of the evolution (Ne core burning / C shell burning).

Computational challenges

- Need to capture initially small effect \Rightarrow need very high spatial resolution: 20000 – 100000 mesh points until oxygen depletion
- Need detailed core structure \Rightarrow large number of isotopes to trace weak reactions determining Y_e : 45 isotopes until O depletion, 203 afterwards.

Hard to compute very late stages until onset of core collapse (extremely slow runs)^{*}.



Implications

- Mapping of M_{ZAMS} to remnant depends on the assumed mass loss algorithm and efficiency factor η .
- Initial condition for SN explosion simulations might be biased depending on the wind scheme.
- "Explodability windows" in mass might shift, with potentially large implications for the population of black holes and neutron stars.

Possible ways out:

- Observational constraints on wind mass loss rates using colliding wind in binaries, Be X-ray binaries, SN shock interaction with circumstellar material, and circumstellar material chemical composition, etc.
- Quantification of the systematic uncertainty due to the treatment of winds using all the algorithms available for each initial condition adopted.